



Particle detection

with semiconductor detectors

Pre-work

- Watch: what is a semiconductor?
<https://www.youtube.com/watch?v=gUmDVe6C-BU>
- Watch: how does a transistor work?
<https://www.youtube.com/watch?v=lcrBqCFLHIY>
- Read 10.1.4 (mobility) and 10.3 (the np semiconductor junction and depletion depth) in [Techniques for nuclear and particle physics experiments](#)
- Submit a picture of band gaps in semiconductors, insulators, and conductors.
- Answer the questions in the quiz

By the end of today you will be able to...

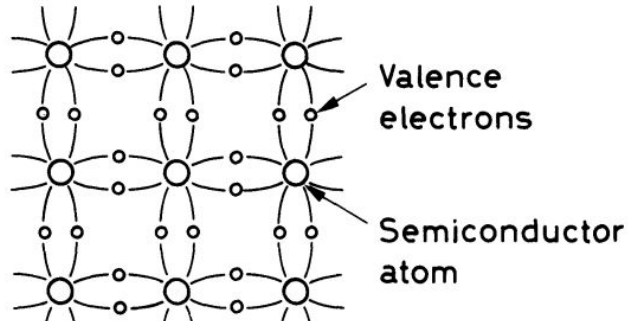
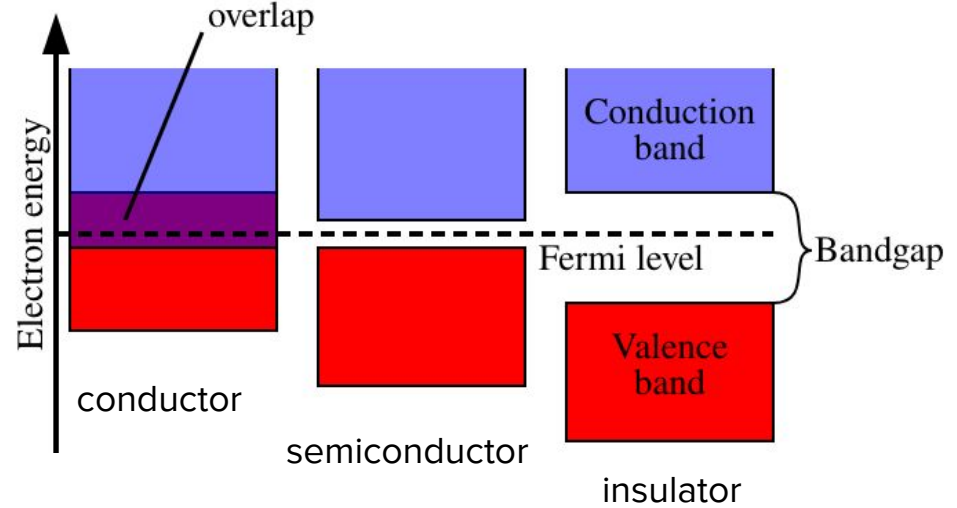
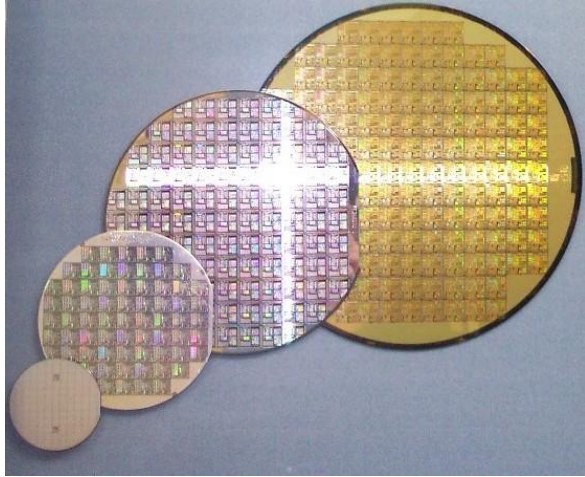
- Explain the difference between a semiconductor, insulator, and conductor
- Explain how a semiconductor can be used to detect particles
- Explain signal creation, detection, propagation, and processing in semiconductor detectors
- Argue for different detector technologies in different applications.

What will not be fully covered until later in the course:

- How to use particle detectors to construct particle tracks
- How to use particle detectors for particle identification

What is a semiconductor?

Semiconductors

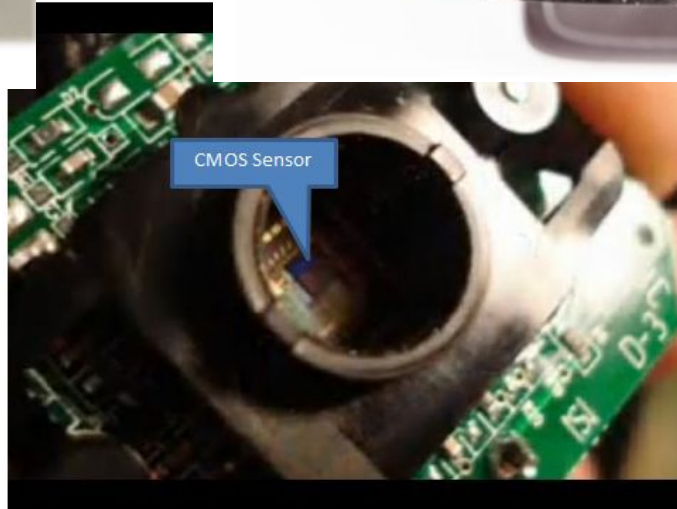
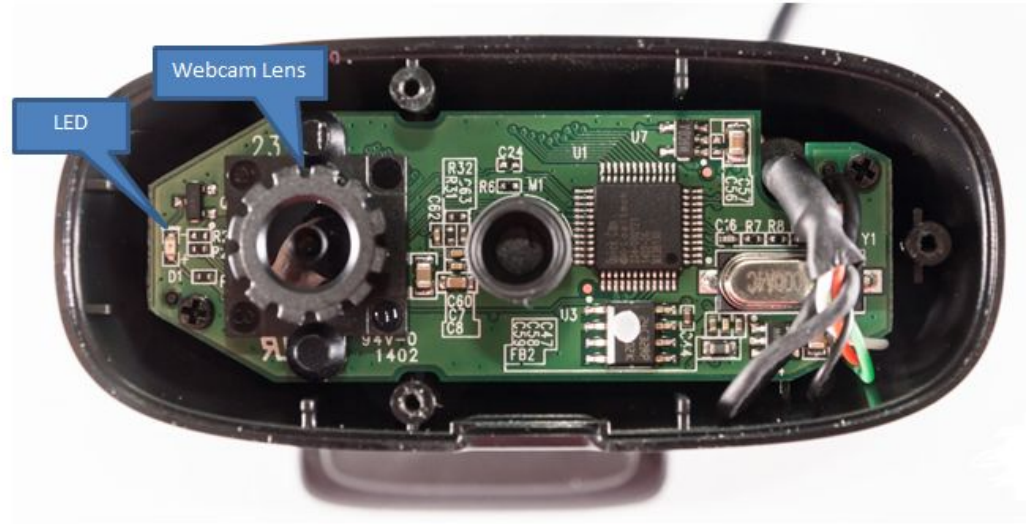


Solid state material: atomic energy levels merge to energy bands

In your webcam

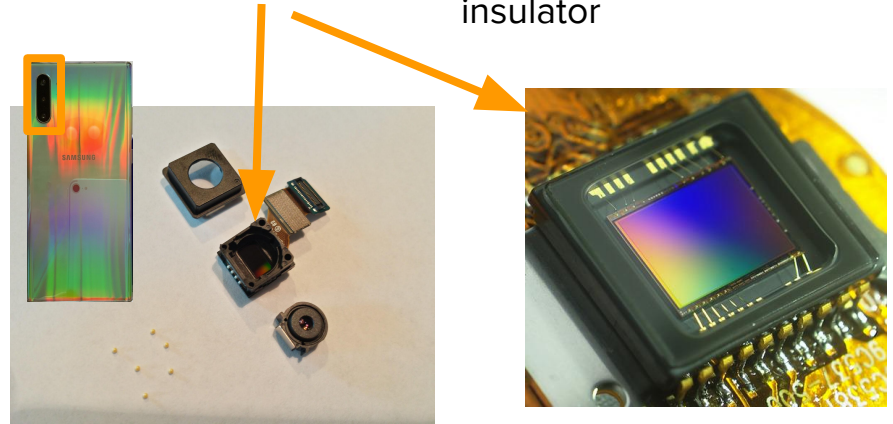
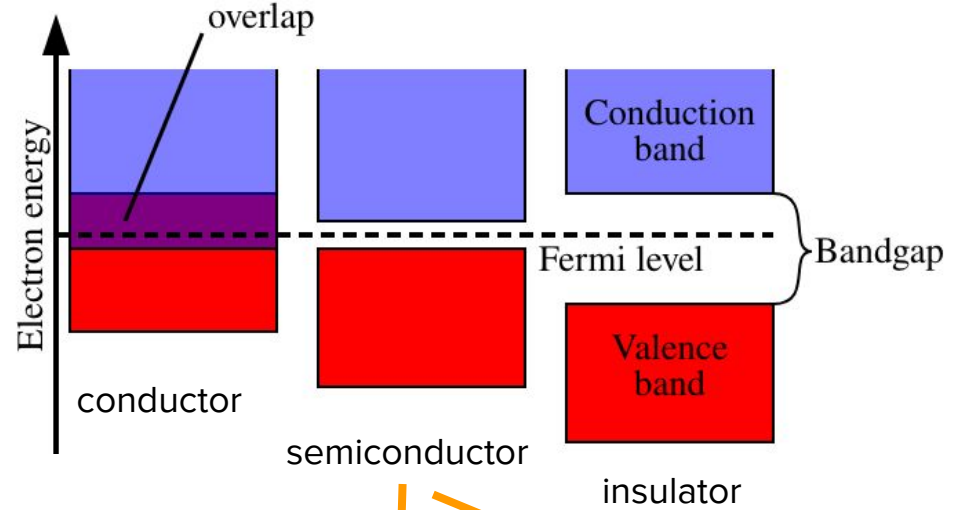
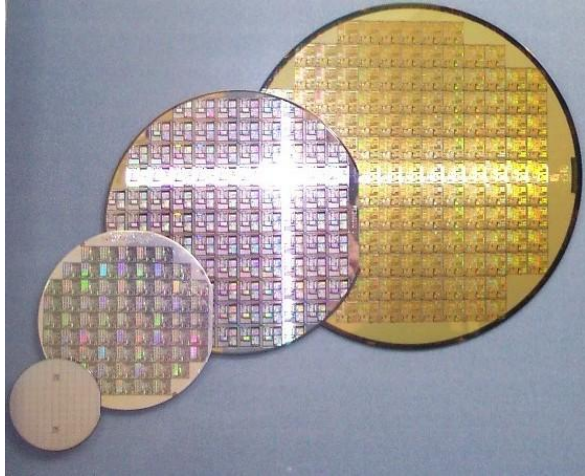


Original camera



... you have a particle detector!

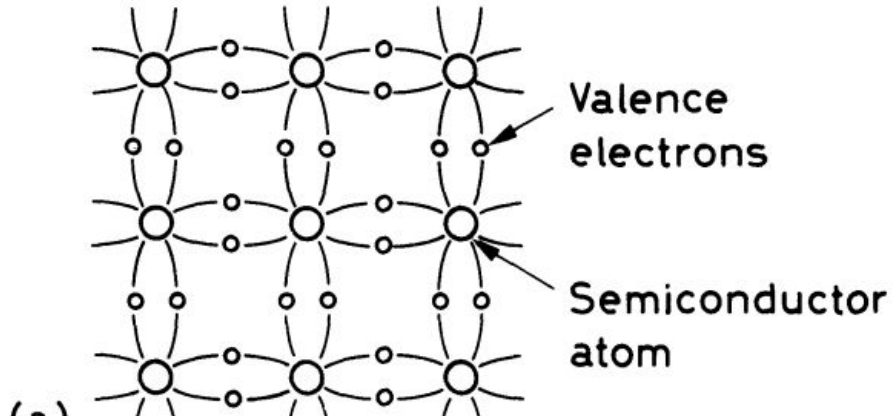
Semiconductors



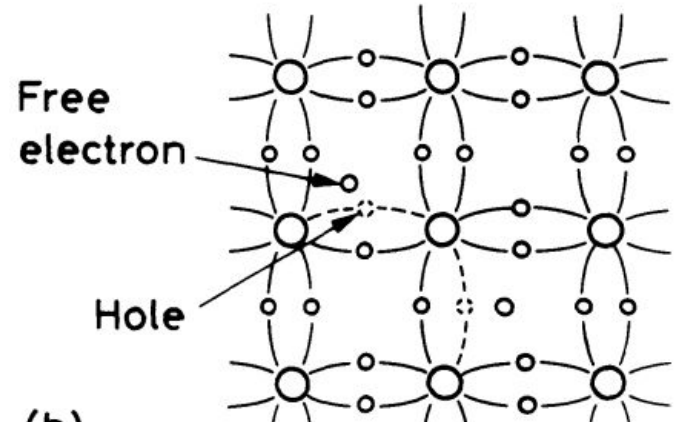
A semiconductor device is easy to manipulate, easily controlled, and makes an excellent sensor

What happened when silicon was heated in the video?

Silicon bonds



Silicon at 0 K: all electrons are bonded



Silicon at higher temperature: through thermal excitations, some electrons are excited and holes are left behind.

Solid state physics: semiconductor statistics

Fermi-Dirac

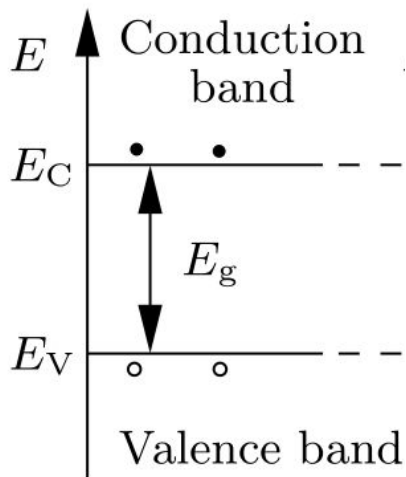
Fermi-Dirac

$$f_e(E) = \frac{1}{e^{(E-E_F)/kT} + 1}$$

Probability that an electron state in the conduction band is filled

$$f_e(E) = \frac{1}{e^{(E-E_F)/kT} + 1}$$

Intrinsic semiconductor



Intrinsic semiconductor

Concentration of acceptor/donor impurities N_A/N_D or N_C/N_V conduction/valence band edge density of states

No doping yet!

Excited electrons n

Number of holes p

$$np = n_i^2$$

\ll

thermally generated eh pairs

Carrier concentration

$$n_i = N_c e^{-(E_c - E_F)/kT} = N_v e^{-(E_F - E_v)/kT}$$

N_c : effective density of states at conduction band edge

N_v : effective density of states at valence band edge

Excited electrons n

Number of holes p

Intrinsic carrier concentration

$$n_i = N_c e^{-(E_c - E_F)/kT} = N_v e^{-(E_F - E_v)/kT}$$

The product of electron
and hole concentrations

Band gap energy in silicon $E_g = 1.12$ eV

$$n_i = 1.45 \times 10^{10} \text{ cm}^{-3}$$

depends
only on the
gap
energy

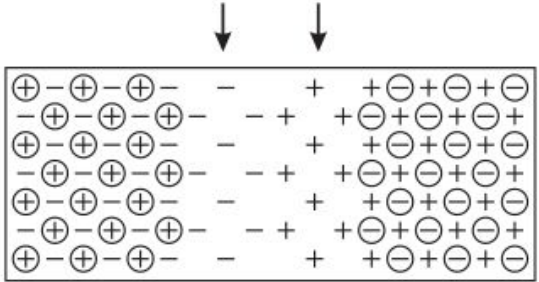


$$np = n_i^2 = N_c N_v e^{-(E_c - E_v)/kT} = N_c N_v e^{-E_g/kT}$$

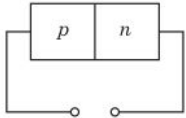
Depletion zone in intrinsic semiconductor

Electrons and holes diffuse through the junction.

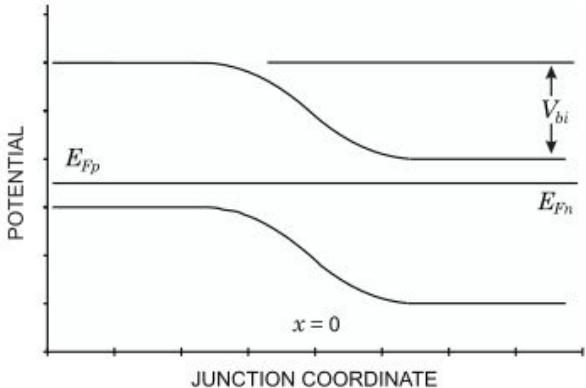
FIXED CHARGE OF ATOMIC CORES



THERMAL EQUILIBRIUM



No doping yet!



This results in a potential difference called the built-in voltage between the p and n regions

A semiconductor device is easy to
manipulate

Charge carriers: jargon alert

Now, we dope! =)

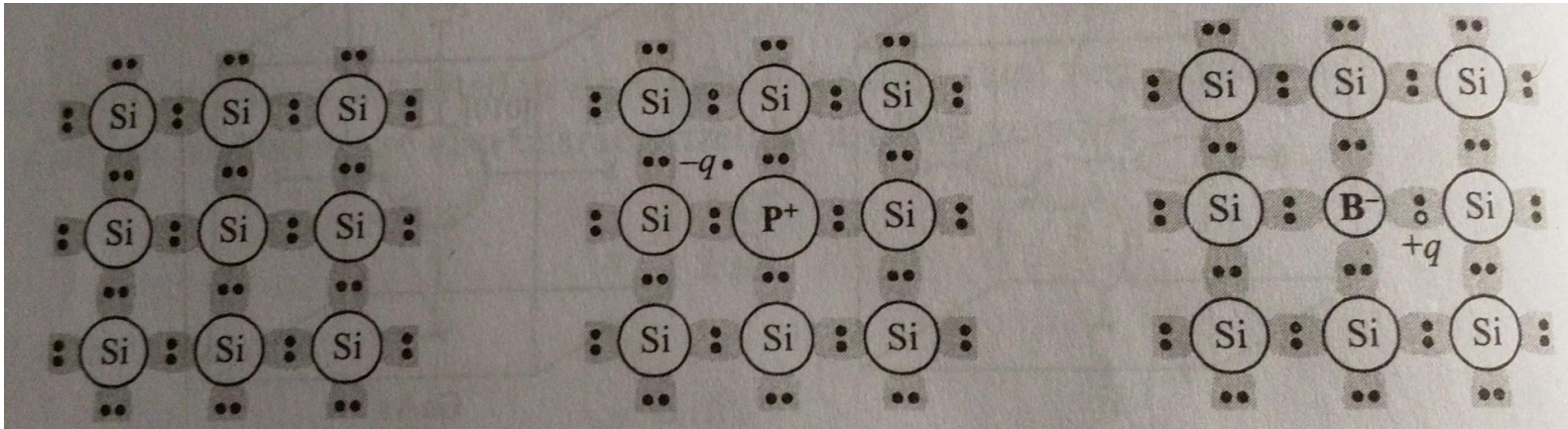
why?

Intrinsic or **pure** or **undoped** semiconductors:

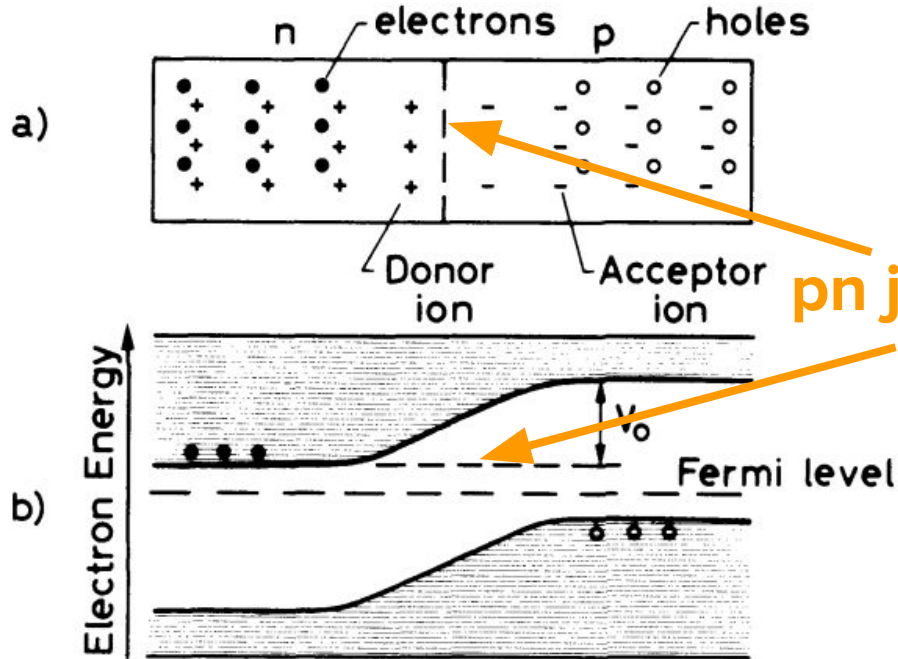
Equal amount of e^- and h carriers

N-type semiconductors are doped with a **donor impurity** that results in **electrons** becoming the **majority carrier** and holes the minority carrier.

