



Particle detection

Passage of radiation through matter

KATRIN: neutrino mass measurement

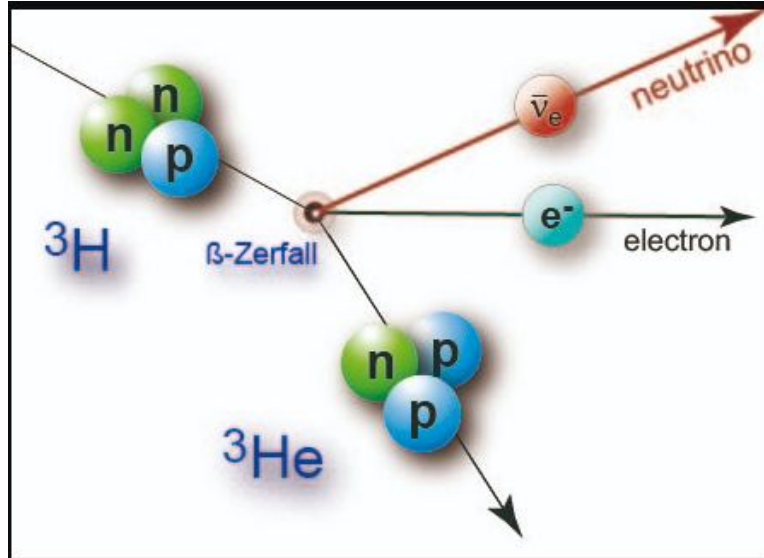


$m_\nu < 1.1$ eV (90% confidence level): twice as precise as previous measurements!

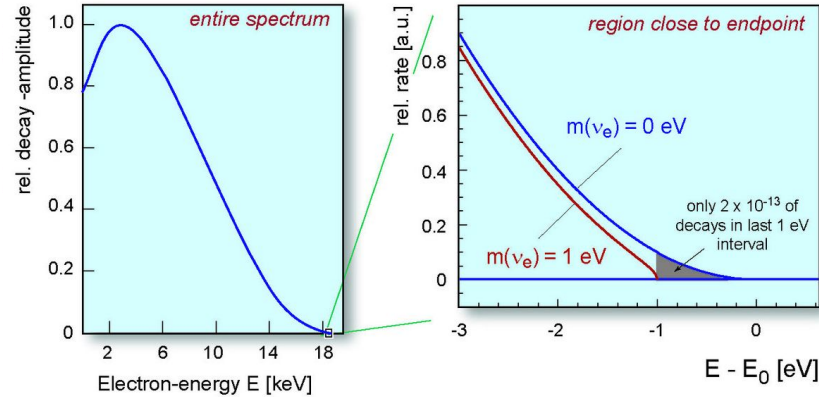
Recently published! <https://arxiv.org/abs/1909.06048>

KATRIN: neutrino mass measurement

Karlsruhe tritium neutrino experiment



$m_\nu < 1.1 \text{ eV}$ (90% confidence level)



Derive
neutrino mass
information
from electron
energy

Recently published!

<https://arxiv.org/abs/1909.06048>

https://www.katrin.kit.edu/img/spdctrum_rdxax_1200x678.jpg

Pre-work: before the lecture

Pre-work due before 2021-02-04

Watch: x-rays and proton beams for curing cancer

Curing Cancer with Proton Beams – with Suzie Sheehy

<https://www.youtube.com/watch?v=ZQ7kyocgiho>

Answer the following questions in the quiz: Why do protons stop and X-rays don't?
What are X-rays?

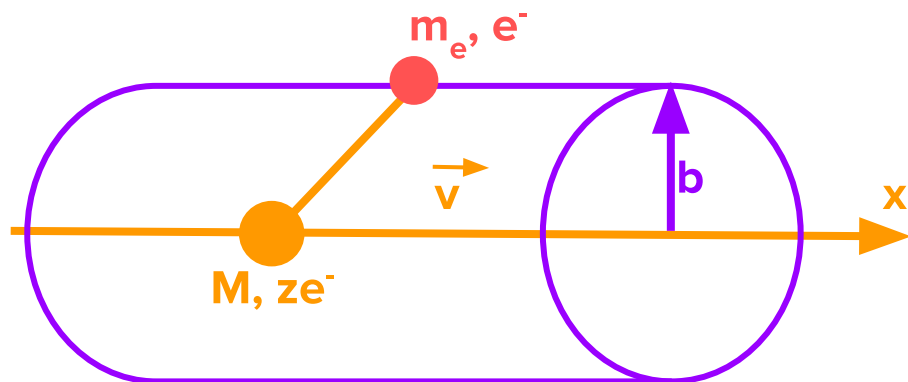
Compute the energy loss classically without quantum effects. Answer the question in the quiz: What is f_{avg} similar to in the [Bethe formula](#)?

See following slides for the exercises. This follows p22 [WR Leo](#).

Read from the book [Pixel detectors](#): p29-35 (2.2.2.1 and 2.2.2.2).

Classical energy loss: a collision

This exercise consists of 5 parts.