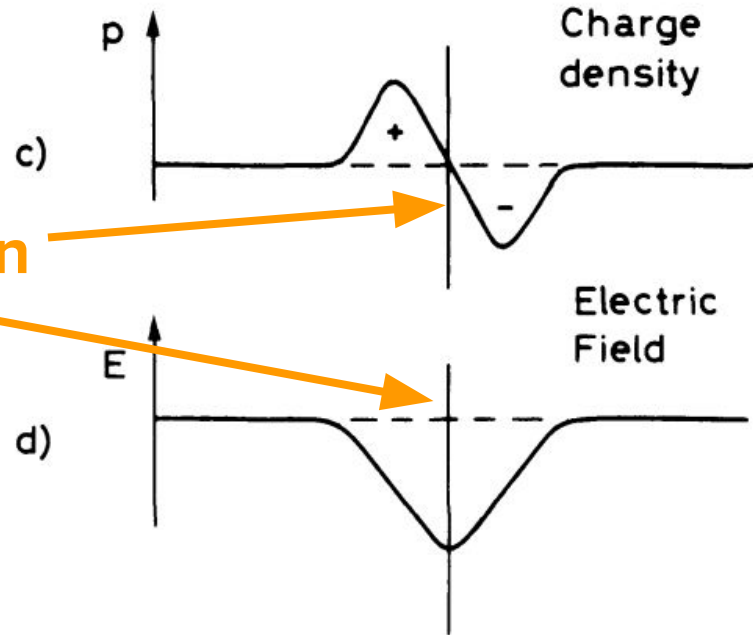
P-type semiconductors are doped with an **acceptor impurity** that results in **holes** becoming the **majority carrier** and electrons the minority carrier



join a p-doped and n-doped semiconductor



pn junction



Dope that stuff! Improve performance

Net carrier
concentration

$$\Delta n = n - p = N_d^+ - N_a^-$$

$$E_F - E_i = -k_B T \log \frac{N_a - N_d}{n_i}$$

Doping concentration N : commonly expressed in terms of the measurable **resistivity** of the material.

$$\rho = \frac{1}{\mu e N}$$

What was μ ?

Dope that stuff! Improve performance

Net carrier
concentration

$$\Delta n = n - p = N_d^+ - N_a^-$$

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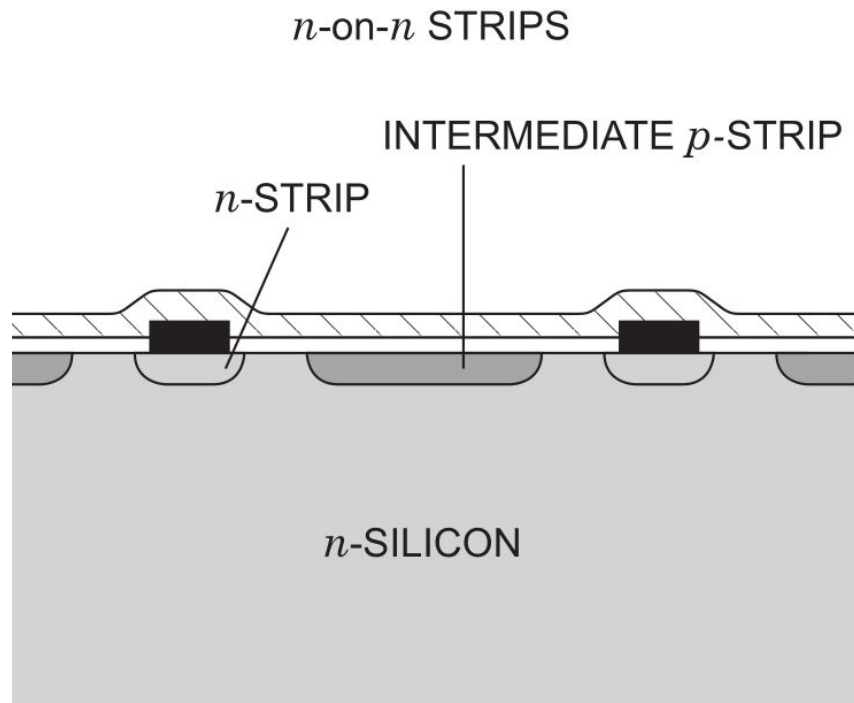
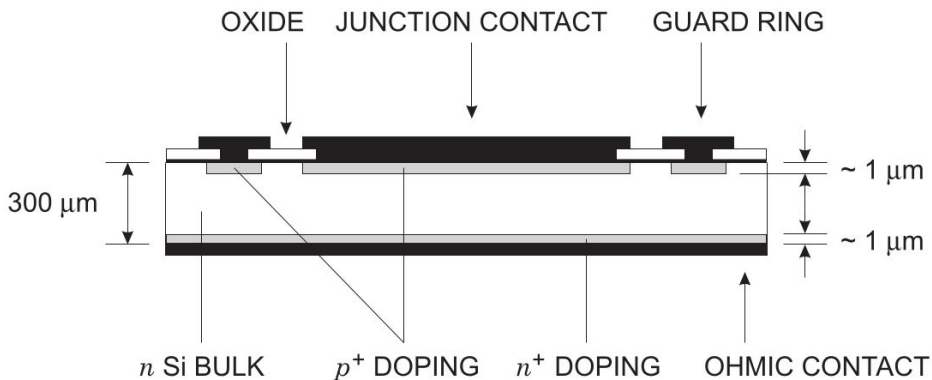
What was μ ?

$$\vec{v} = \mu \vec{E}$$

Example silicon diode and strip detector

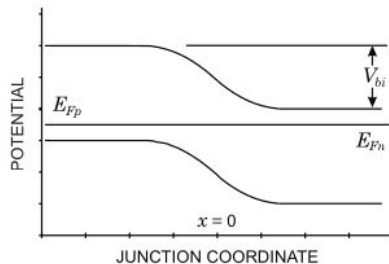
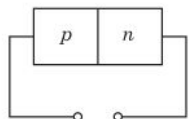
Highly doped surface layer in lightly doped “bulk”

Ohmic contact = back electrode

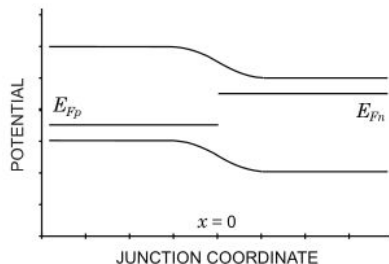
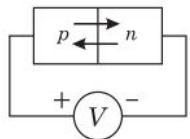


Fermi level depends on doping concentrations

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FORWARD BIAS



$$\frac{E_c - E_F}{k_B T} = \log \left(\frac{N_c}{N_d - N_a} \right)$$



$$N_d \gg N_a,$$



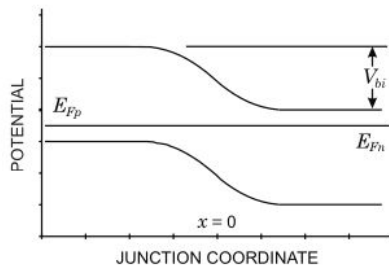
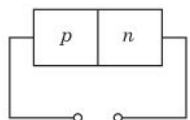
$$\text{small } E_c - E_F$$



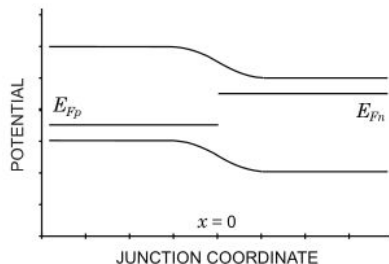
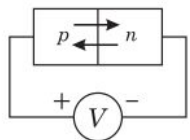
Fermi level close to conduction band edge

Fermi level depends on doping concentrations

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FORWARD BIAS



$$\frac{E_c - E_F}{k_B T} = \log \left(\frac{N_c}{N_d - N_a} \right)$$

$$N_d \gg N_a,$$

$$\text{small } E_c - E_F$$

Fermi level close to conduction band edge

What happened at the pn junction in the video?
Can we detect particles like that?

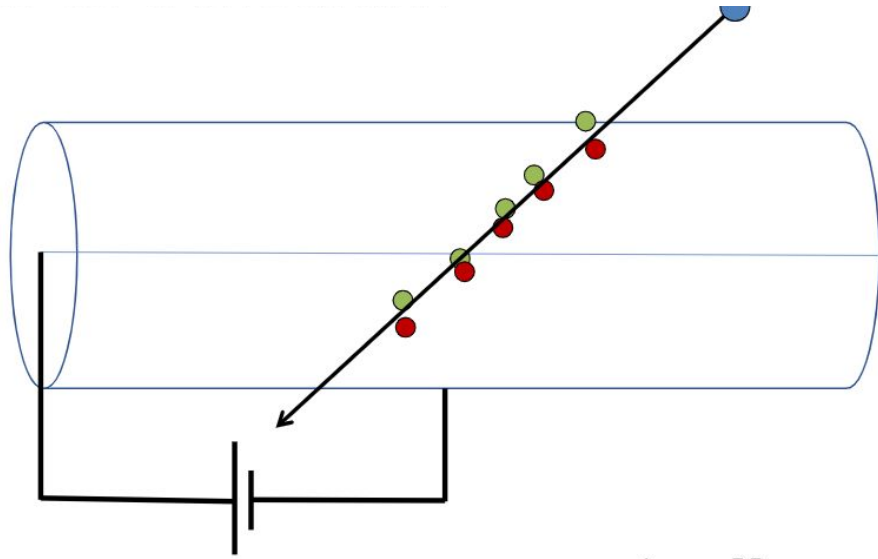
A signal in a gas detector

Last week:



A signal in a gas detector

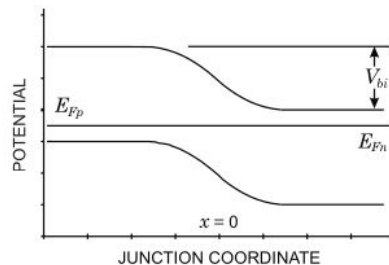
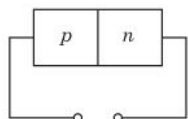
Last week:



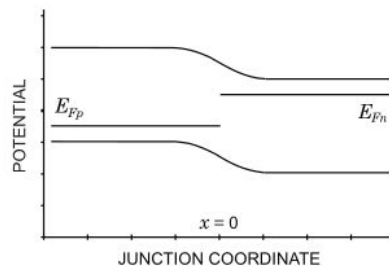
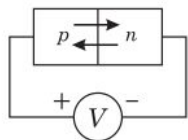
$$E = \frac{1}{s} \frac{V}{\ln(b/a)}$$

Adjust the pn diode to particle detection mode

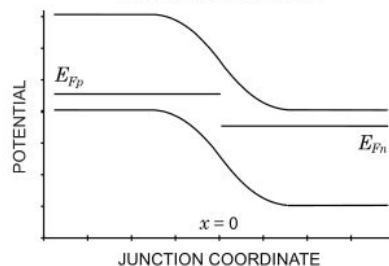
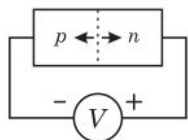
THERMAL EQUILIBRIUM



FORWARD BIAS



REVERSE BIAS



Increase depletion zone by applying a reverse bias

Particle detection mode!

Reverse bias

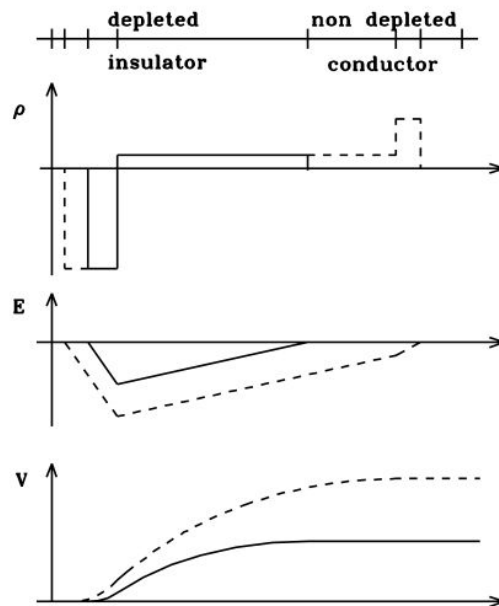
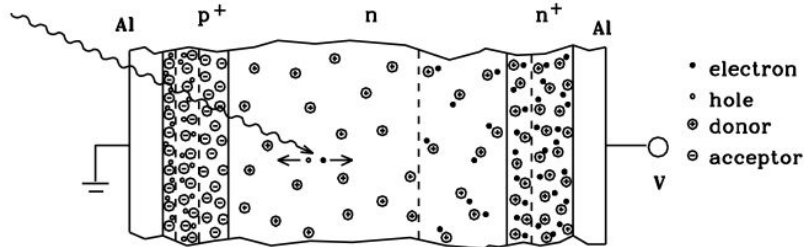
Ideal for particle detection!

Suppress current across junction

Incoming particles **ionize** the atoms creating electron-hole pairs

e⁻h⁺ pairs travel in electric field

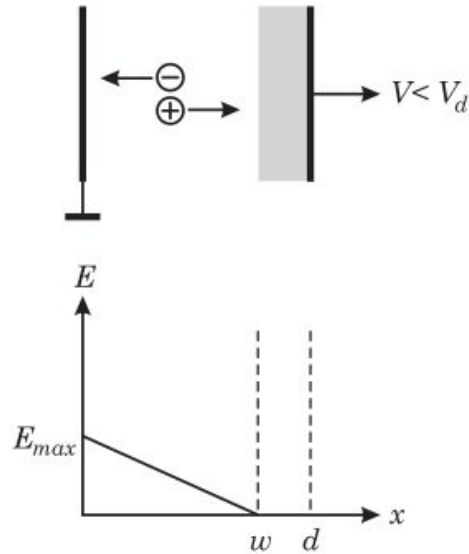
Charge **induced** in electrodes



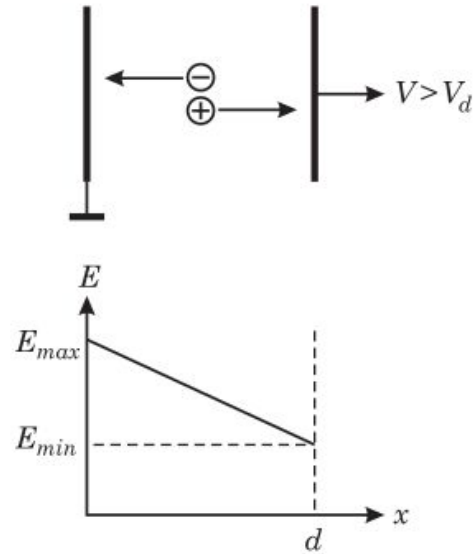
Remember Poisson:

$$\frac{d^2 V}{dx^2} = -\frac{\rho(x)}{\epsilon}$$

Full vs partial depletion



Partial depletion



Overdepletion

Depletion zone

$$d = x_n + x_p = \left(\frac{2 \epsilon V_0 (N_A + N_D)}{e N_A N_D} \right)^{1/2}$$

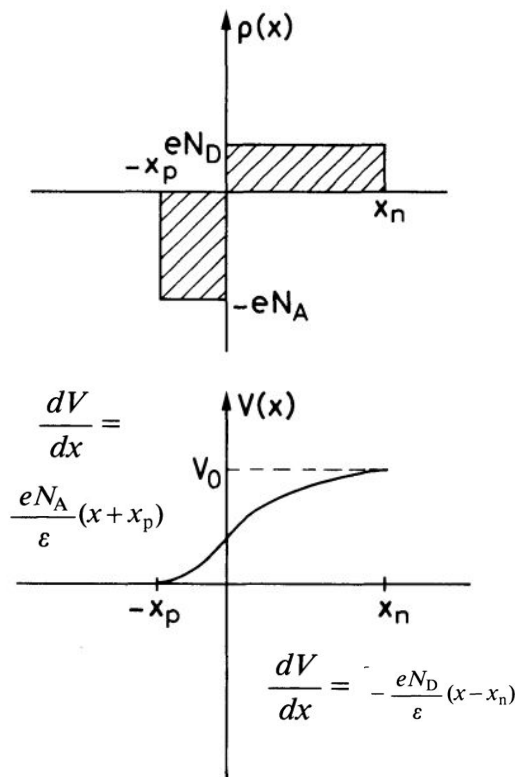
$$N_A \gg N_D$$

d = depletion depth

$$d \approx x_n \approx \left(\frac{2 \epsilon V_0}{e N_D} \right)^{1/2}$$

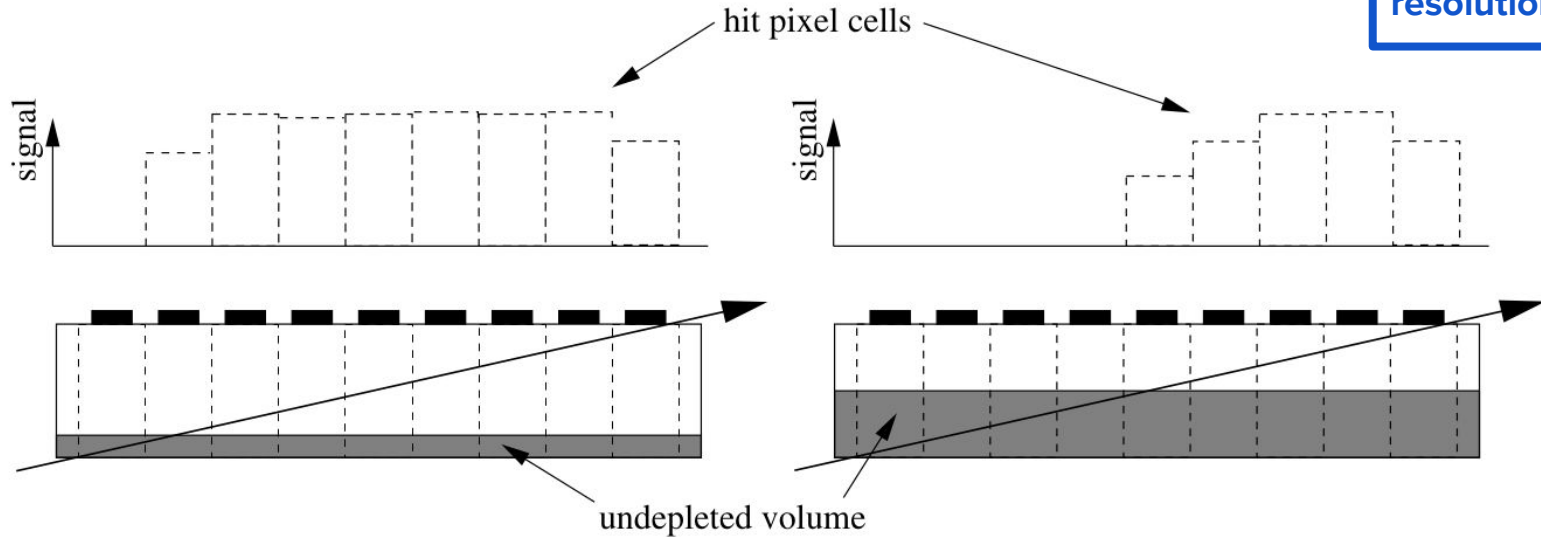
We want a large depletion zone! Why?

$$\rho(x) = \begin{cases} eN_D & 0 < x < x_n \\ -eN_A & -x_p < x < 0 \end{cases},$$



Depletion width: the higher, the more signal

Silicon has an excellent position resolution



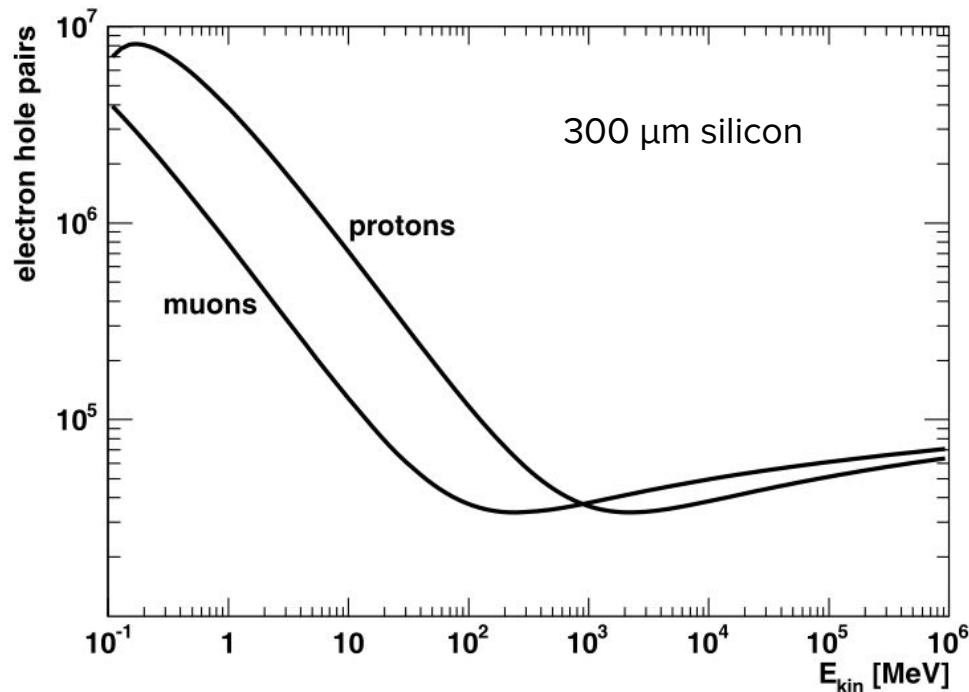
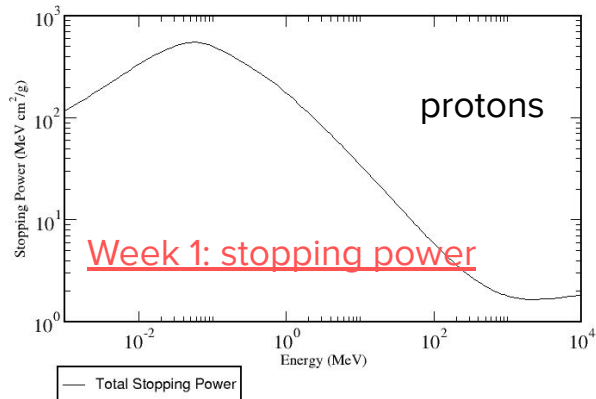
Why does the particle come in at such a shallow angle? Hint: it is on purpose.

dE/dx: electron hole pairs created in silicon

Silicon density at ~ 293 K:
 2.3290 g/cm^3

Silicon mean ionization
energy: 3.6 eV

SILICON

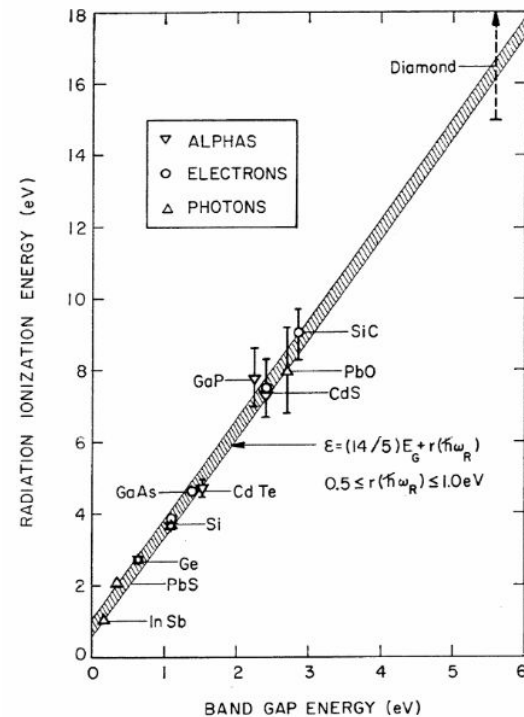


23000 e⁻ in 300 μm silicon \rightarrow 2.76 MeV / cm

From <https://physics.nist.gov>

Band gap vs mean ionization energy

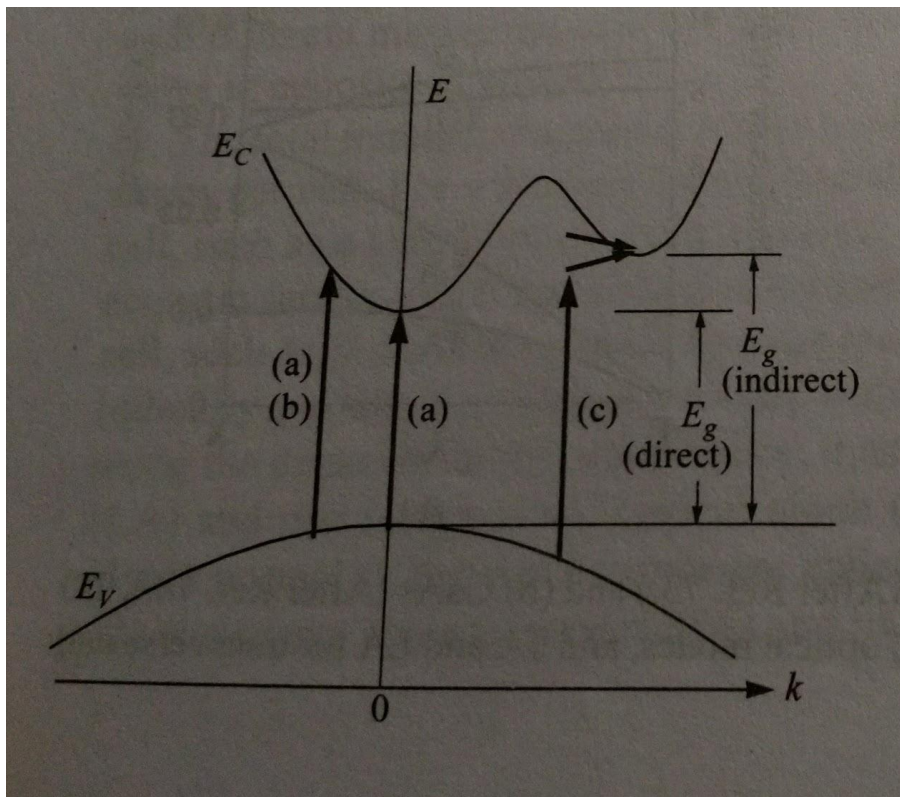
What is the silicon band gap energy?
The ionization energy?
Why?



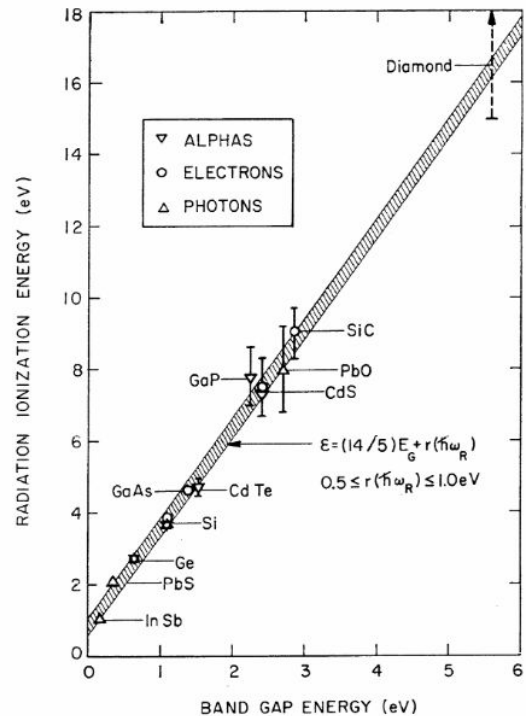
C.A. Klein, J. Applied Physics 39 (1968) 2029

Helmut Spieler
LBNL

Band gap vs mean ionization energy



What is the silicon band gap energy? The ionization energy? Why?

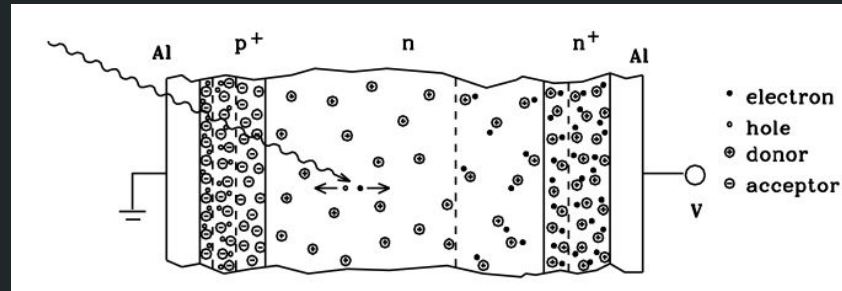


C.A. Klein, J. Applied Physics 39 (1968) 2029

Helmut Spieler
LBNL

A semiconductor device is easy to manipulate, easily controlled, and makes an excellent sensor

Apply a reverse bias to the diode with pn junction to deplete the sensor of free carriers. The larger this zone, the more sensitive volume for detecting particles.



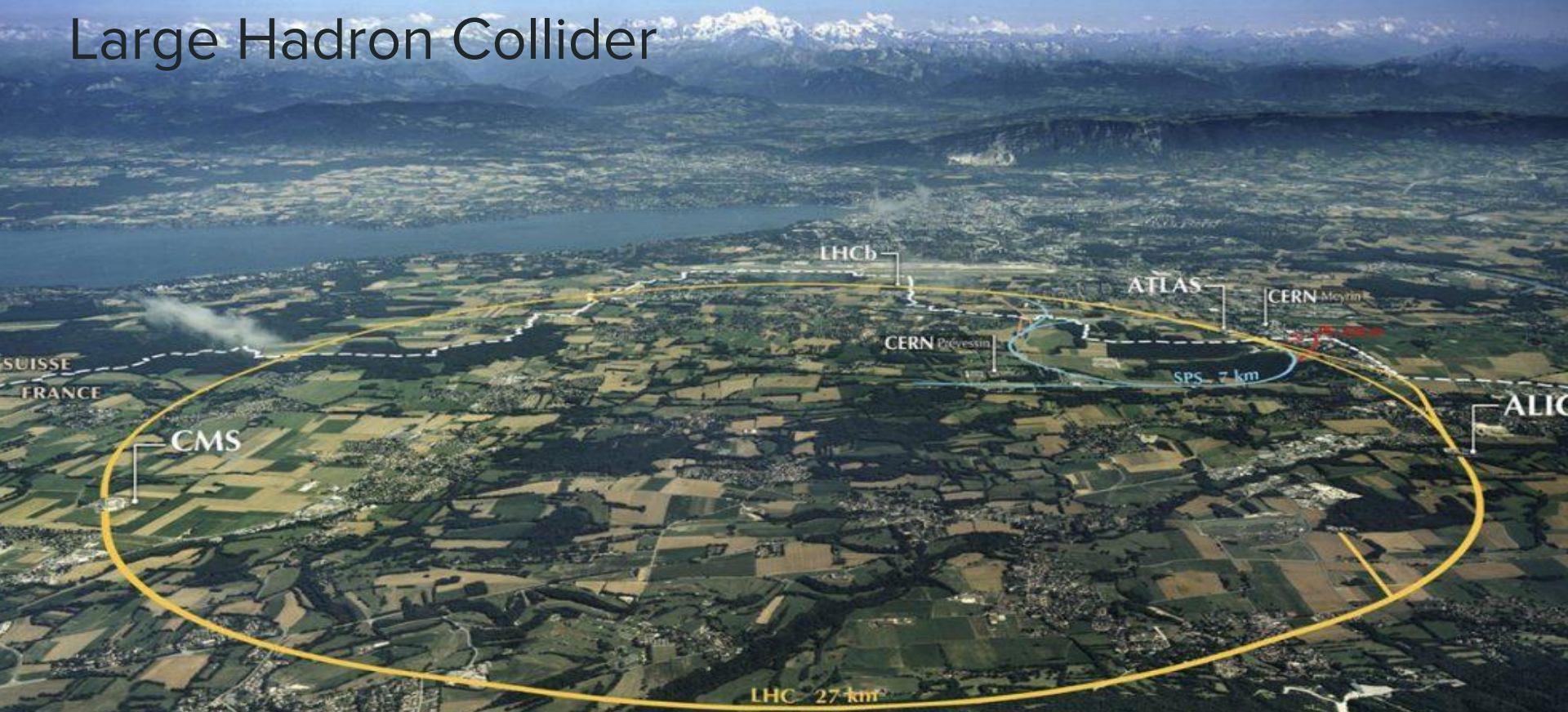
Dope a silicon detector with acceptors or donors to control the depth w of your depletion zone.

$$w = \sqrt{2\varepsilon\mu_n\rho_n V_b}$$

Resistivity is inversely proportional to the doping concentration

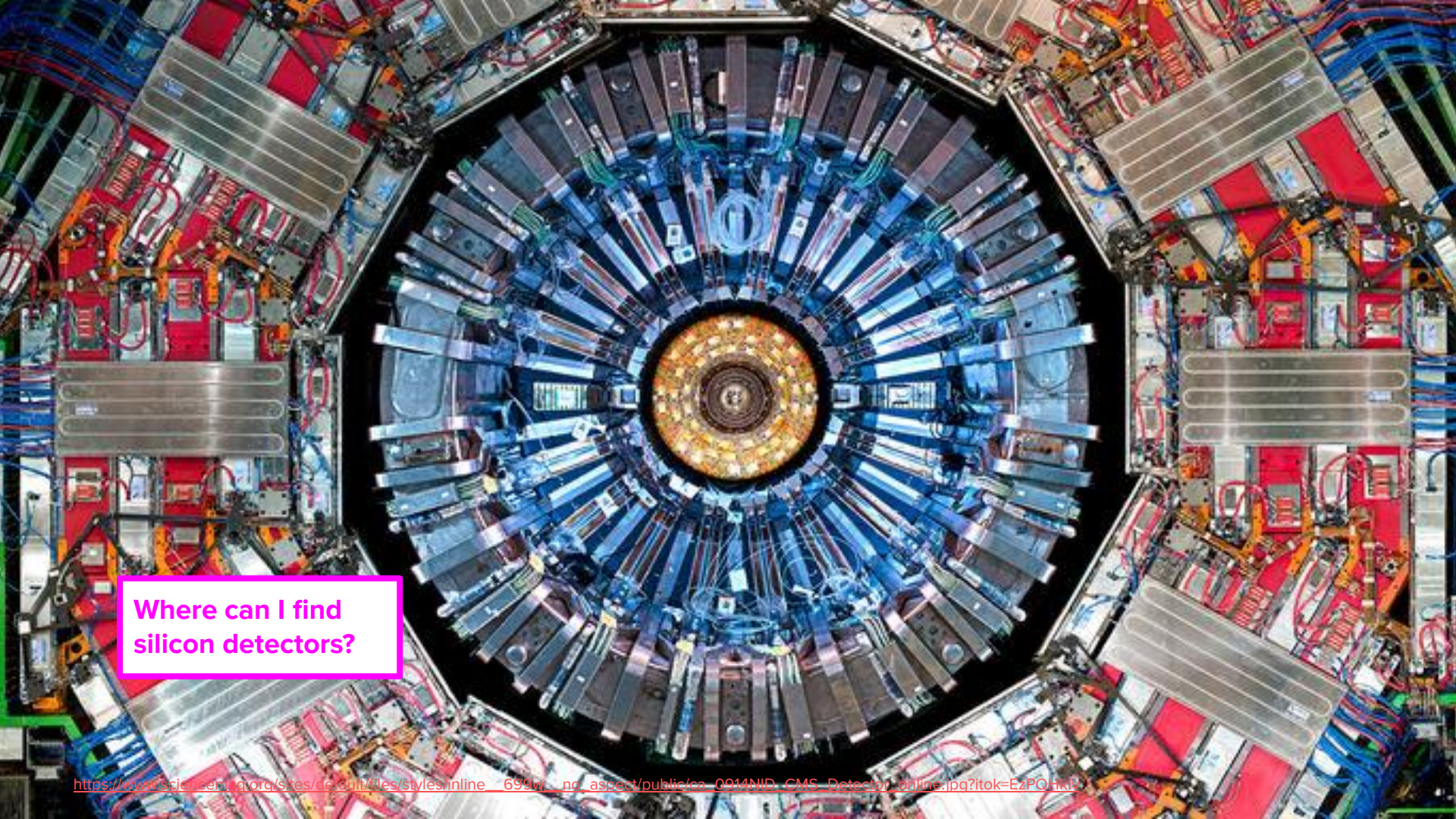
Semiconductors make excellent sensors: applications

Large Hadron Collider



<https://www.youtube.com/watch?v=NhXMXIXOWAA>

https://cdn.zmescience.com/wp-content/uploads/2015/05/cern_lhc-aerial.jpg
<http://sites.uci.edu/energyobserver/files/2012/10/lhc-aerial.jpg>

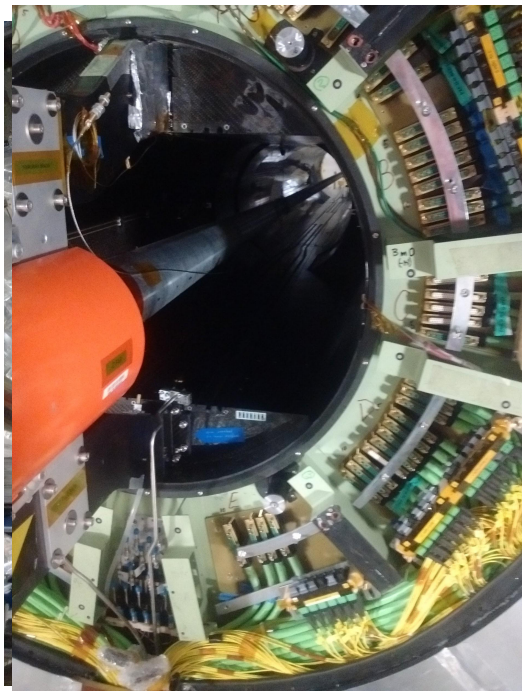
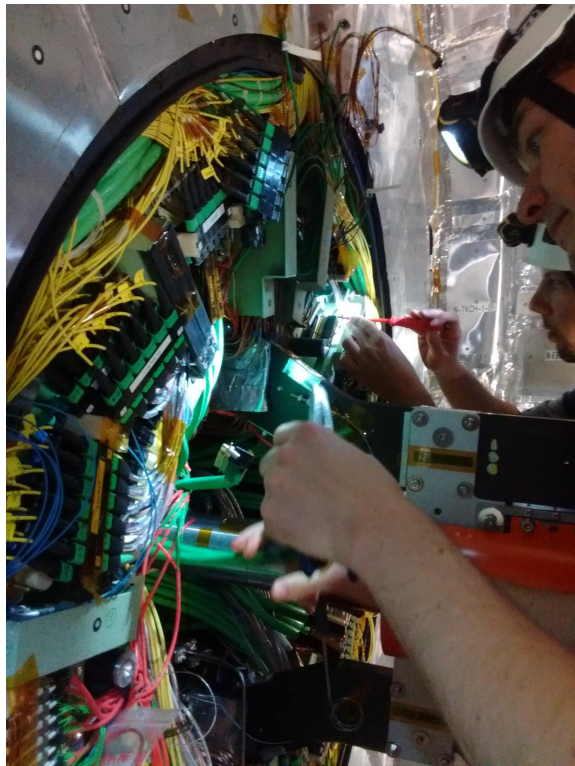


Where can I find
silicon detectors?