Gauss's Law:

$$\Phi_E = \oiint_S \mathbf{E} \cdot d\mathbf{A}$$

Bohr computed the energy loss of a heavy charged particle of mass M with charge ze^- in a collision with an electron. Using impulse

$$J = \int F dt$$

1. Show that the energy gained by the electron $\Delta E(b) = J^2/(2m_e)$ in this collision with the particle of mass $M \gg m_e$ is equal to

$$\Delta E(b) = 2z^2e^4/(m_e b^2v^2)$$

You can assume an infinite cylinder centered on the heavy particle trajectory with the radius passing through the position of the electron. 6

Classical energy loss: validity of formula

At very large b , the time is no longer short, so the impulse calculation is invalid. Similarly at $b=0$, the formula explodes.

2. Show that the maximum transferable energy assuming maximum relativistic kinetic energy transfer in a head-on collision with momentum conservation gives a minimum b : $b_{\min} = ze^2/(\gamma m_e v^2)$ with $\gamma = 1/(1-v^2/c^2)^{-1/2}$.

3. Show that assuming collision time $t \ll \tau = 1/f_{\text{avg}}$ for an average orbital electron frequency f_{avg} , which allows to assume the electron is at rest, gives a maximum b : $b_{\max} = \gamma v/f_{\text{avg}}$. This is a distance at which the minimum kinetic energy is transferred, $E_{\min} = I$ with I the mean ionization potential.

Classical energy loss: stopping power in material

4. Assuming a volume element of $dV = 2\pi b db dx$ and an electron density N_e , compute the energy lost to all electrons between b and db over a distance dx :

$$-dE(b) = \Delta E(b) N_e dV.$$

5. Write Bohr's classical formula for the energy loss or stopping power $-dE/dx$.

Solutions: will be given next week

Lecture on 2021-02-04 09:00

Particles from outer space

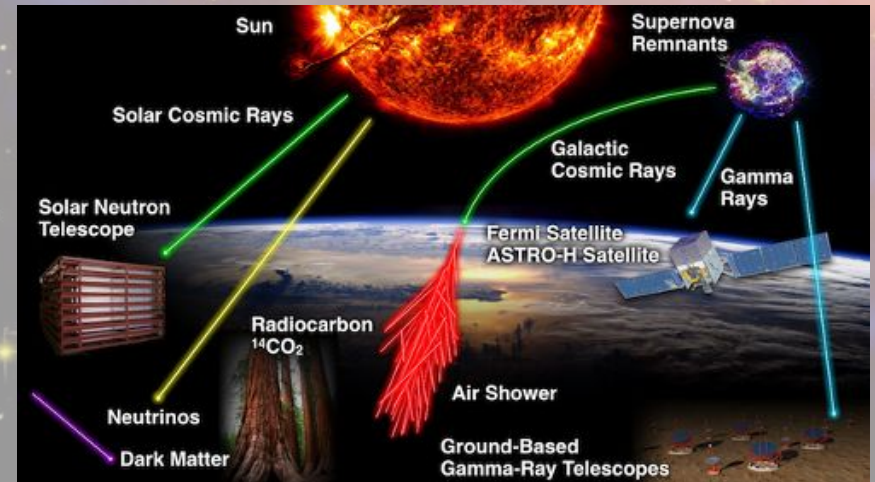
10000 times a second you have particles from cosmic rays passing through you



HESS: high energy stereoscopic system, in Namibia, can detect gamma rays

http://www2.cnr.fr/sites/en/image/hess_new_large_hd.jpg

What are these particles and how do they behave?
What are we and what is the universe made of?

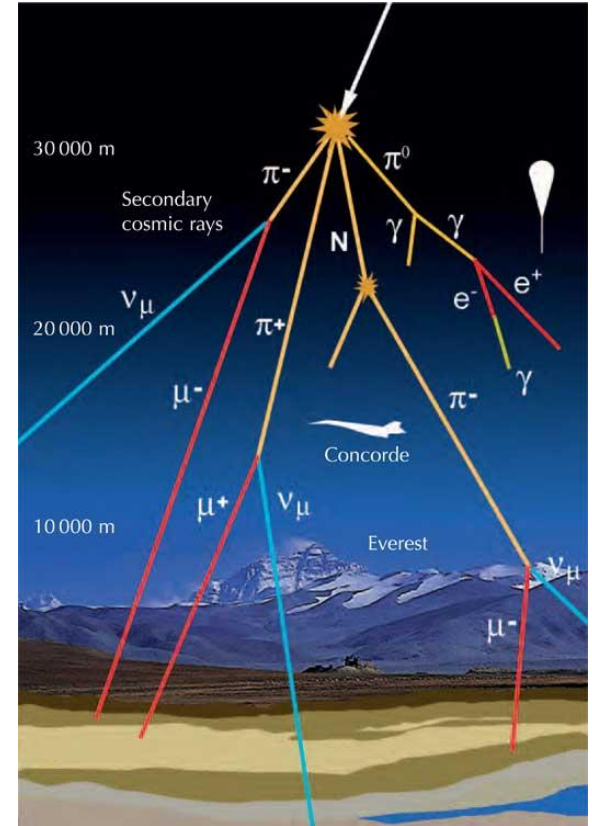


http://www.isee.nagoya-u.ac.jp/en/assets_c/2016/03/study01_1thumb-500xauto-153.png

What are cosmic rays?
What do we detect?