CMS pixel detector

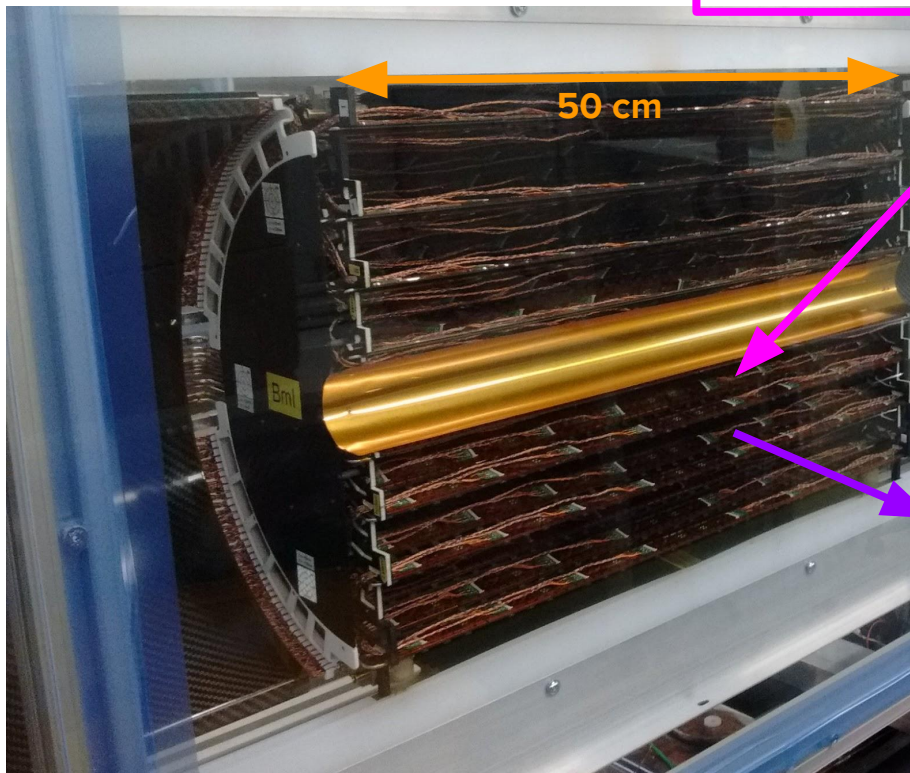
CMS pixel detector removal 2019



How about that ATLAS? ;-)

CMS pixel detector

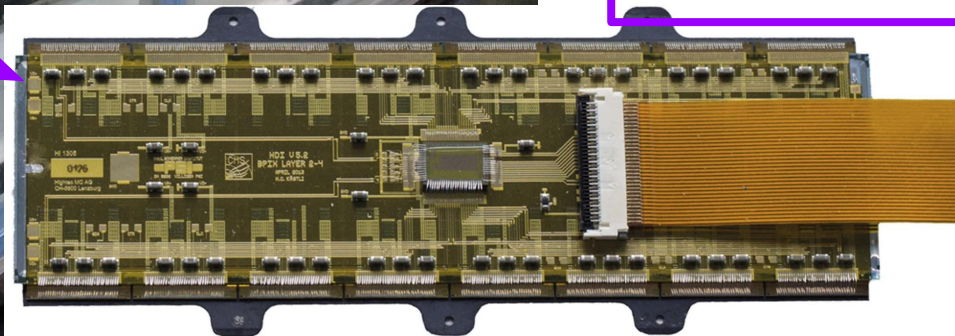
Innermost layer is currently being replaced at CERN



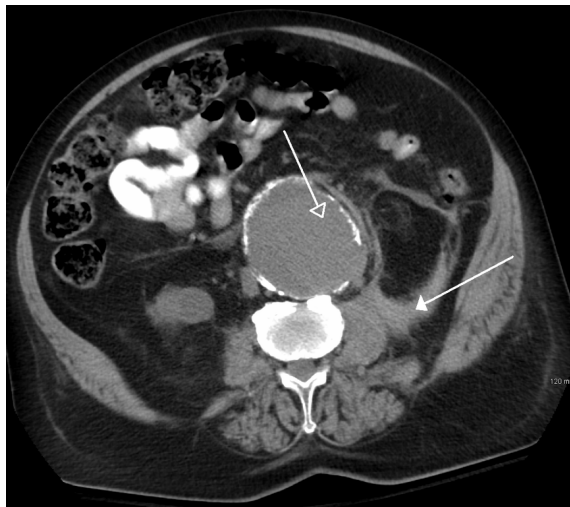
Services: DCDC converters, conversion to optical signal...



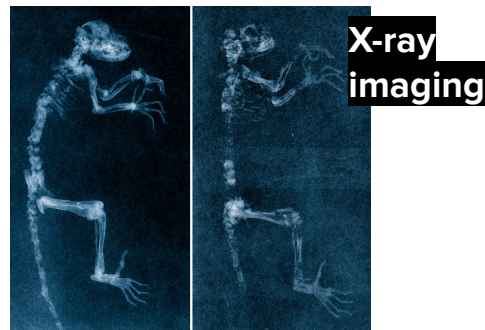
One "pixel detector"



Applications of silicon detectors

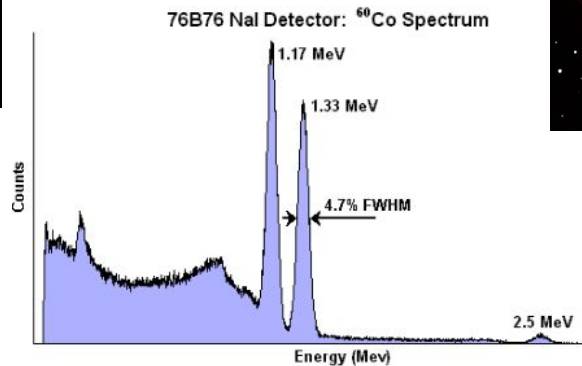


Medical imaging

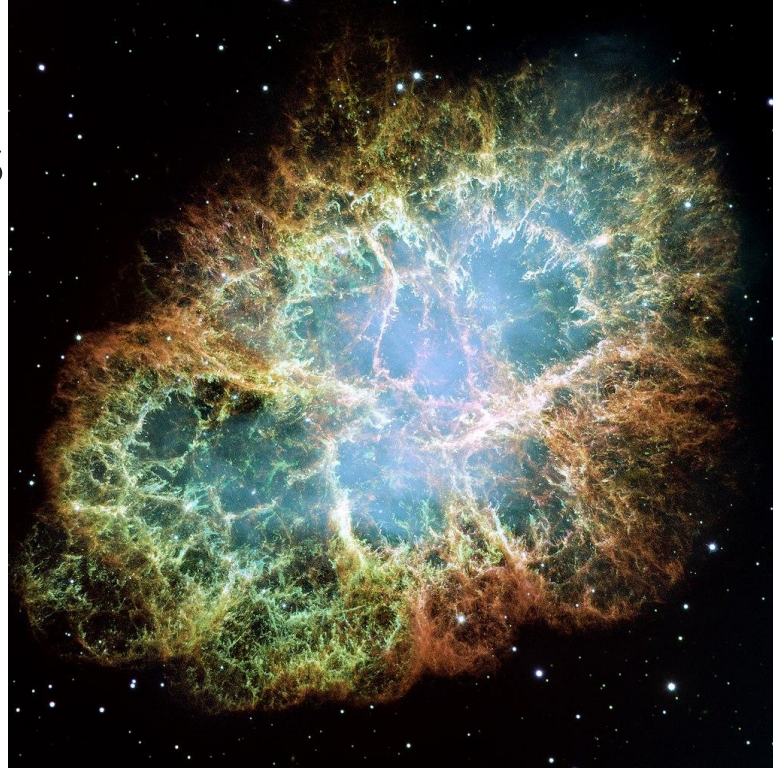


X-ray imaging

Plate A 10 cm Plate B

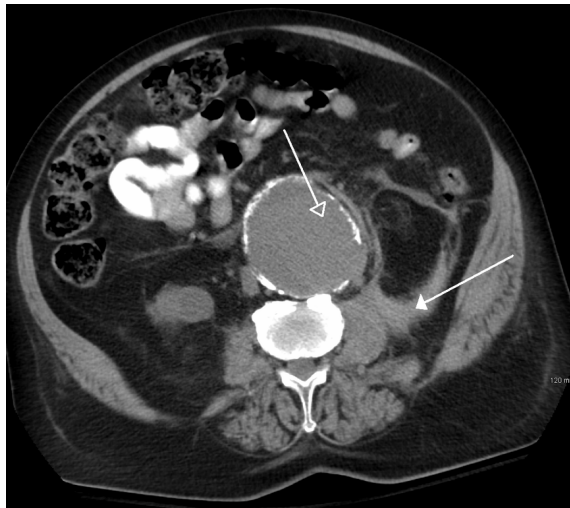


Week 1 X-ray spectroscopy



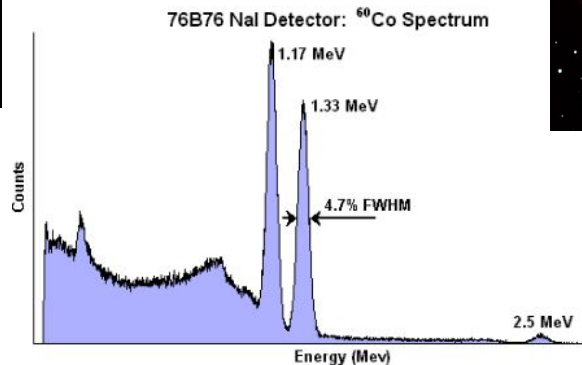
X-ray astronomy

Applications of silicon detectors

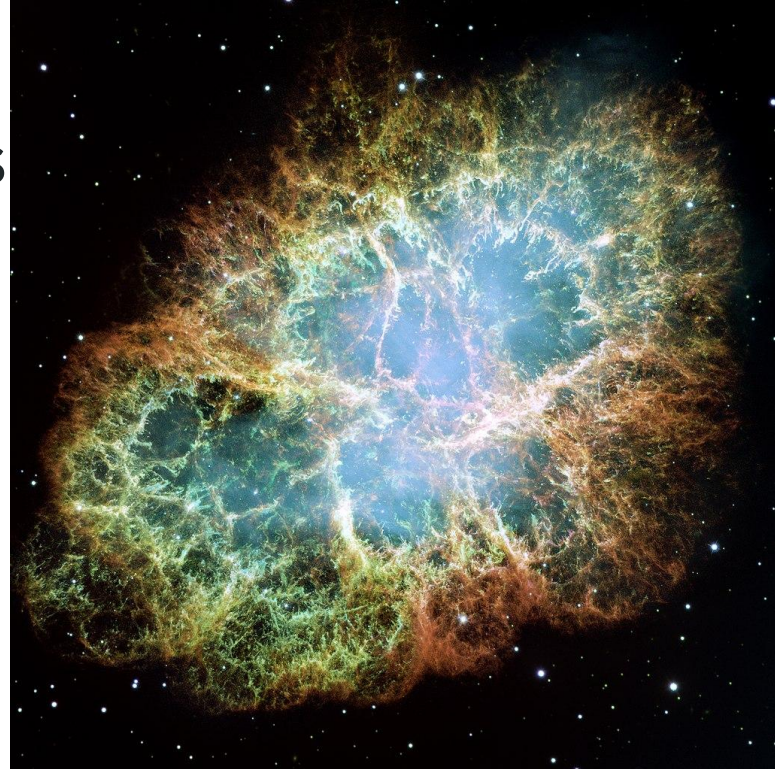


Medical imaging

Need no superfast taking of pictures every second: rather integrate over time to measure energy.



Week 1 X-ray spectroscopy

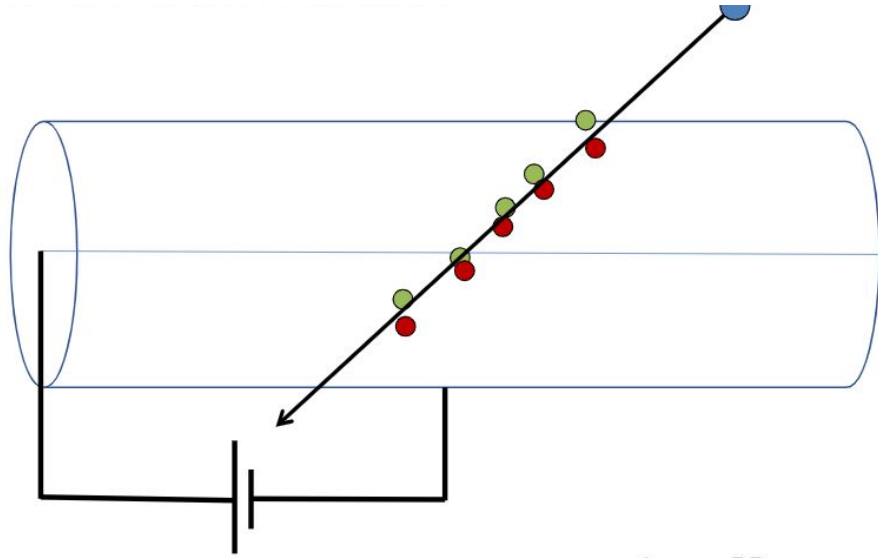


X-ray astronomy

How do we detect particles with a semiconductor?

A signal in a gas detector

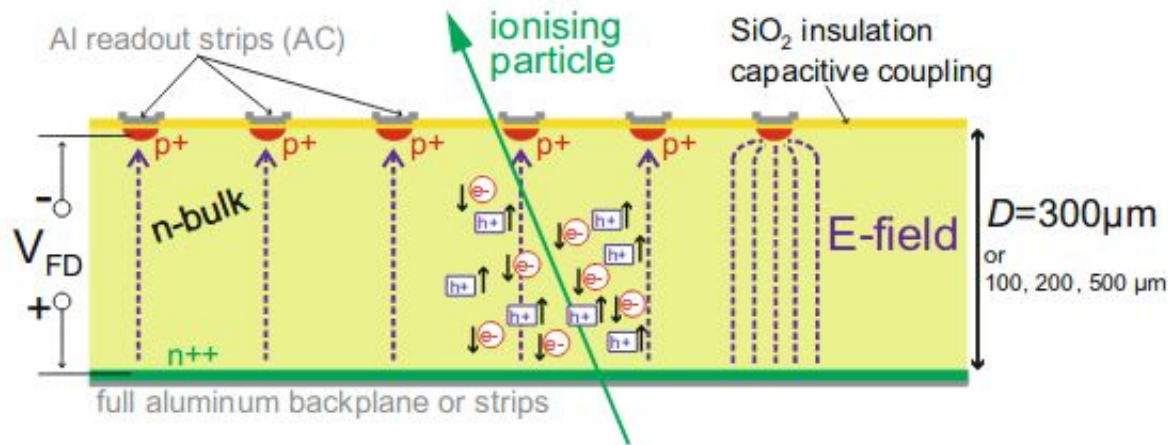
Last week:



$$E = \frac{1}{s} \frac{V}{\ln(b/a)}$$

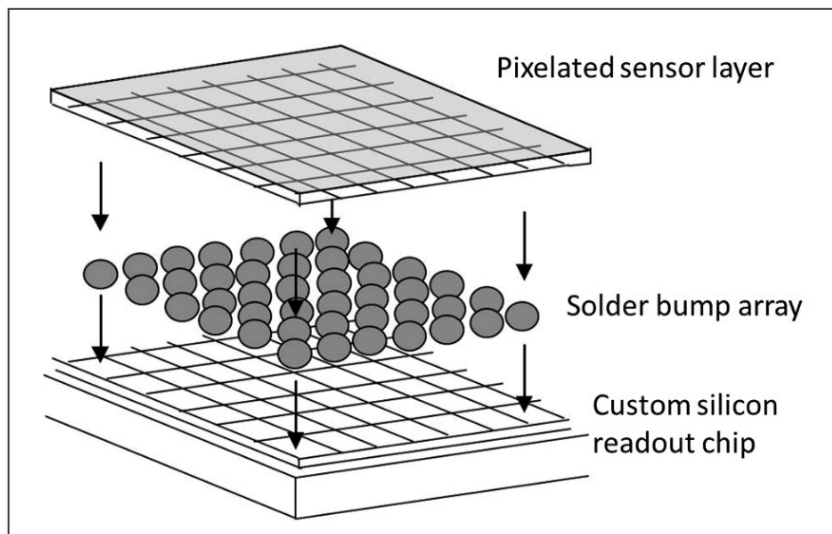
Ideal signal detection with silicon sensors

- A minimum ionizing particle (MIP) traveling through a fully depleted region (V_{FD}) creates electron hole pairs
- The charges drift to opposite directions under the electric field
- Within nanoseconds, charges are collected at the readout

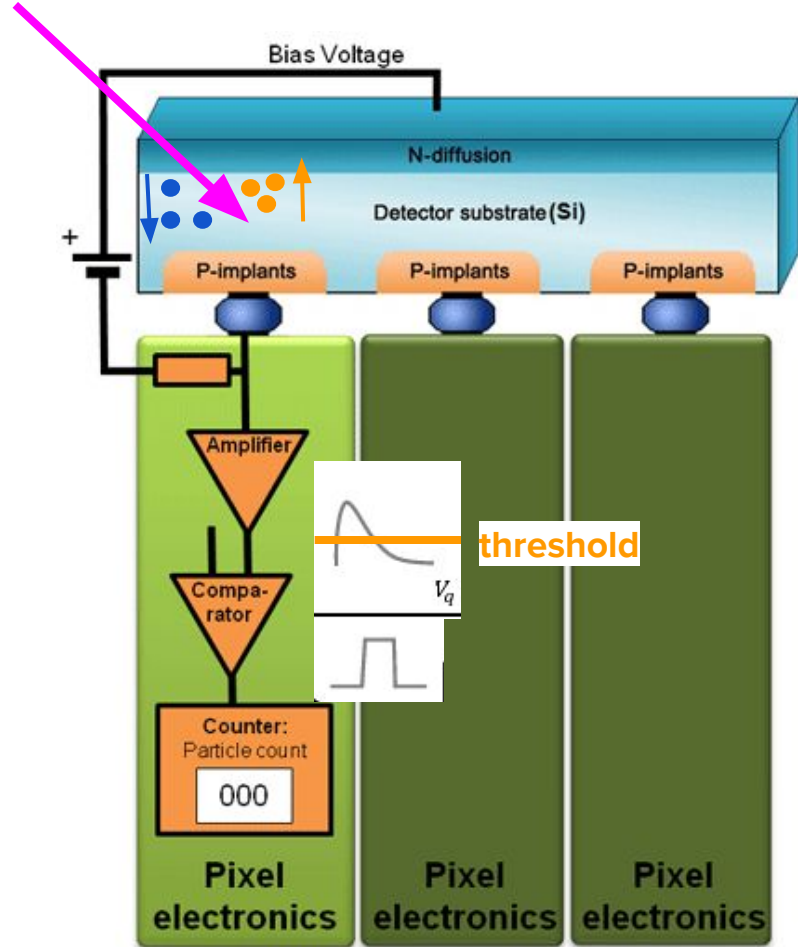


p-in-n silicon sensor

A hybrid pixel detector

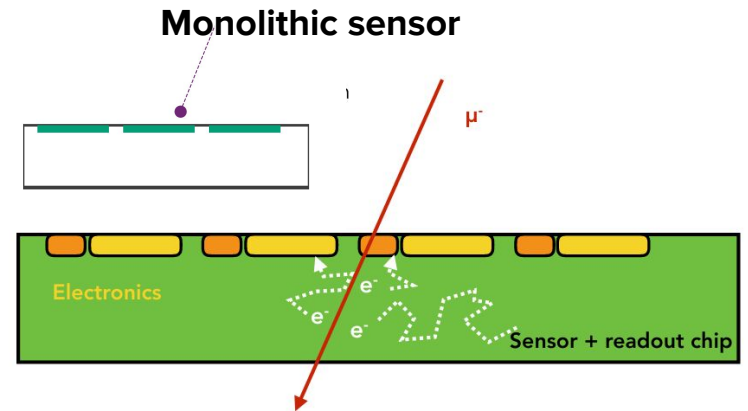
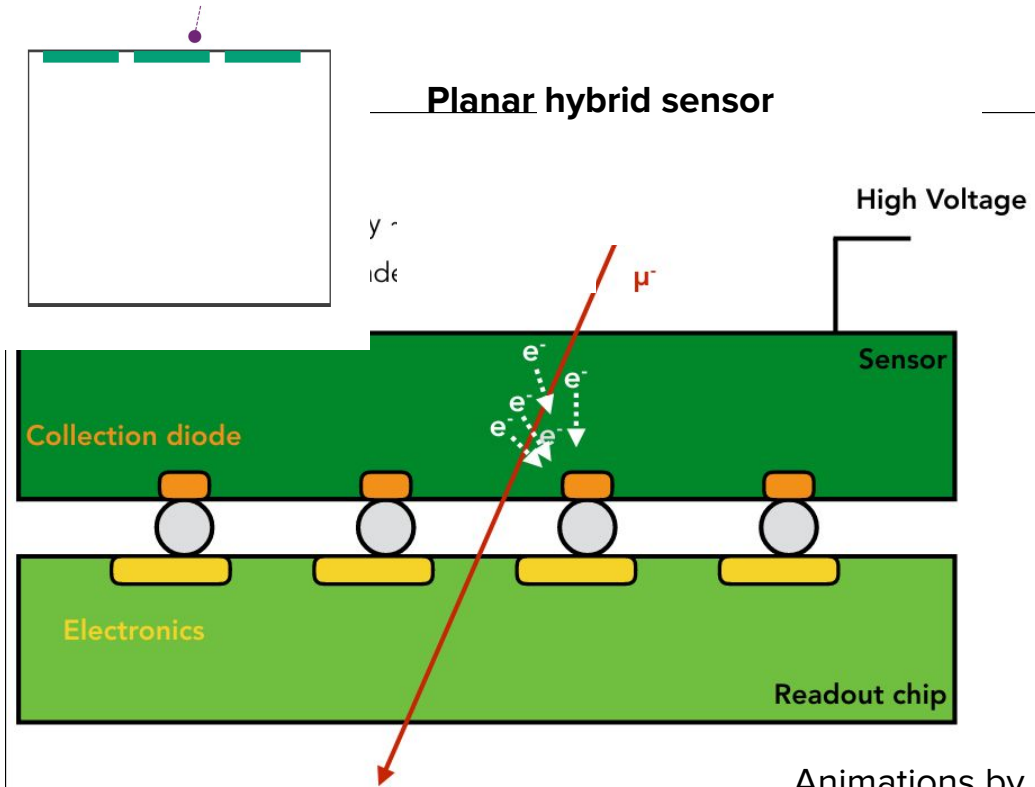


diffusion?



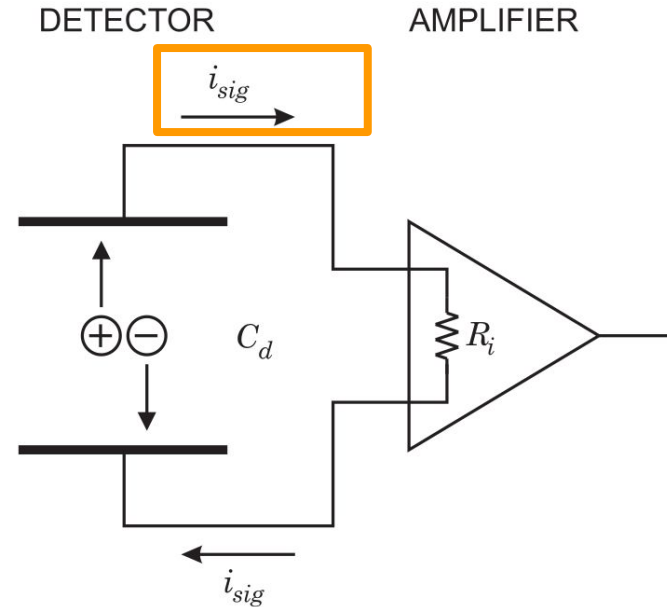
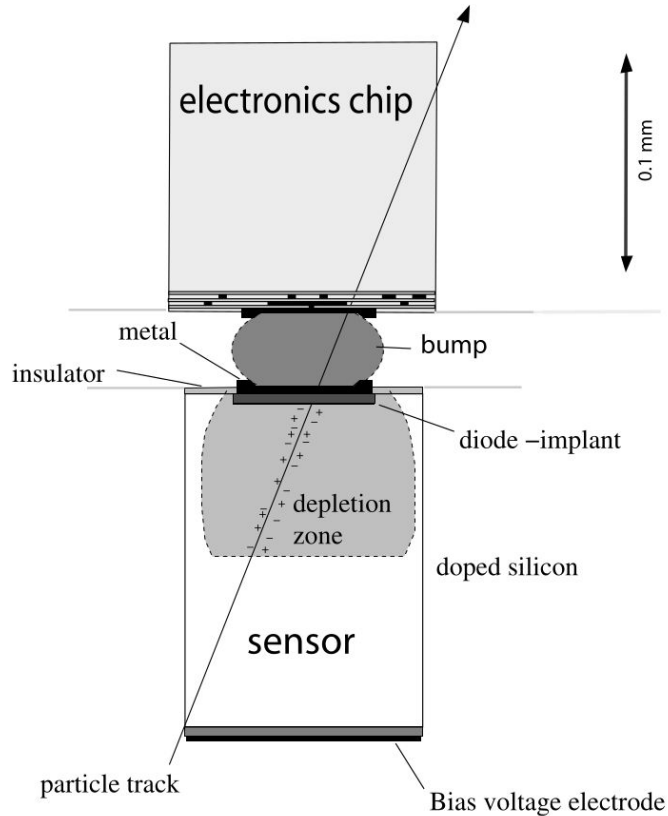
How did signal travel in gas detectors?

Drift vs diffusion



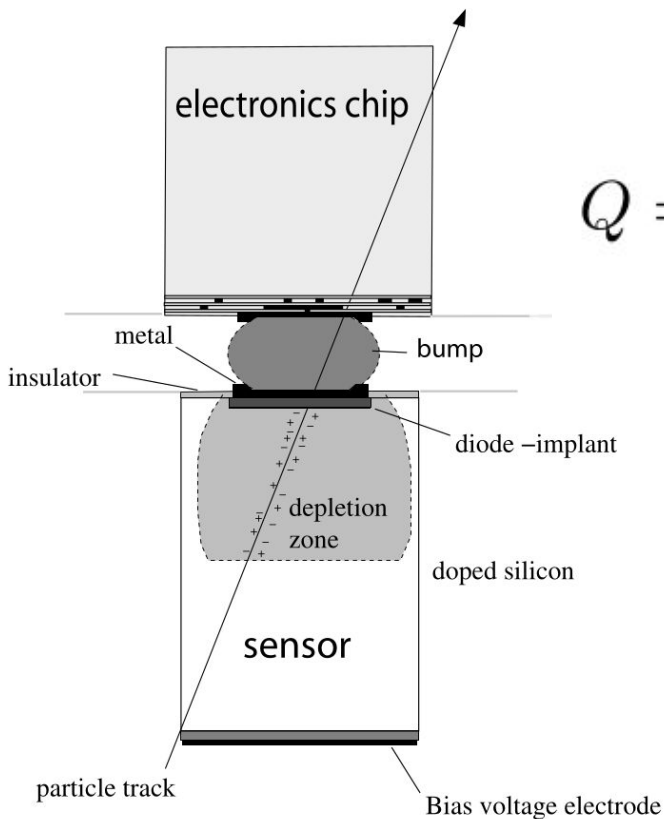
Charge “collection”

Charge is **induced** in the electrodes



Charge is **induced** in the electrodes

\mathbf{E}_w weighting field
 \neq electric field



$$Q = \int_{t_1}^{t_2} i(t) dt = e [\phi_w(\mathbf{x}_1) - \phi_w(\mathbf{x}_2)]$$

$$i = e \mathbf{E}_w \mathbf{v}$$

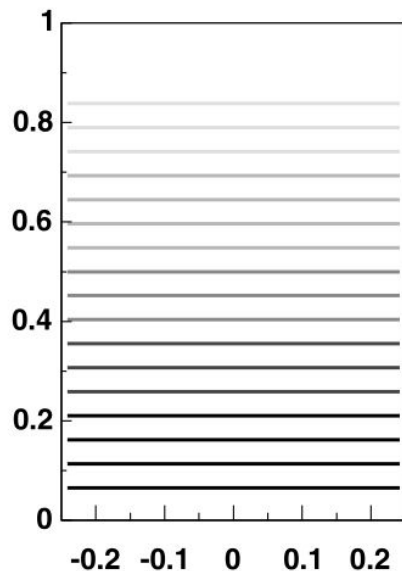
Weighting potential ϕ_w : set one electrode to unit potential, all others zero \rightarrow solve Poisson

From Ramo and Shockley

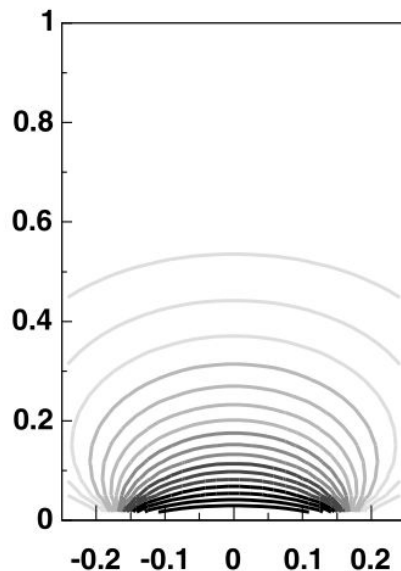
$$\frac{d^2 V}{dx^2} = - \frac{\rho(x)}{\epsilon}$$

Charge is **induced** in the electrodes

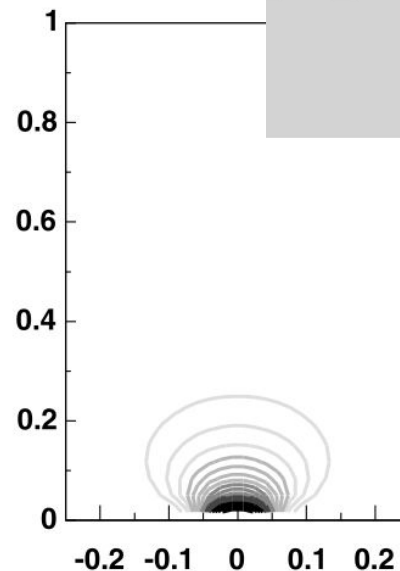
E_w is **not** the electric field! It is the **weighting potential**



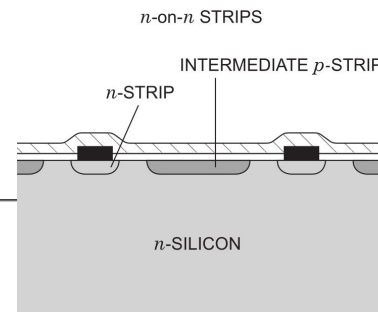
Infinite parallel plate



Collection electrode $1/3$
wafer thickness

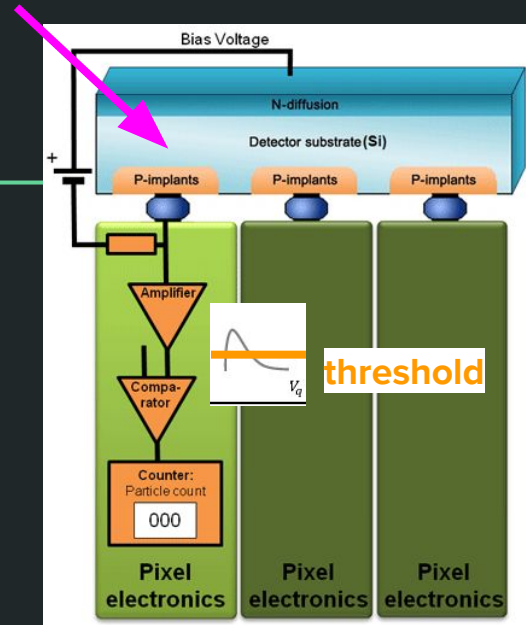


Collection electrode $1/10$
wafer thickness



How do we detect particles with a semiconductor?

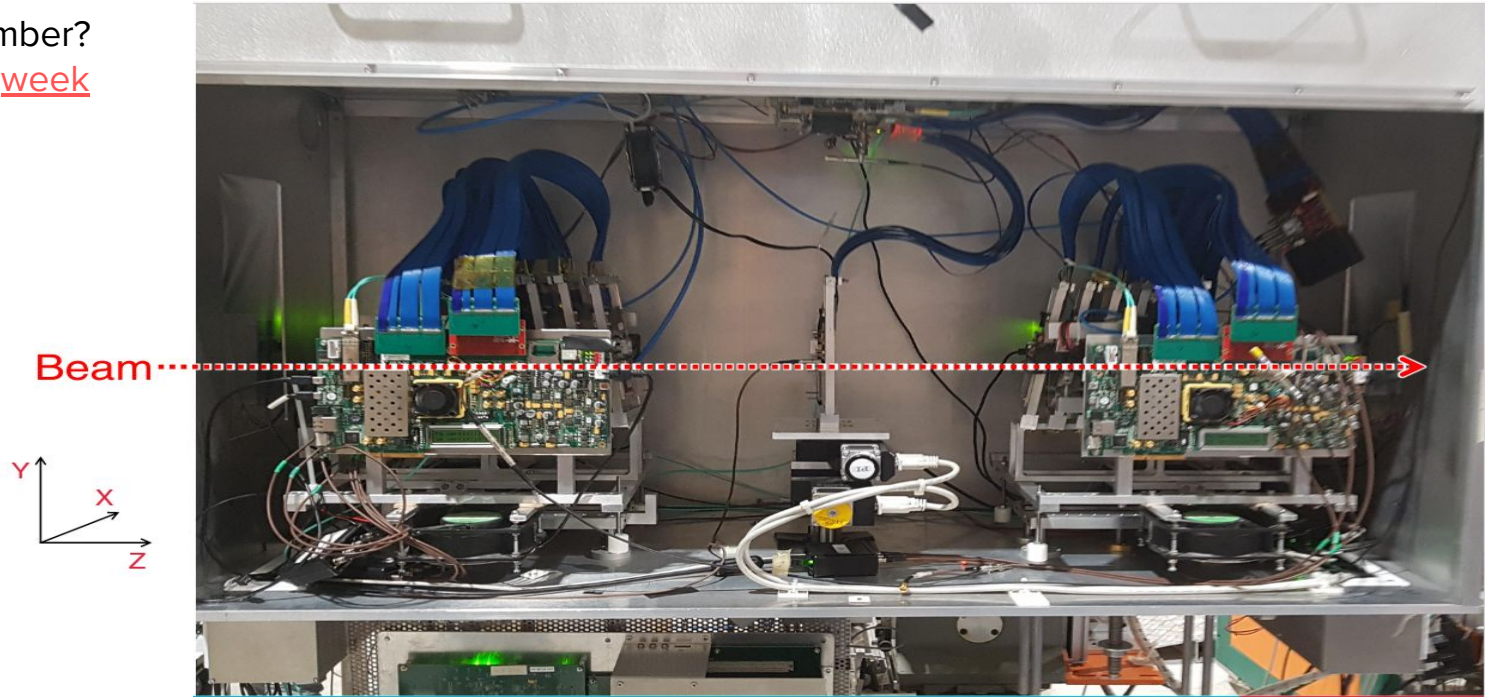
Ionizing radiation creates electron-hole pairs in a sensitive **depleted volume** of silicon. A current is **induced** in electrodes, amplified, and detected as a signal when over threshold.



How do you measure if a
semiconductor is an excellent sensor?

Sensor characterization with Timepix telescope

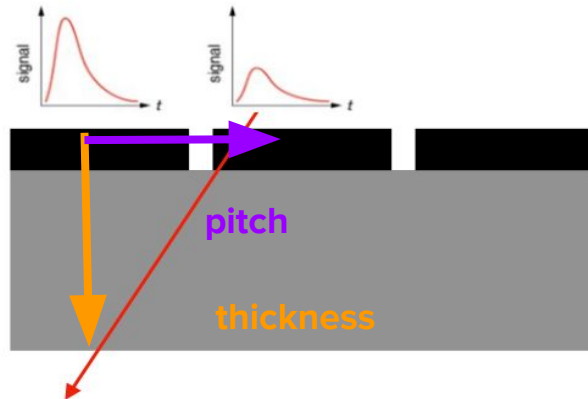
Do you remember?
Mentioned in [week](#)
[1](#)



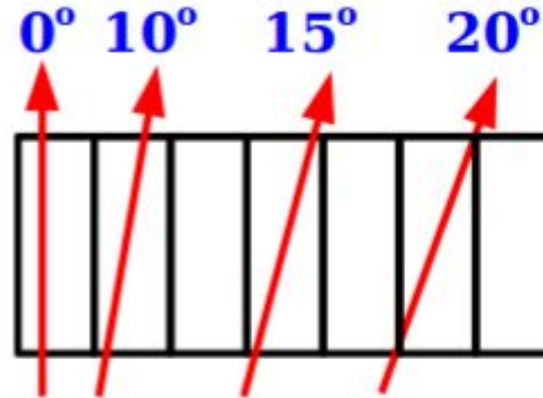
[K. Akiba et al 2019 JINST 14 P05026](#)
[K. Heijhoff et al 2020 JINST 15 P09035](#)

Pixel cell dimensions and spatial resolution

- Charge sharing depends on pixel **pitch**
- **Best resolution: cluster size of 2 . why?**



Position = weighted average
Resolution < pitch / $\sqrt{12}$

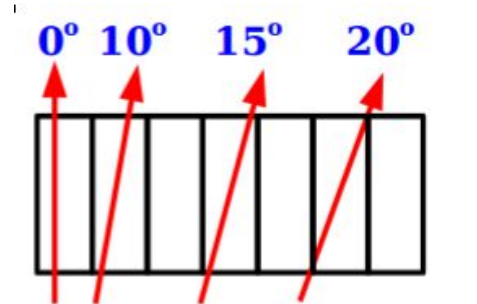
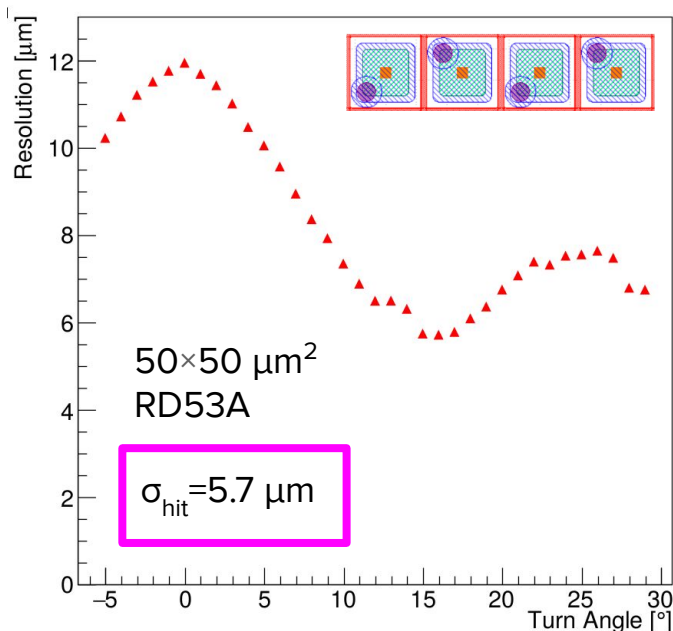


Optimal charge sharing at an angle $\text{atan}(\text{pitch}/\text{depth})$:
 $50 \times 50 \mu\text{m}^2$: $\text{atan}(50/150) = 18.4^\circ$

What happens if too many pixels share the charge?

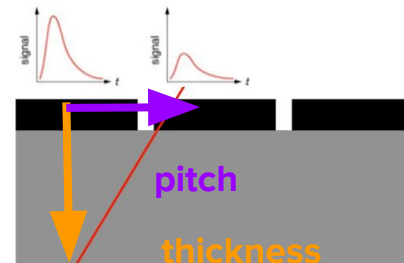
Pixel cell dimensions and spatial resolution

- Charge sharing depends on pixel **pitch**
- **Best resolution: cluster size of 2**



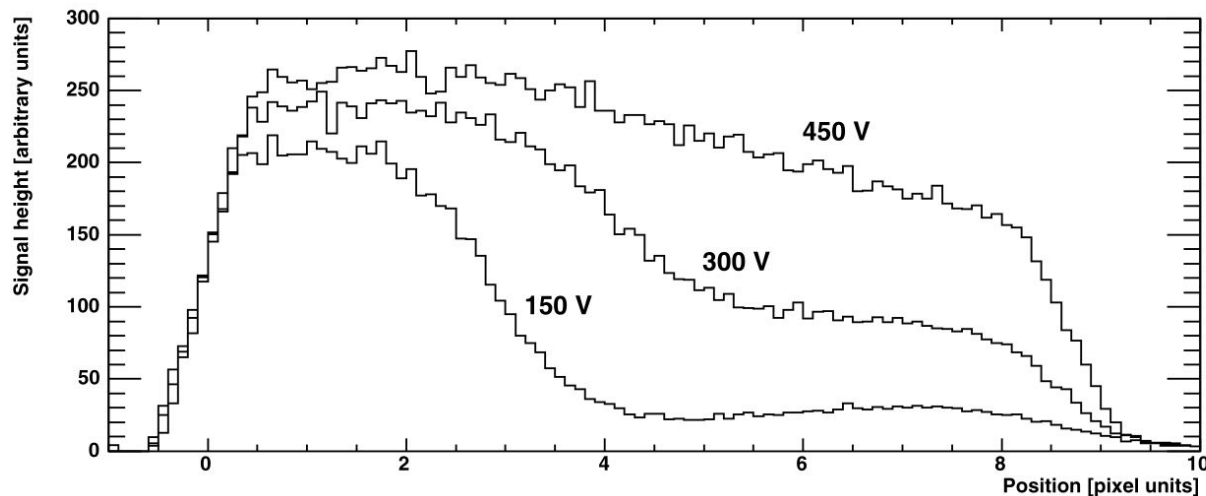
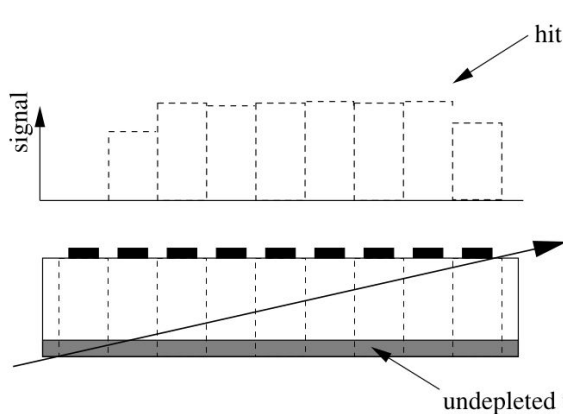
2 pixel charge sharing for 18°

Optimal charge sharing at an angle $\text{atan}(\text{pitch}/\text{depth})$:
 $50 \times 50 \mu\text{m}^2$: $\text{atan}(50/150) = 18.4^\circ$