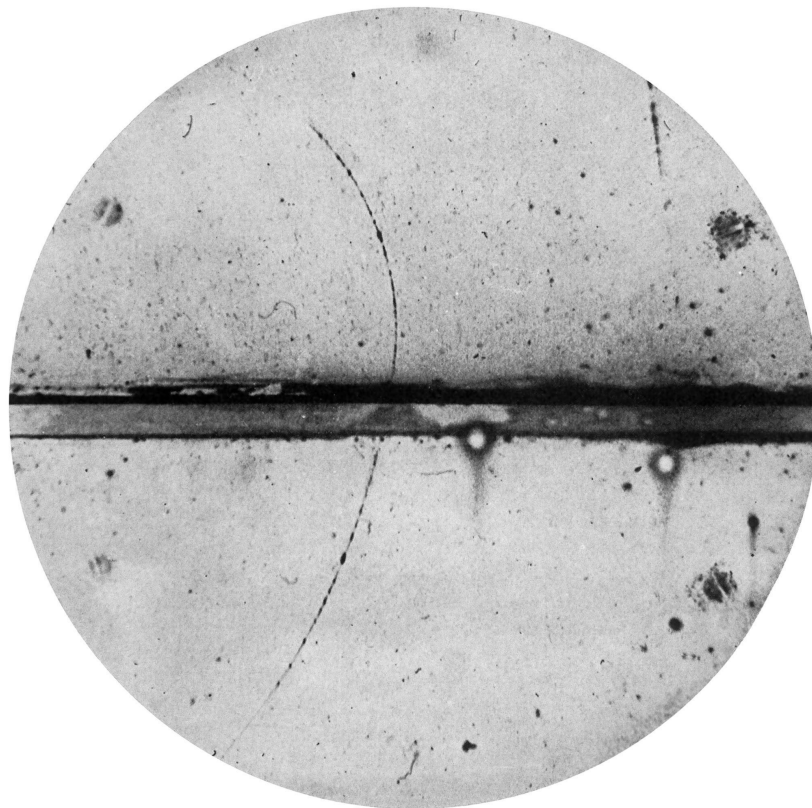
Passage of radiation through matter

Why do the muons reach the earth and the electrons do not?



To detect and identify particles, we need material and we do not want material

What do you see?



Discovery of antimatter

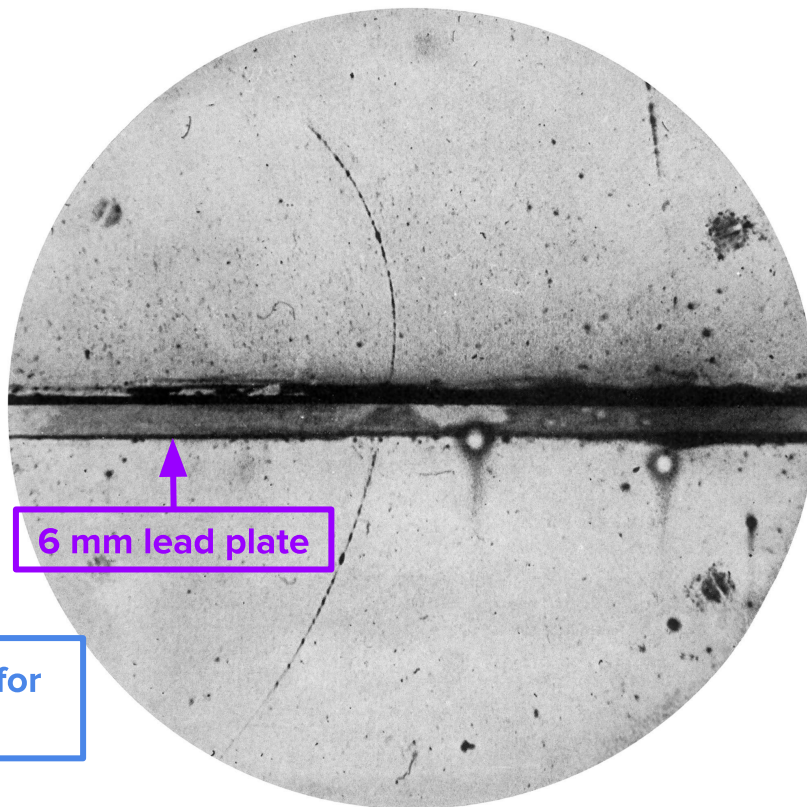
<https://upload.wikimedia.org/wikipedia/commons/6/69/PositronDiscovery.jpg>