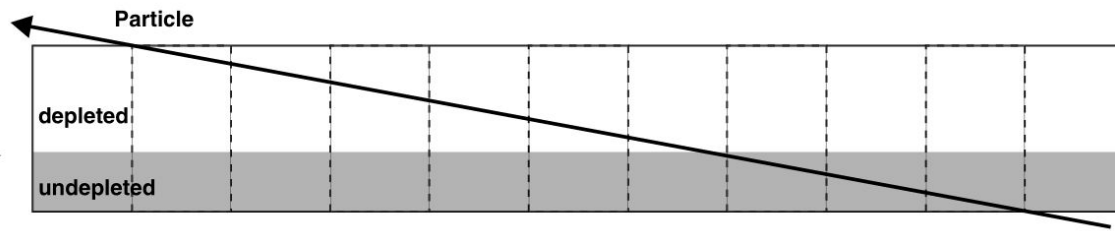
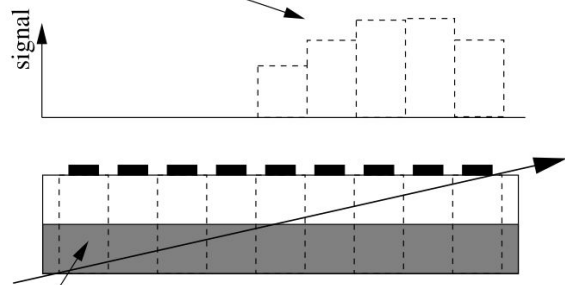


Position = weighted average
Resolution < pitch / $\sqrt{12}$

Measuring depletion depth



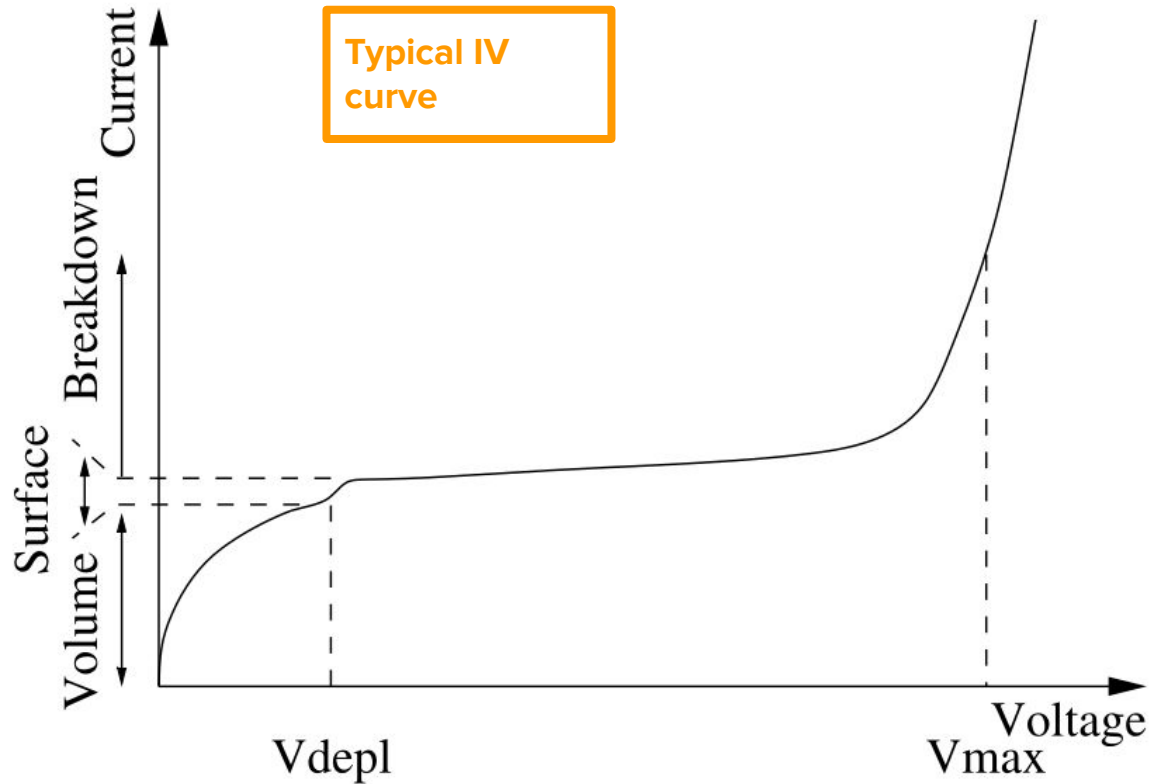
hit pixel cells



Grazing angle method

Undepleted volume

Measuring the depletion depth



How to detect particles with your
phone?

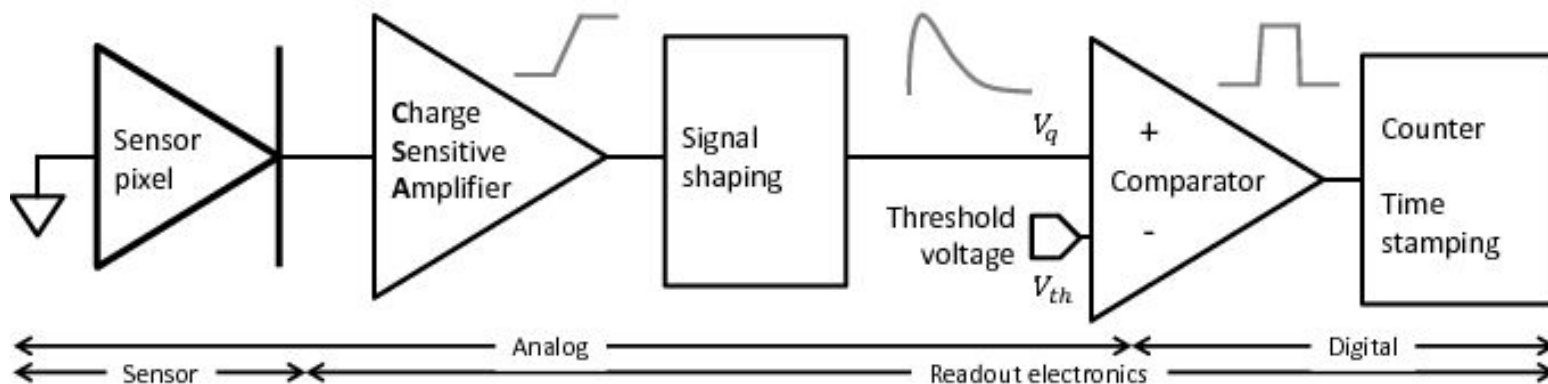
Are semiconductors easy to control?

Noise: capacitance

$$C = \epsilon_0 \epsilon_{\text{Si}} \frac{A}{d}$$

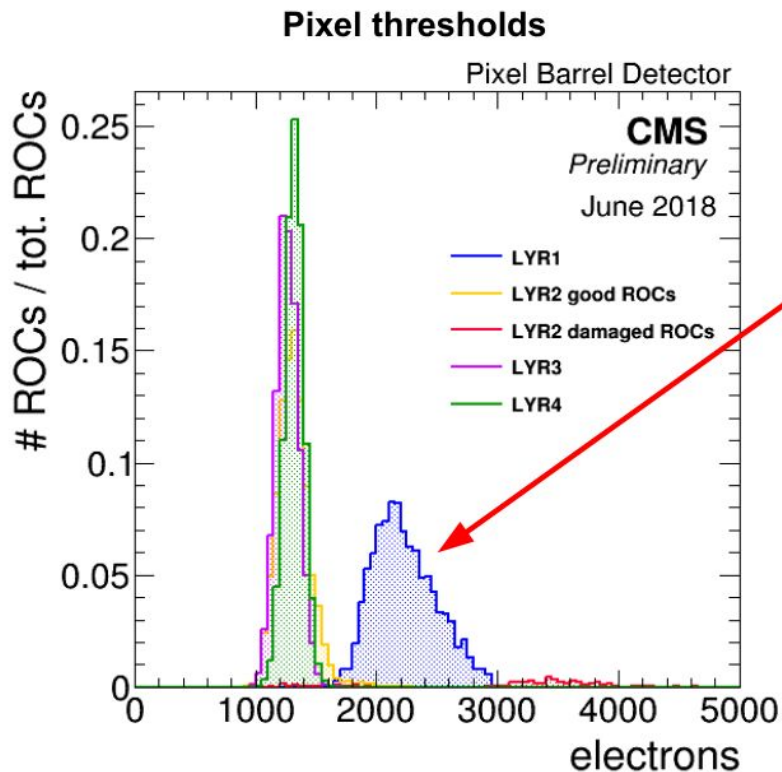
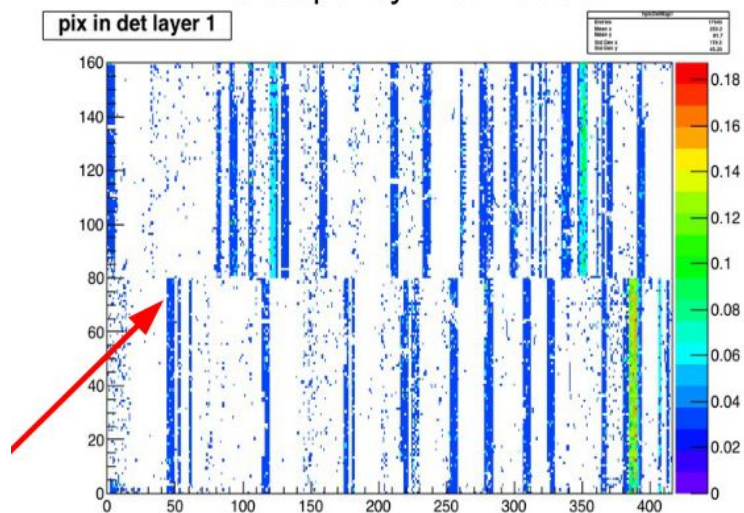
For a 20,000 μm^2 pixel cell on 300- μm -thick silicon this gives a contribution of about 7 fF

Total capacitance of each pixel to backside determines noise on preamplifier



Noise: cross talk

Cross-talk or interpixel capacitance between pixels -- hits where one does not expect hits -- here entire columns!



Thresholds are tuned to pick up signal and not too much noise. **Why are layer 1 thresholds higher?** 70

Noise: leakage current

Leakage current or dark current: thermal excitations **or** diffusion of free charge carriers from undepleted into depleted region

→ generation of electron-hole pairs

This is highly temperature dependent:

$$J_{\text{vol}} \propto T^2 e^{-E_g(T)/2kT} .$$

Roughly doubles every
8 K.

Semiconductor = silicon?

Actually, no!

Until 1959

Germanium most common
Silicon worse conductivity
because of unstable
states at surface

Other materials: diamond (C) 5.5 eV,
Gallium arsenide (GaAs) 1.43 eV, Silicon
dioxide SiO_2 9 eV ... many more!

Table 10.2. Average energy for electron-hole creation in silicon and germanium

	Si	Ge
300 K	3.62 eV	—
77 K	3.81 eV	2.96 eV

Lower band gap :
What does that
mean?

Actually, no!

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Other materials: diamond (C) 5.5 eV,
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Table 10.2. Average energy for electron-hole creation in silicon and germanium

	Si	Ge
300 K	3.62 eV	—
77 K	3.81 eV	2.96 eV

Lower band gap :
→ larger signal
→ larger leakage
current

Table 2.3 Representative detector materials. Mobilities μ are in units of $\text{cm}^2\text{V}^{-1}\text{s}^{-1}$ and $\mu\tau$ products in cm^2V^{-1} .

Material	E_g (eV)	E_i (eV)	ε	μ_e	μ_h	$(\mu\tau)_e$	$(\mu\tau)_h$	ρ	$\langle Z \rangle$
Si	1.12	3.6	11.7	1350	450	> 1	> 1	2.33	14
Ge	0.67	2.96	16	3900	1900	> 1	> 1	5.33	32
GaAs	1.43	4.2	12.8	8000	400	$8 \cdot 10^{-5}$	$4 \cdot 10^{-6}$	5.32	31.5
Diamond	5.5	13	5.7	1800	1200			3.52	6
4H-SiC	3.26	8	9.7	1000	115	$4 \cdot 10^{-4}$	$8 \cdot 10^{-5}$	3.21	10
GaN	3.39	8 – 10		1000	30			6.15	19
InP	1.35	4.2	12.4	4600	150	$5 \cdot 10^{-6}$	$< 10^{-5}$	4.78	32
CdTe	1.44	4.43	10.9	1100	100	$3 \cdot 10^{-3}$	$2 \cdot 10^{-4}$	5.85	50
Cd _{0.9} Zn _{0.1} Te	1.572	4.64	10	1000	120	$4 \cdot 10^{-3}$	$1.2 \cdot 10^{-4}$	5.78	49.1
HgI ₂	2.15	4.2	8.8	100	4	$3 \cdot 10^{-4}$	$4 \cdot 10^{-5}$	6.4	62
TlBr	2.68	6.5	30	30	4	$5 \cdot 10^{-4}$	$2 \cdot 10^{-6}$	7.56	58
a-Si	1.9	6	12	1 – 4	0.05	$2 \cdot 10^{-7}$	$3 \cdot 10^{-8}$	2.3	14

Actually, no!

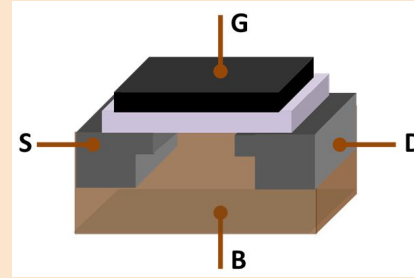
Until 1959

Germanium most common
Silicon worse conductivity
because of unstable
states at surface



1959 Mohamed Atalla and Dawon Kahng

develop surface passivation: coat silicon surface
with SiO_2 :



Metal-oxide-semiconductor field-effect transistor

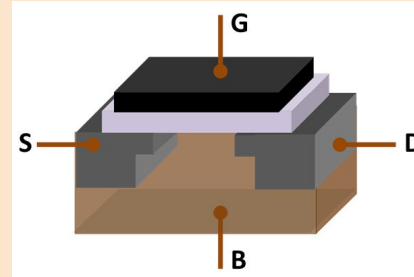
Silicon has a high quality, robust native
oxide. This property is unique among
semiconductors.

Actually, no!

Until 1959

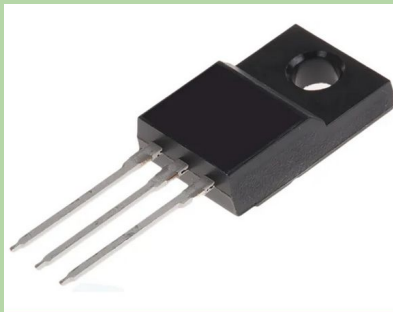
Germanium most common
Electric signal could not reach electrode in silicon: worse conductivity because of unstable surface states

1959 Mohamed Atalla and Dawon Kahng
develop surface passivation: coat silicon surface with SO_2 :



Metal-oxide-semiconductor field-effect transistor

1960-2018 $13e22$ MOSFETs produced: most frequently manufactured device in history!



Op voorraad - levertijd 3 à 5 werkdagen

Prijs Each (In a Pack of 5)

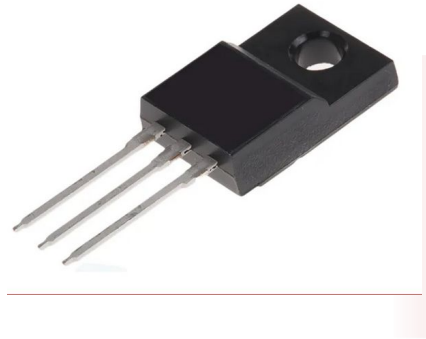
€ 0,78
(excl. BTW)

Aantal stuks	Per stuk
5 - 20	€ 0,78
25 - 95	€ 0,686
100 - 245	€ 0,60
250 - 495	€ 0,57
500 +	€ 0,556

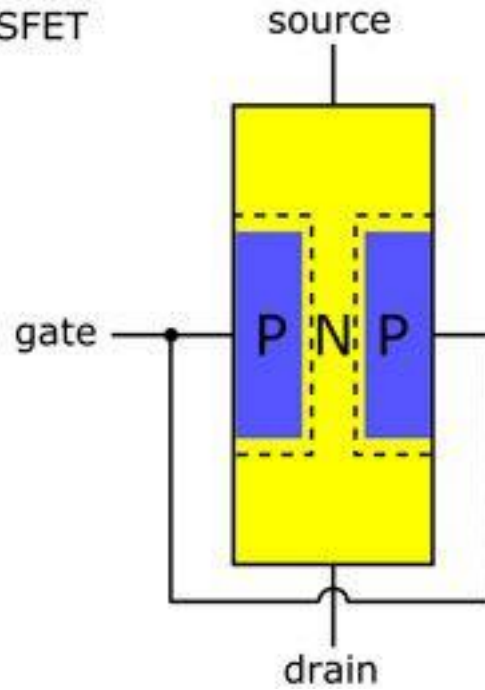
*prijsindicatie

What are the three pins?

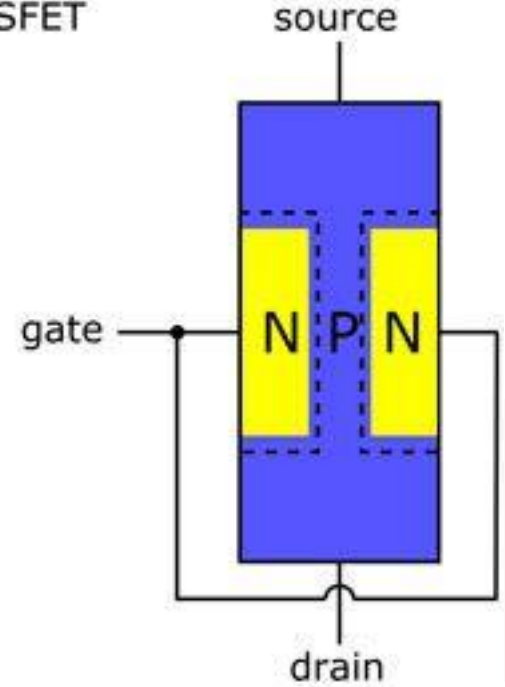
MOSFET



N-channel
MOSFET



P-channel
MOSFET



Source

Drain

Gate



n

n

P

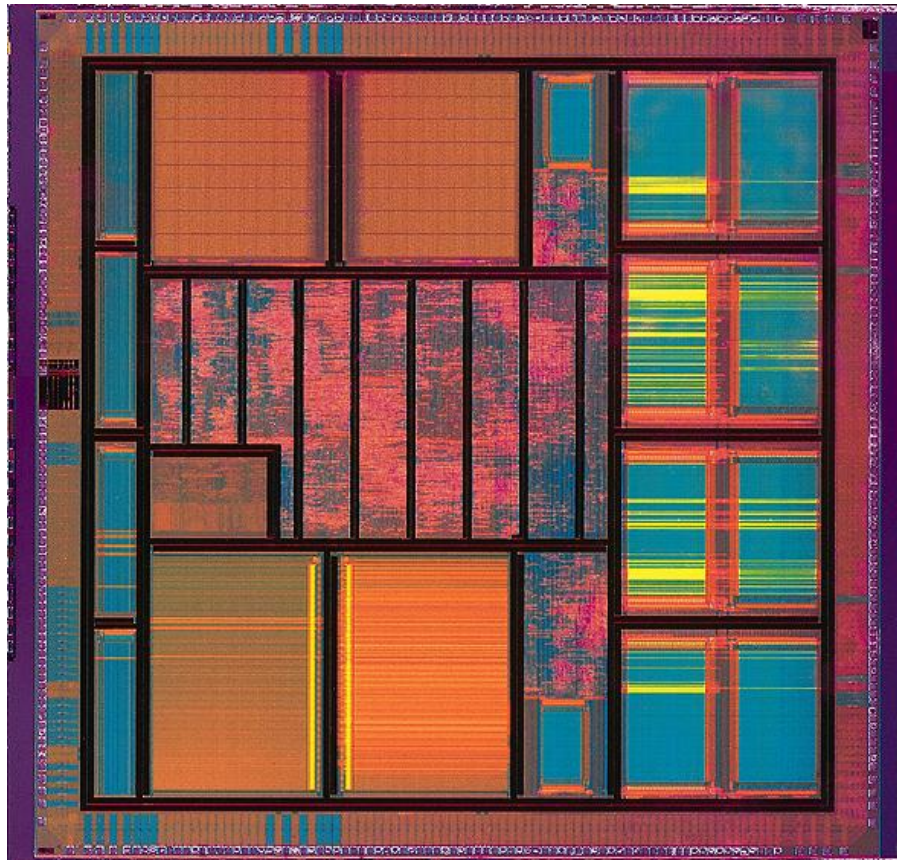


Those tiny things

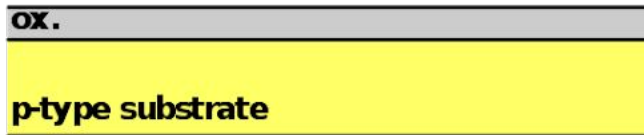
Very large scale
integration: VLSI

Combines many transistors
on a single chip to create
an integrated circuit

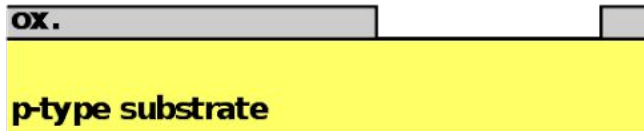
Feature size: 22 nm in the
movie. We are only at 180
nm today (CMS pixel still
has 250 nm!)