C.D. Anderson <https://journals.aps.org/pr/pdf/10.1103/PhysRev.43.491>

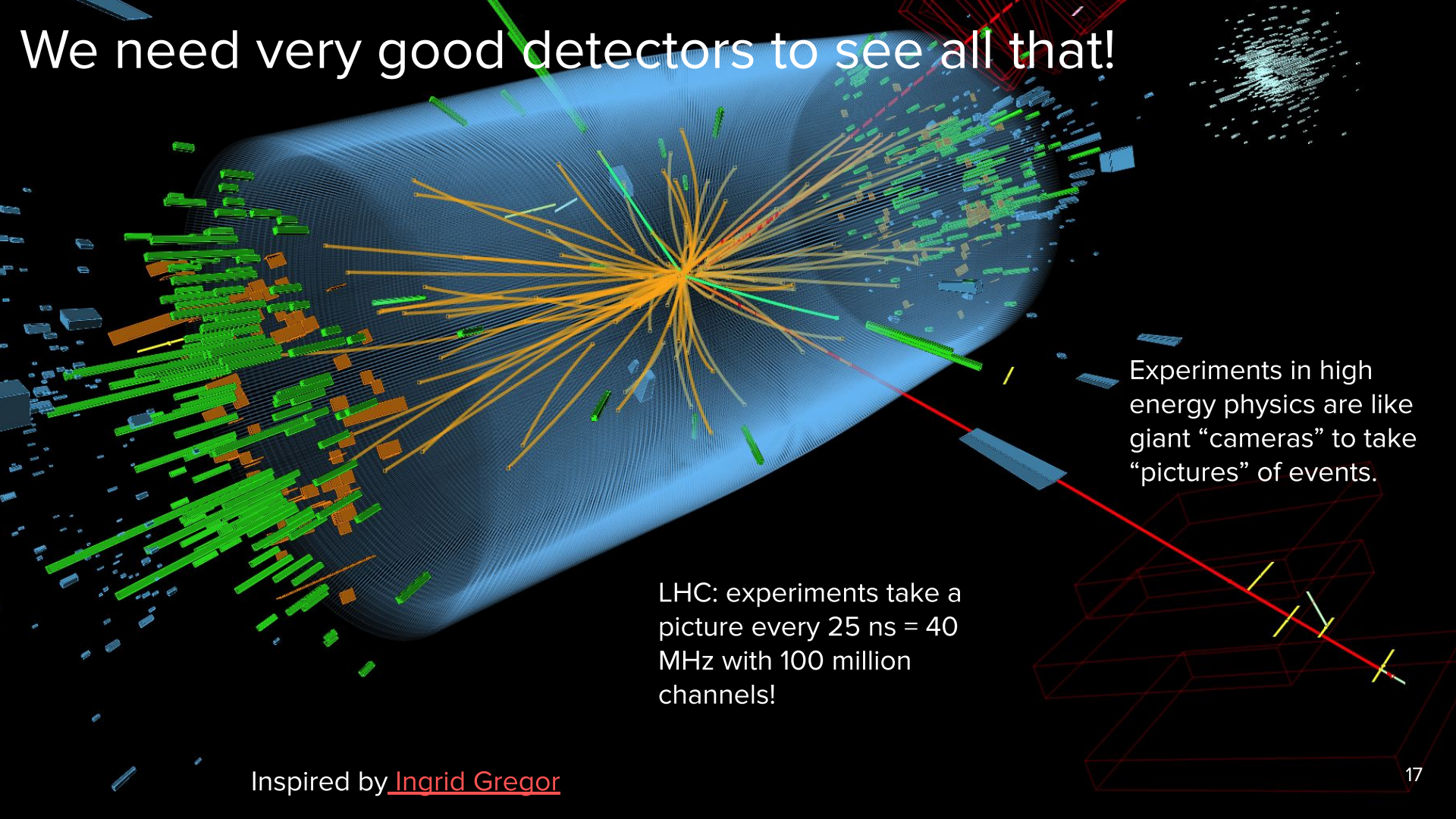
The first positron ever observed!

Wilson cloud chamber: gaseous mixture of supersaturated water or alcohol. Energetic particle ionizes gas and ions form condensation centers visible as a 'cloud'.

15000 Gauss = 1.5T magnetic field Wilson chamber for detecting cosmic rays



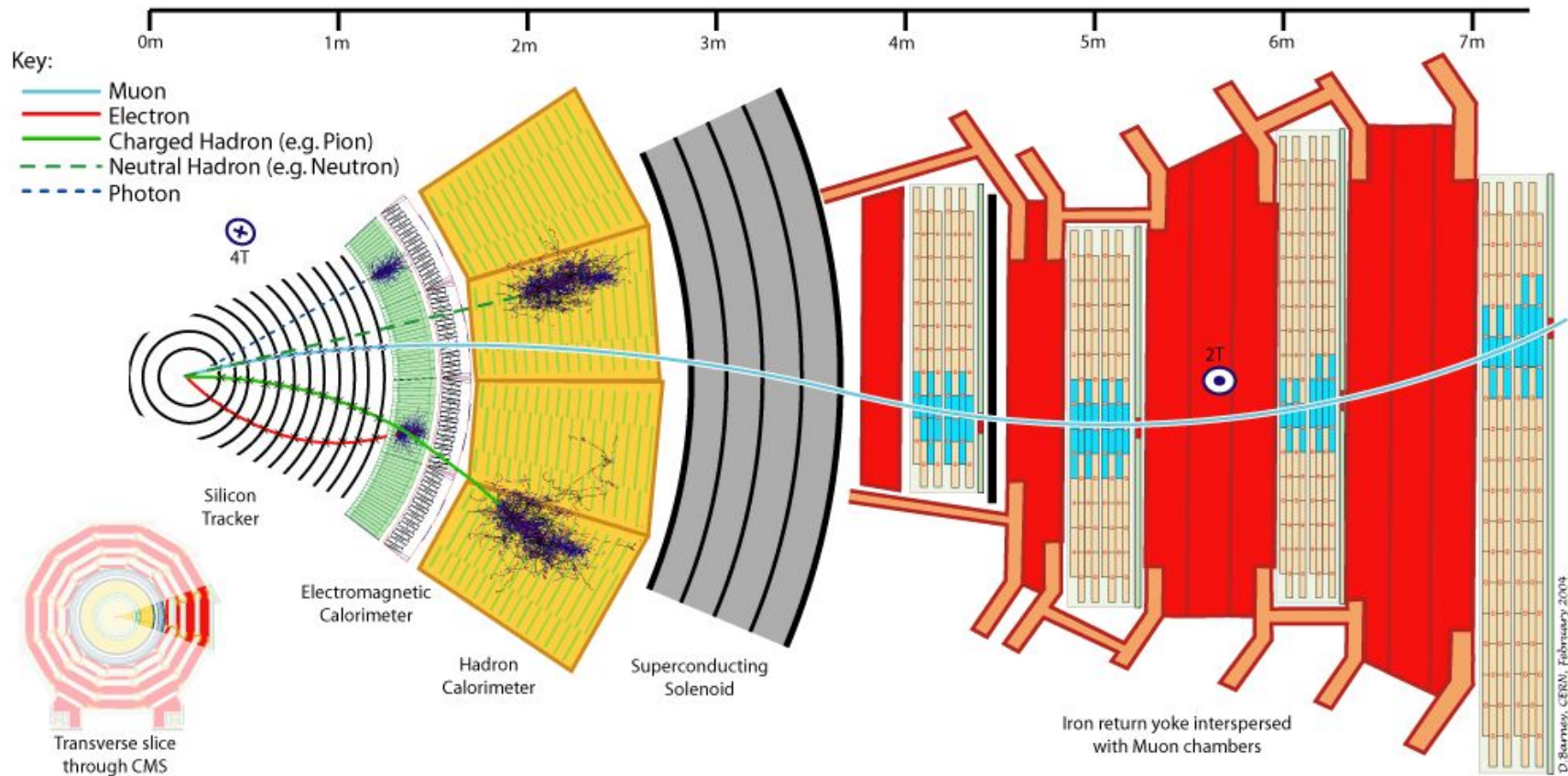
We need very good detectors to see all that!



Experiments in high energy physics are like giant “cameras” to take “pictures” of events.

LHC: experiments take a picture every 25 ns = 40 MHz with 100 million channels!

What do you notice in the particle behavior here?



To detect and identify particles, we use their interaction with matter and we do not want too much material



d=15m

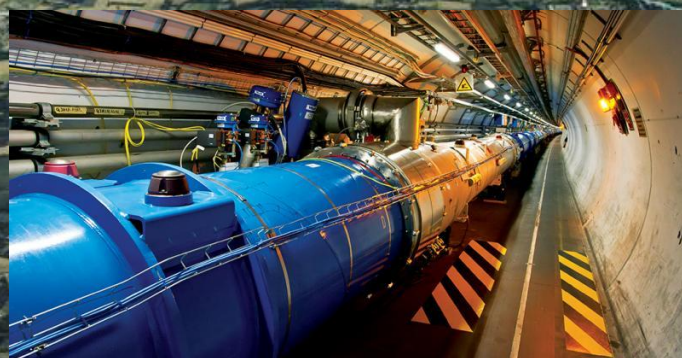
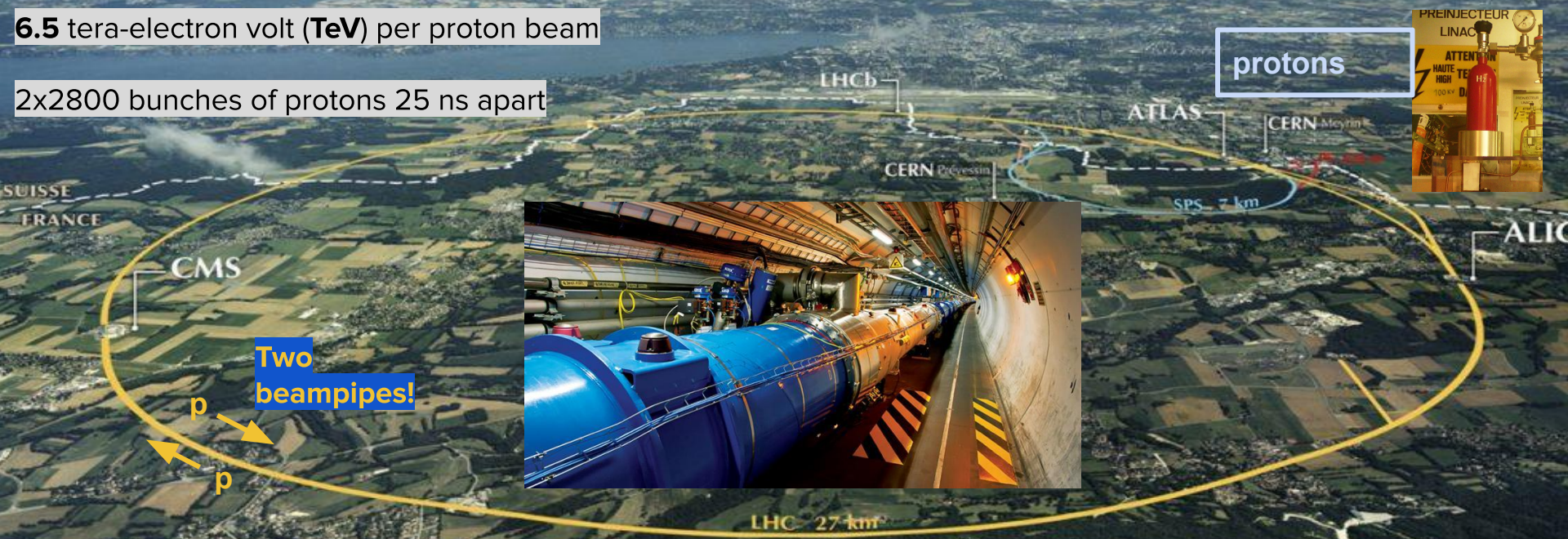
100 meters underground in the CMS cavern