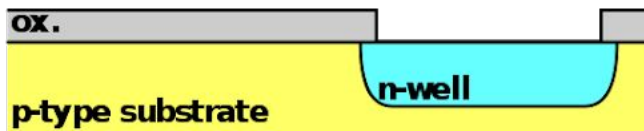
1. Grow field oxide



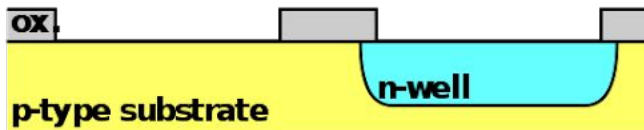
2. Etch oxide for pMOSFET



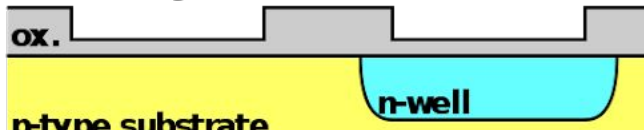
3. Diffuse n-well



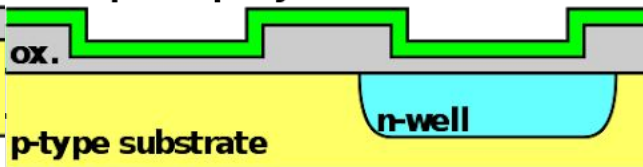
4. Etch oxide for nMOSFET



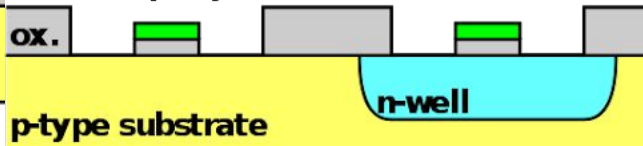
5. Grow gate oxide



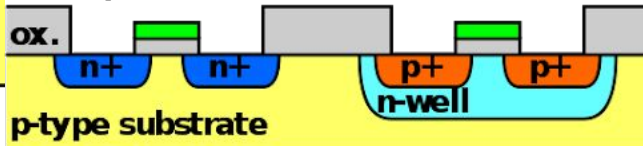
6. Deposit polysilicon



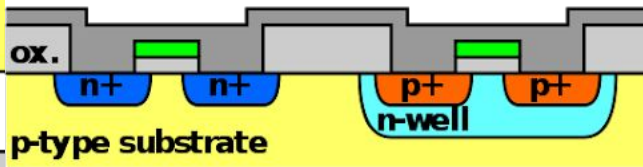
7. Etch polysilicon and oxide



8. Implant sources and drains



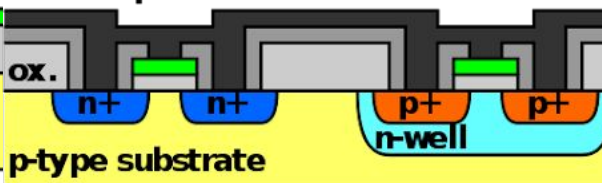
9. Grow nitride



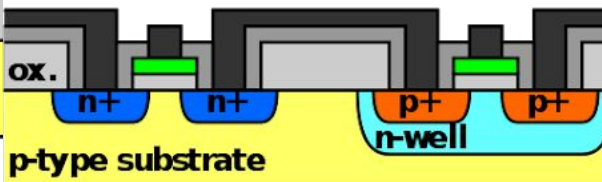
10. Etch nitride



11. Deposit metal



12. Etch metal



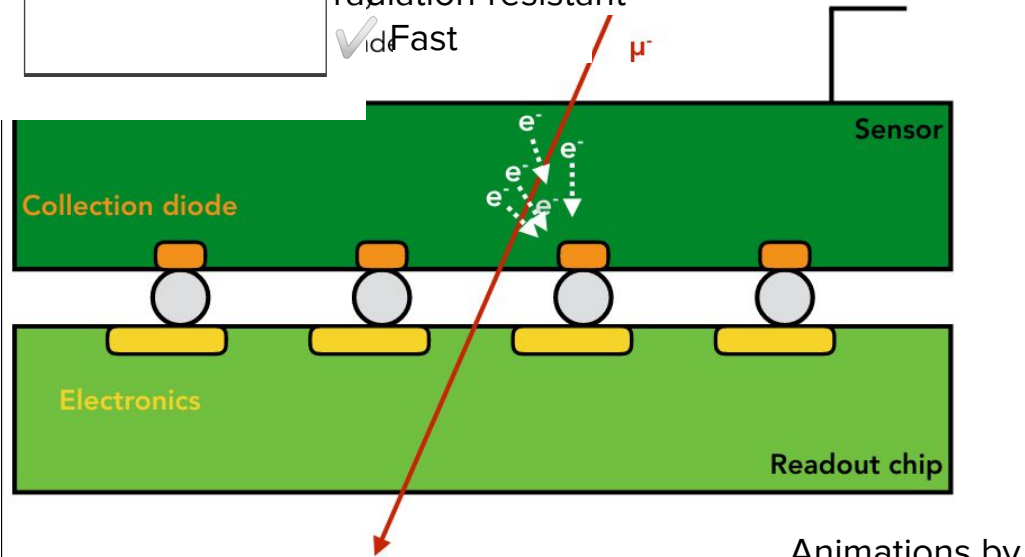
A long process to create a CMOS sensor!

A semiconductor device is easy to manipulate, easily controlled, and makes an excellent sensor

Two types of silicon sensors: hybrid vs monolithic

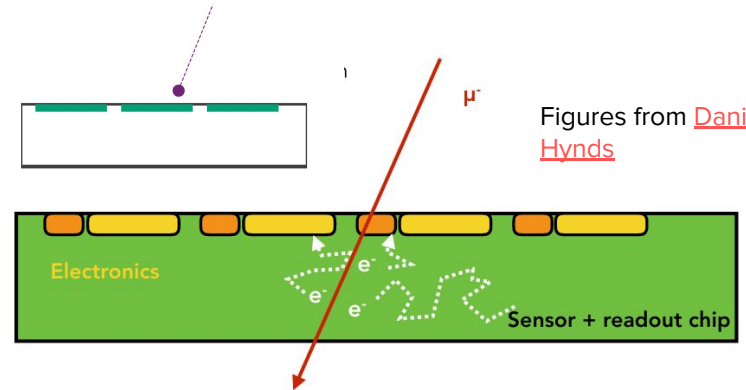
Planar hybrid sensor:

- ✗ Limited thickness for stability
- ✗ Limited pixel pitch
- ✗ Bump bonding costly
- ✓ But widely used and reliably radiation-resistant
- ✓ Fast



Monolithic sensor:

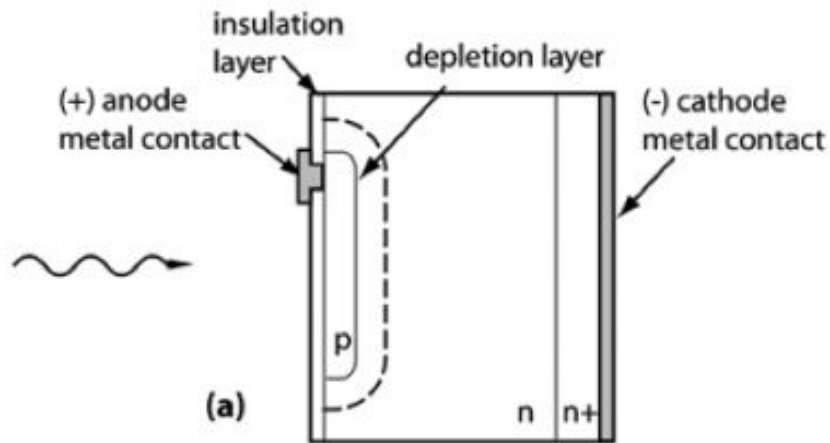
- ✓ Very little material
- ✓ Low noise
- ✗ Limited depletion region: slow charge collection by diffusion **but this is improving**
- ✗ Not very radiation-hard **but this is improving**



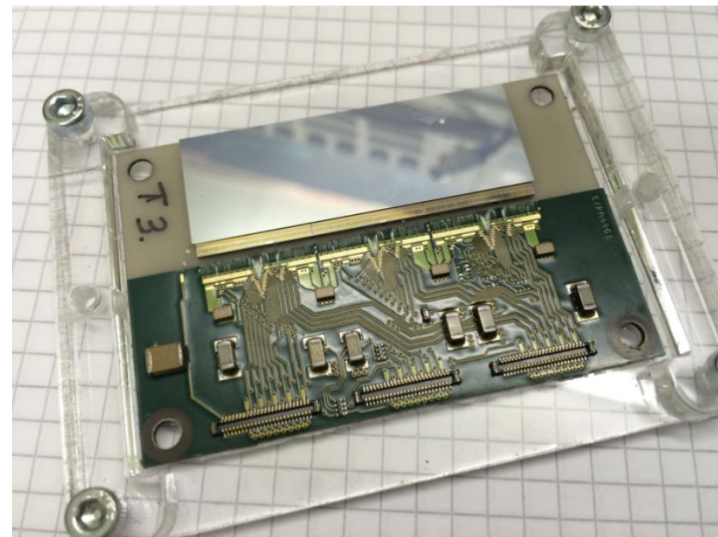
Figures from [Daniel Hynds](#)

Other semiconductor detectors

Timepix -- and Medipix -- were developed at Nikhef



Silicon photomultiplier -- [from week 2](#)



An LHCb sensor on Timepix3
Application **S**pecific **I**ntegrated **C**ircuit
(ASIC)

Why silicon not gas?

Next generation silicon detectors

Voorjaar
2019

DIM ENSIES IES

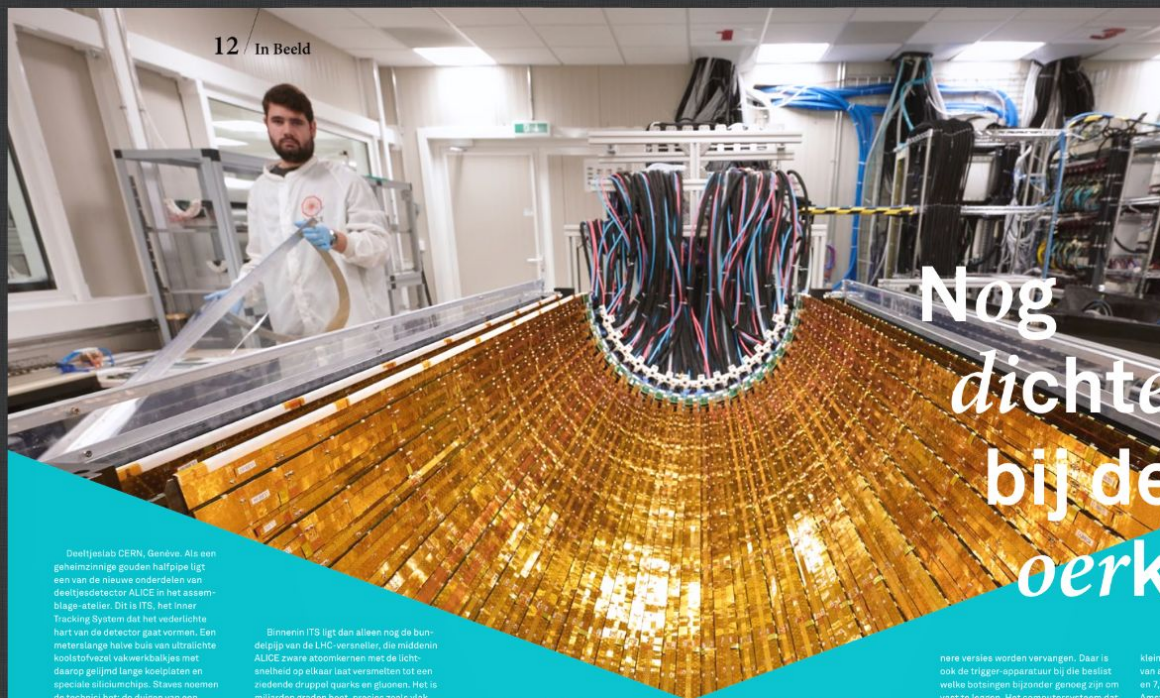
Nikhef

Nationaal
instituut voor
subatomaire fysica

UPGRADES
Van CERN tot Gran Sasso:
overal wordt nu gesleuteld
aan de experimenten

ASTRODEELTJES
Vele vensters
op het universum

12 / In Beeld



Nog
dichter
bij de
oerkr

Deeltjeslab CERN, Genève. Als een geheimzinnige gouden halfronde ligt een van de nieuwe onderdelen van de deeltjesdetector ALICE in het assemblage-atelier. Dit is ITS, het Inner Tracking System dat het vederlichte hart van de detector gaat vormen. Een meterlange halve buis van ultralichte koolstofvezel vakwerkbalkejes met daarop gelijmd lange koeplaten en speciale siliciumchips. Staves noemen de technici het de dingen van een high-tech ton.

Het goud is overigens geen goud, maar polyimide-folie met ragdunne koperen voedingskabels voor de sensoren. Dur genoeg om vrijkomende deeltjes in het experiment niet in de weg te zitten.

Het Inner Tracking System wordt een steeltel van zeven van dit soort lichtere

Binnenin ITS ligt dan alleen nog de bundel van de LHC-versnellers, die middenin ALICE zware atoomkernen met de licht-snelheid op elkaar laat vermalen tot een zeldzame druppel quark en gluonen. Het is miljoenen graden heet, precies zoals vlak na de oerknal waarin het heelal zelf moet zijn ontstaan. ITS moet de deeltjes betrappen die uit die ziedende oersop ontstappen en de fysici vertellen wat er daarbinnen precies gaande is.

Op oudere opties van botsende koolkernen is een verbijsterende warboel van rondvliegende deeltjes te zien. Veel eigen-scherven van de draupel maakte en

geven, is de verwachting. De nieuwe inner tracker doet dat van nog dichterbij dan de eerdere versie. Daardoor kunnen de kern-botsingen preciezer worden bekeken.

Honderd keer beter

ITS is te zien als een van de grootste digitale camera's ter wereld. Het oppervlak is in totaal tienwinkanta meter. En zijn ruim

dezelfde plak silicium zitten. Dat schiet kabels en elektronica in de detector.

In de komende meetperiode kan ITS gemakkelijk honderd keer zoveel meetgegevens verzamelen als alles wat ALICE in eerdere rondes binnenhaalde. Dat is cruciaal voor het bestuderen van zeldzame maar veelzeggende processen in de zeldzame oersop.

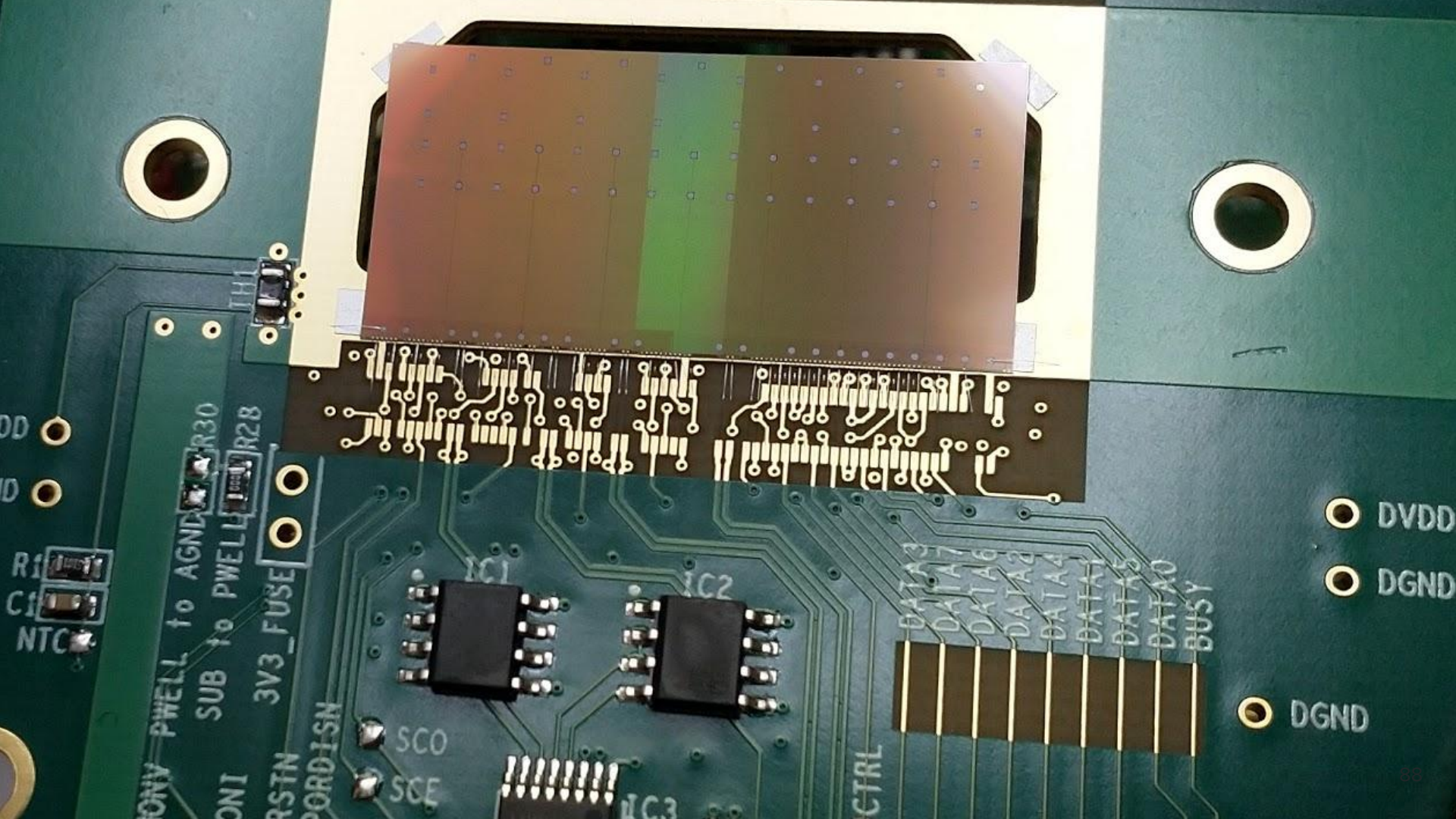
ere versies worden vervangen. Daar is ook de trigger-apparatuur bij die beaalt welke botsingen bijzonder genoeg zijn om vast te leggen. Het computersysteem dat data verzamelt en toegankelijk maakt, wordt eveneens vernieuwd.

De upgrade-periode is een hectische tijd. Het binnenste van de grote ondergrondse detector is vorig jaar meteen leeggehaald. Tot alle nieuwe onderdelen zijn ingebouwd, staan de karakteristieke rode stalen deuren wijd open in het ondergrondse laboratorium. Inmiddels

klein van en 7. Amz ALICE hang gers. al in naar aard ITS

Nikhef R&D group works with LHC experiments and Nikhef electronics group on pixel detectors that are now being installed at the LHC!

Remember from week 1 the ALICE MAPS detector? 0.1% X_0 !



DD
D
R1
C1
NTC

ONV PWELL to AGND R30
SUB to PWELL R28
3V3_FUSE
ONV
RSTN
ORDISN

SC0
SC1
SC2

IC3

CTRL

DATA3
DATA7
DATA6
DATA2
DATA4
DATA1
DATA5
DATA0
BUSY

DVDD
DGND
DGND

Next: how to calibrate a detector?