Large Hadron Collider

Hadron: composite particle made of quarks held together by the strong force

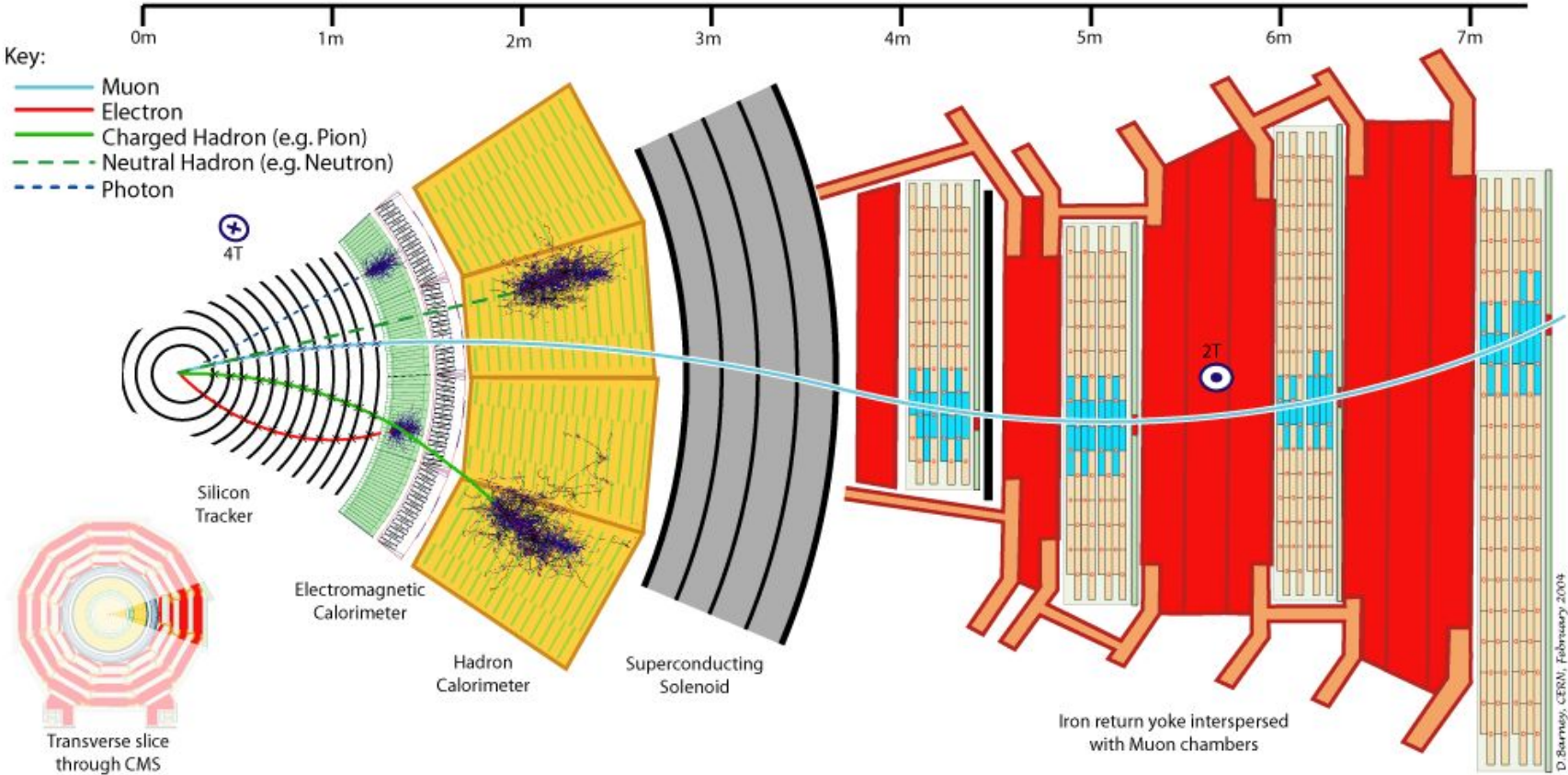
6.5 tera-electron volt (**TeV**) per proton beam

2x2800 bunches of protons 25 ns apart

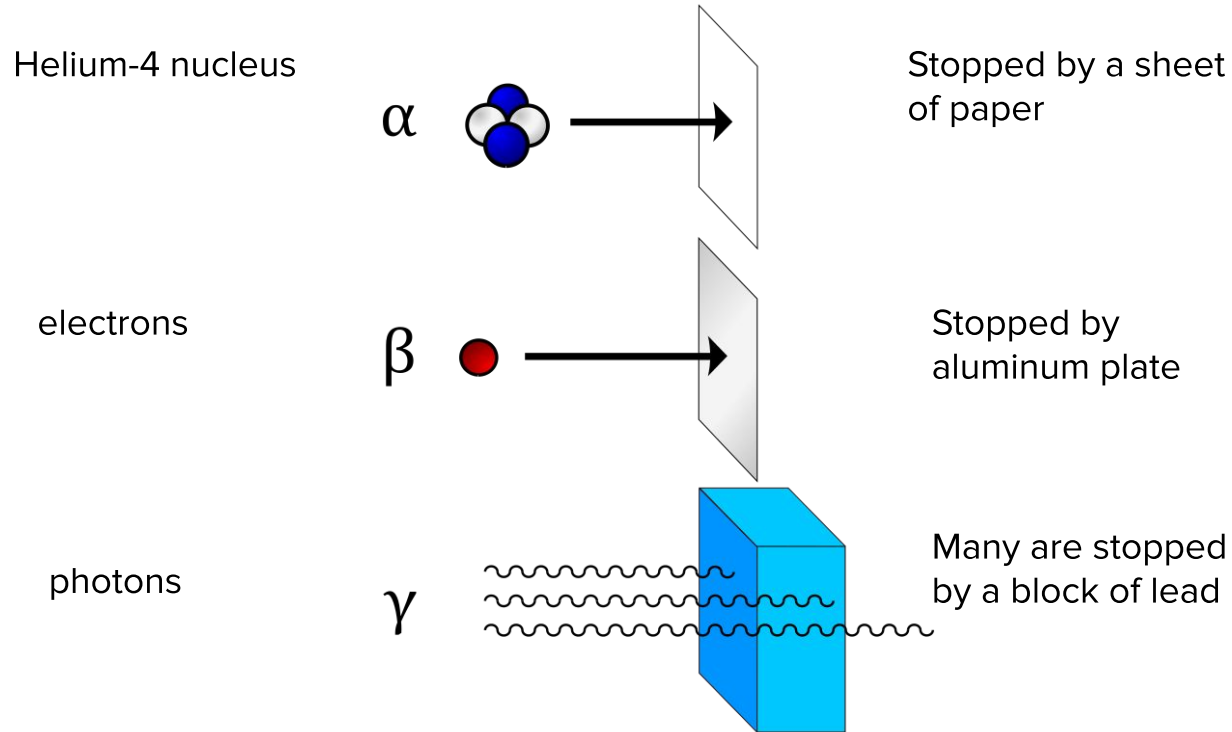




Different detector types for different measurements



Passage of radiation through matter



What do we detect?

How can we detect these particles?

Directly detect:

Decay products

jets

Indirectly detect:

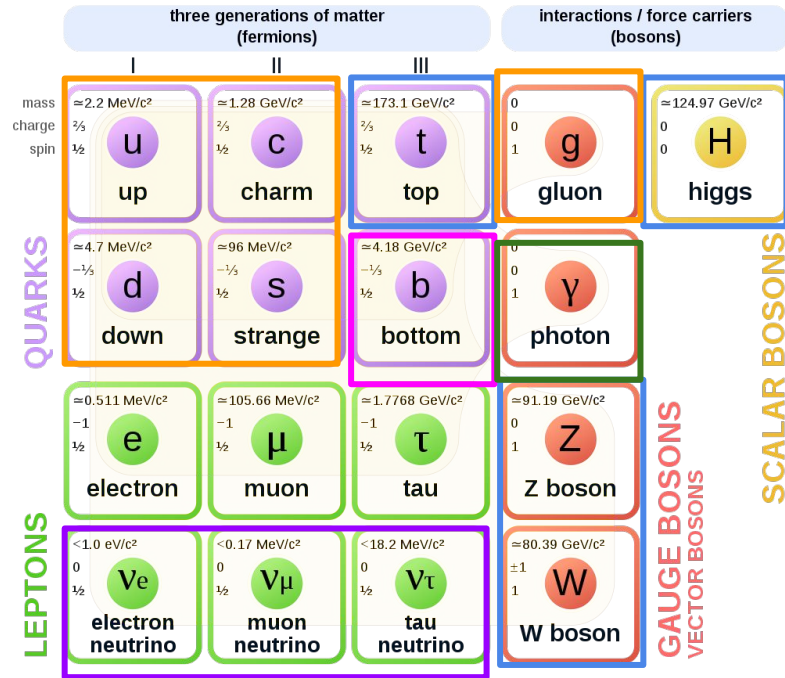
Missing energy

Secondary vertex + jets

Neutral particles

Should be able to detect and identify:
 $e^\pm, \mu^\pm, \gamma, \pi^\pm, K^\pm, p^\pm, K^0, n$
 using mass, charge, interaction