Additional material

Literature

<https://cds.cern.ch/record/2637873/files/CERN-THESIS-2018-153.pdf> Linxing Meng

- Slides from Erik Butz (see also [here](#)), [Simon Spannagel](#), [Freya Blekman](#), [Peter Schleper](#), [Erika Garutti](#), [Ingrid-Maria Gregor](#)
- Book on [Pixel detectors](#) by Rossi, Fischer, Rohe, Wermes
- Book on [Semiconductor Detector Systems](#) by Helmuth Spieler, see also [online notes](#)
- Particle data group review: [passage of radiation through matter](#)
- Book on [particle radiation](#) by Konrad Kleinknecht
- Book on [Techniques in experimental particle physics](#) by WR Leo
- Book on [The art of electronics](#) by Horowitz and Hill
- Book on [Evolution of silicon sensor technology in particle physics](#) by Frank Hartmann
- Book on [Physics of Semiconductor devices](#) by SM Sze and Kwok K Ng
- Book on [Semiconductor radiation detectors](#) by Gerhard Lutz

$$\bar{n}_i = \frac{1}{e^{(\varepsilon_i - \mu)/k_B T} + 1}$$

Do you remember?

$$N(E) = \frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{\frac{3}{2}} \sqrt{E - E_0}$$

$$\bar{n}_i = \frac{1}{e^{(\varepsilon_i - \mu)/k_B T} + 1}$$

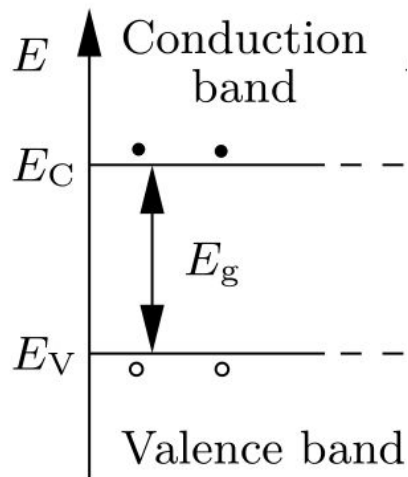
Fermi-Dirac

Do you remember?

Density of states

$$N(E) = \frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2} \right)^{\frac{3}{2}} \sqrt{E - E_0}$$

Intrinsic semiconductor



Intrinsic semiconductor

Concentration of acceptor/donor impurities N_A/N_D

\ll

thermally generated eh pairs

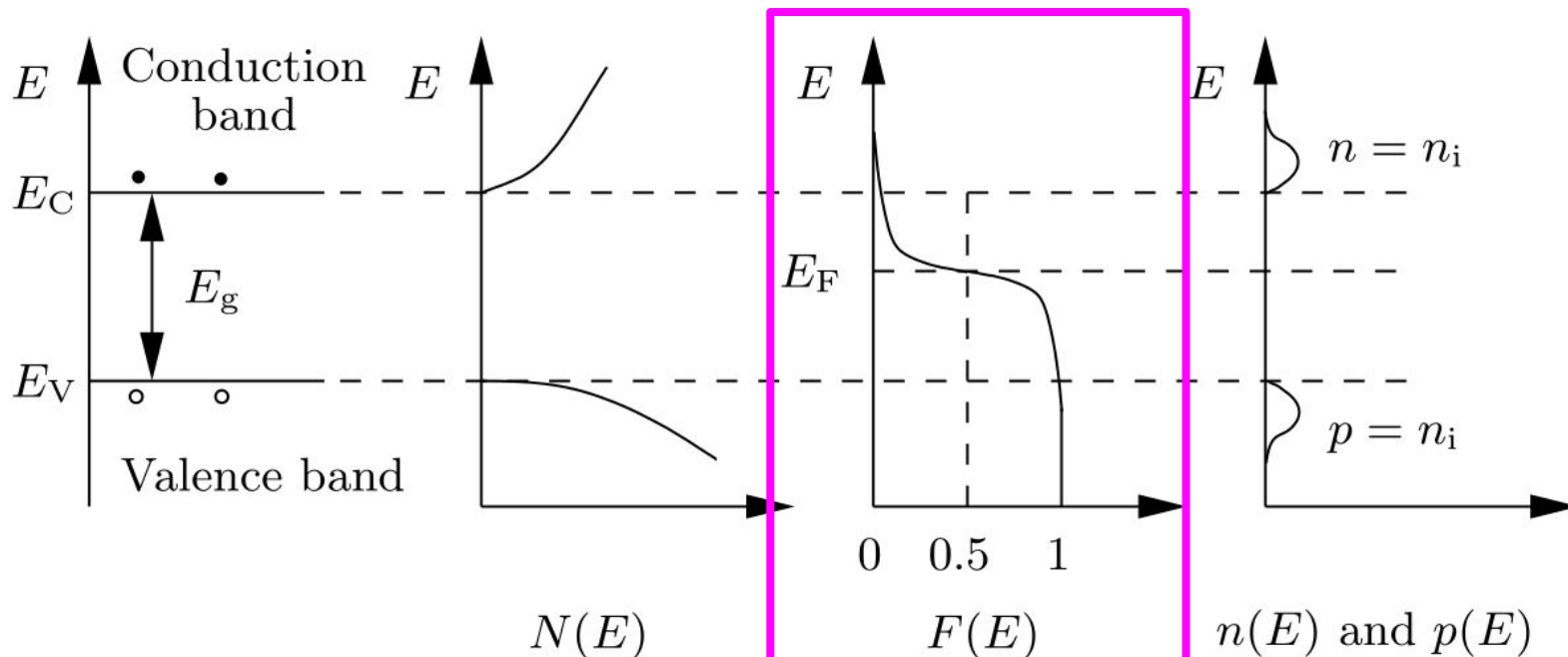
No doping yet!

$$np = n_i^2$$

Excited electrons n

Number of holes p

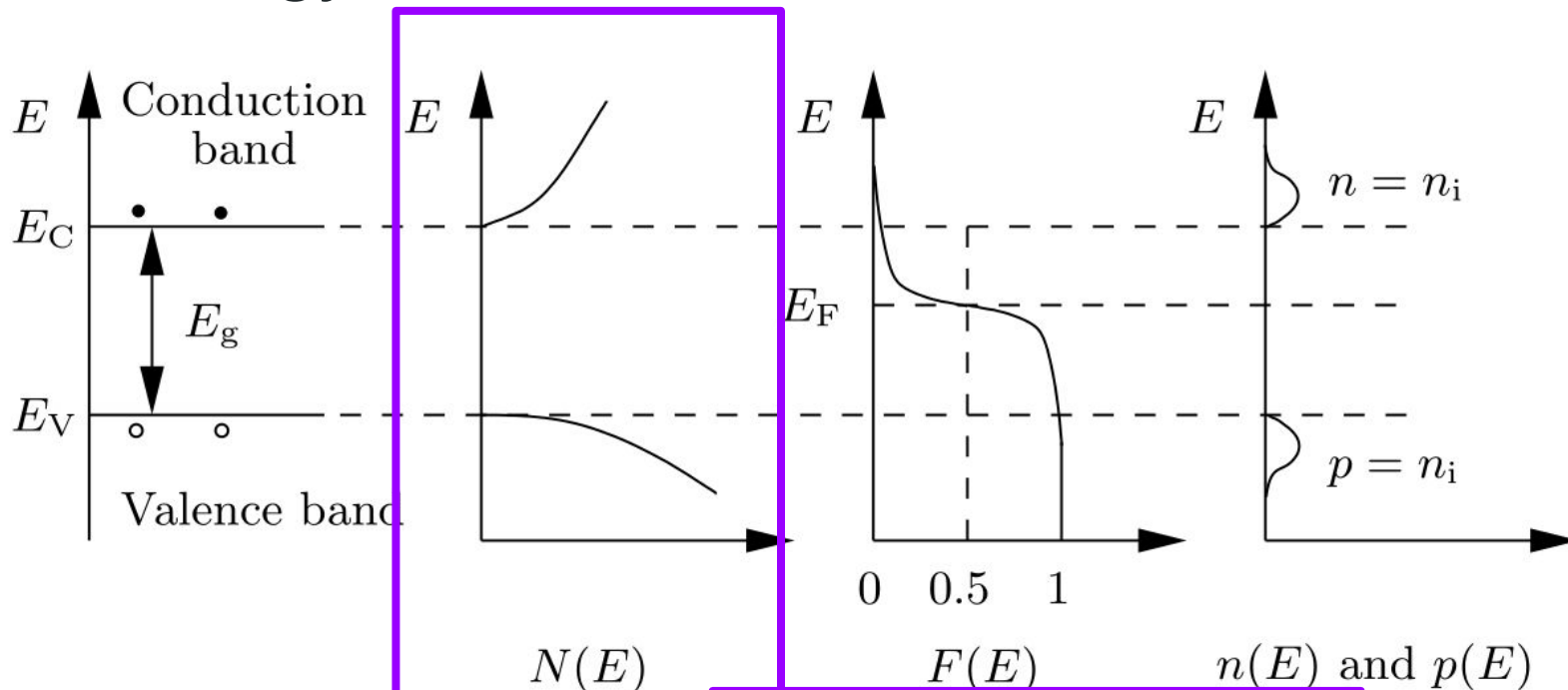
Fermi energy



Fermi function:

$$F(E) = \frac{1}{1 + e^{(E-E_F)/kT}} \approx e^{-(E-E_F)/kT}$$

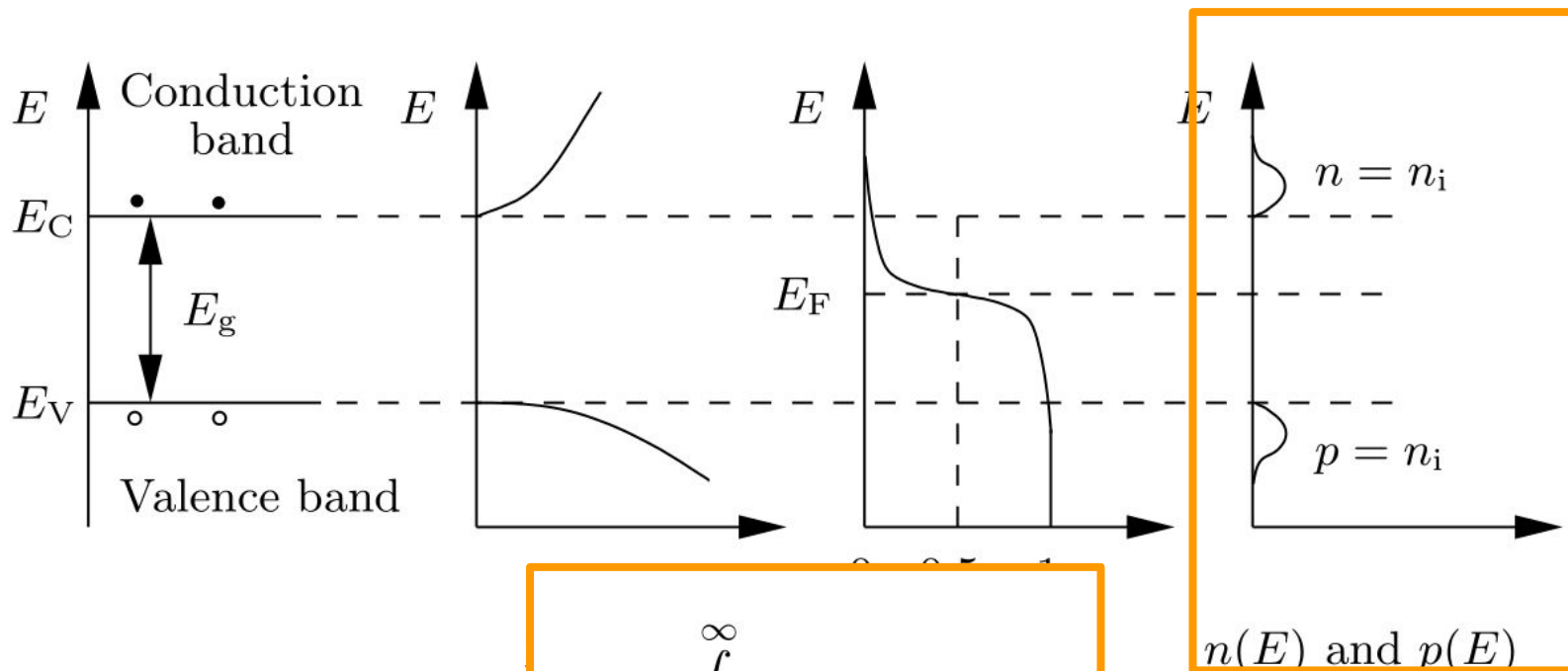
Fermi energy



Density of states

$$N(E) dE = 4\pi \left(\frac{2m_n}{h^2} \right)^{2/3} E^{1/2} dE$$

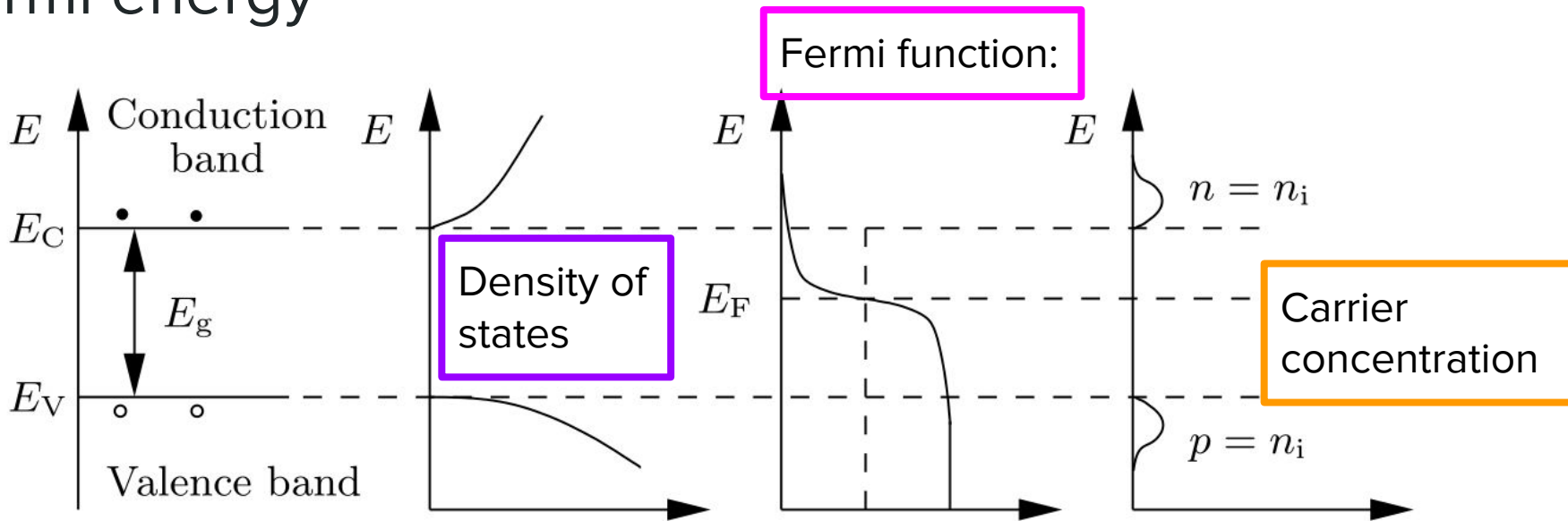
Fermi energy



Carrier concentration

$$n = \int_{E_C \equiv 0}^{\infty} N(E) F(E) dE$$

Fermi energy



Still no doping!

Band gap energy in silicon $E_g = 1.12$ eV

$$n_i = 1.45 \times 10^{10} \text{ cm}^{-3}$$

$$N(E) dE = 4\pi \left(\frac{2m_n}{h^2} \right)^{2/3} E^{1/2} dE$$

0 0.5 1

$$F(E) = \frac{1}{1 + e^{(E-E_F)/kT}}$$

$$\approx e^{-(E-E_F)/kT}$$

$$n = \int_{E_C \equiv 0}^{\infty} N(E) F(E) dE$$

Carrier concentration

$$n_i = N_c e^{-(E_c - E_F)/kT} = N_v e^{-(E_F - E_v)/kT}$$

The product of electron
and hole concentrations

depends
only on the
gap
energy



$$np = n_i^2 = N_c N_v e^{-(E_c - E_v)/kT} = N_c N_v e^{-E_g/kT}$$

The $np = n_i^2$ relationship is maintained: **law of mass action**

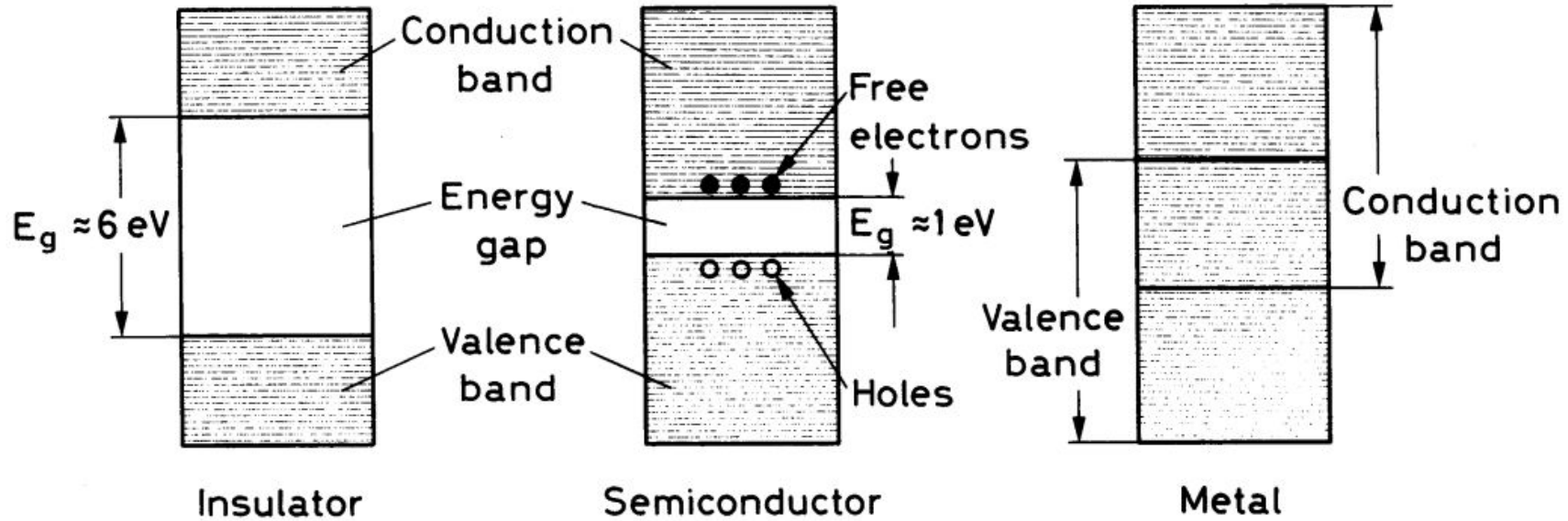
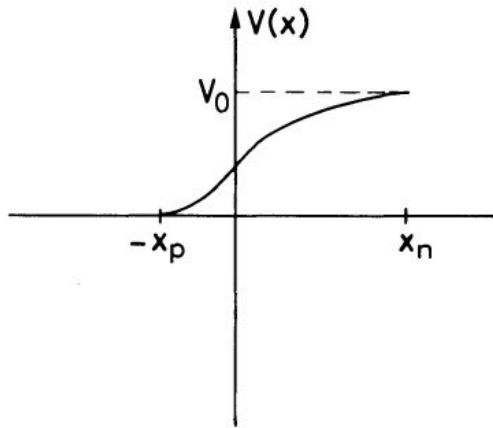
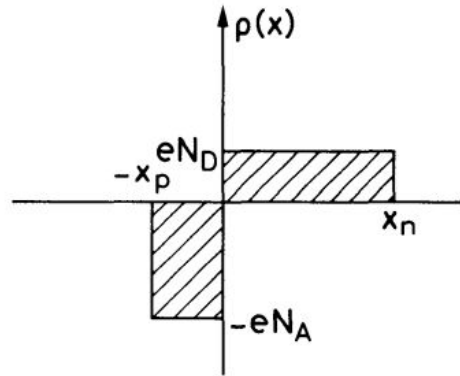


Fig. 10.1. Energy band structure of conductors, insulators and semiconductors



Poisson

$$\frac{d^2 V}{dx^2} = -\frac{\rho(x)}{\epsilon}$$

Electrostatics

$$\frac{dV}{dx} = \begin{cases} -\frac{eN_D}{\epsilon}(x-x_n) & 0 < x < x_n \\ \frac{eN_A}{\epsilon}(x+x_p) & -x_p < x < 0 \end{cases}$$

Actually, no!

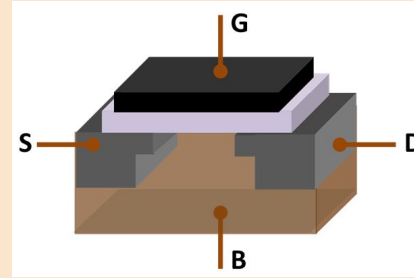
Until 1959

Germanium most common
Silicon worse conductivity
because of unstable
states at surface



1959 Mohamed Atalla and Dawon Kahng

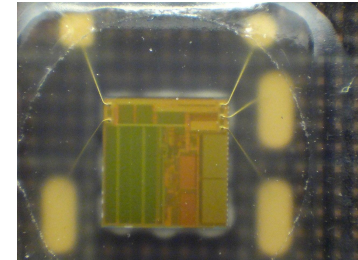
develop surface passivation: coat silicon surface with SO_2 :



Metal-oxide-semiconductor field-effect transistor

1959: first monolithic integrated circuit from silicon

1967: flash memory invented by Kahng and S. Sze → EEPROM



Inside
a sim
card

Silicon has a high quality, robust native oxide. This property is unique among semiconductors.