Standard Model of Elementary Particles



What do we measure and how?

| Observable | Measurable quantity |
|---------------------|--|
| Momentum (p) | Bending radius in magnetic field |
| Speed (v) | Time of flight, Cherenkov radiation |
| Charge (Q) | Bending in magnetic field |
| Lifetime (τ) | Distance traveled before decay |
| Energy (E) | Absorption in calorimeters |
| Mass (m) | Indirectly from momentum |
| Spin | <u>Angular distributions</u> |

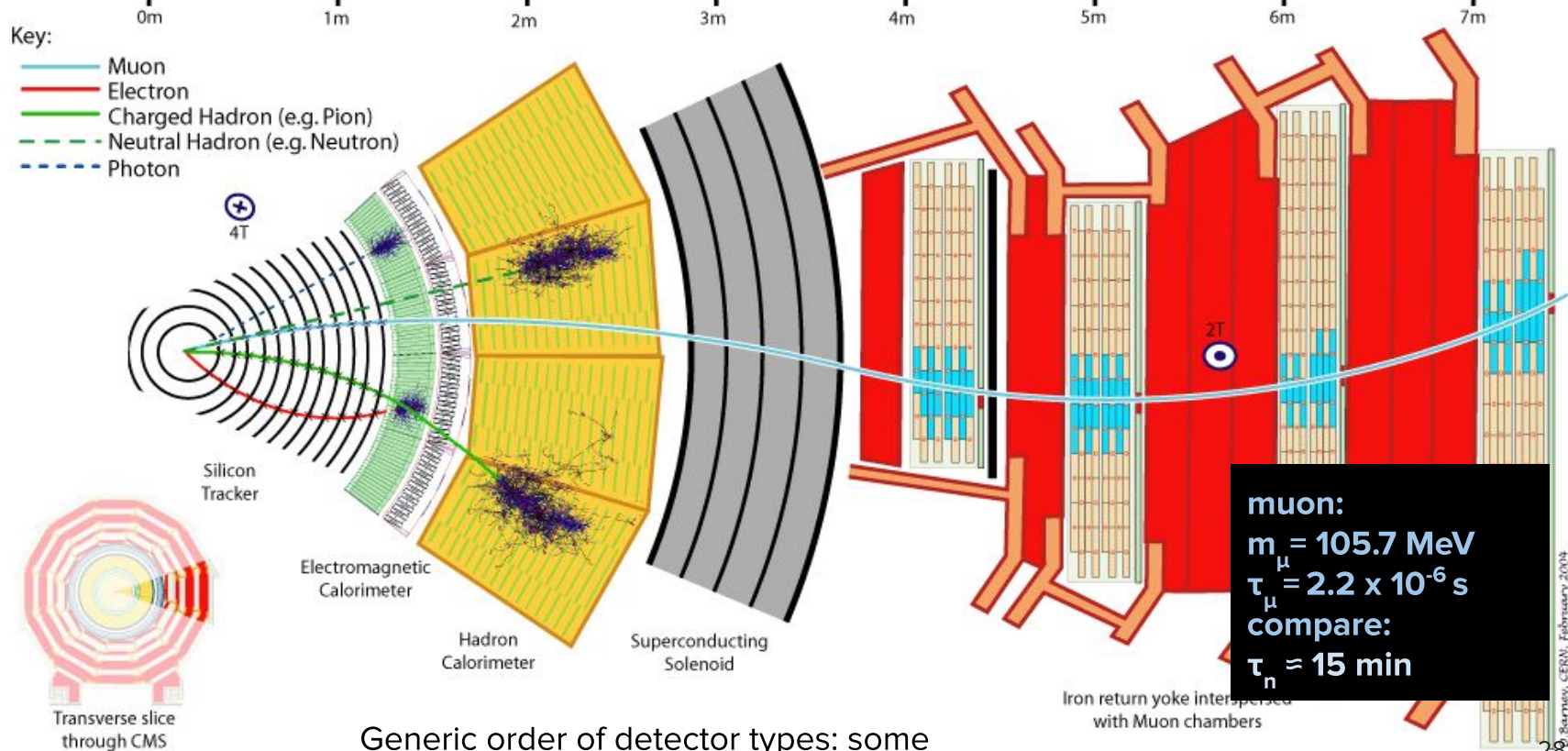
- $d = c\tau\gamma$
- $\gamma = 1/\sqrt{1-\beta^2}$
- $\beta = v/c$
- $E^2 = m^2c^4 + p^2c^2$
- $p = \gamma mv = mv/\sqrt{1-v^2/c^2}$

For some examples of measuring spin see

<https://arxiv.org/pdf/1202.6660.pdf> and
<http://moriond.in2p3.fr/QCD/2013/proceedings/Muehleitner.pdf>

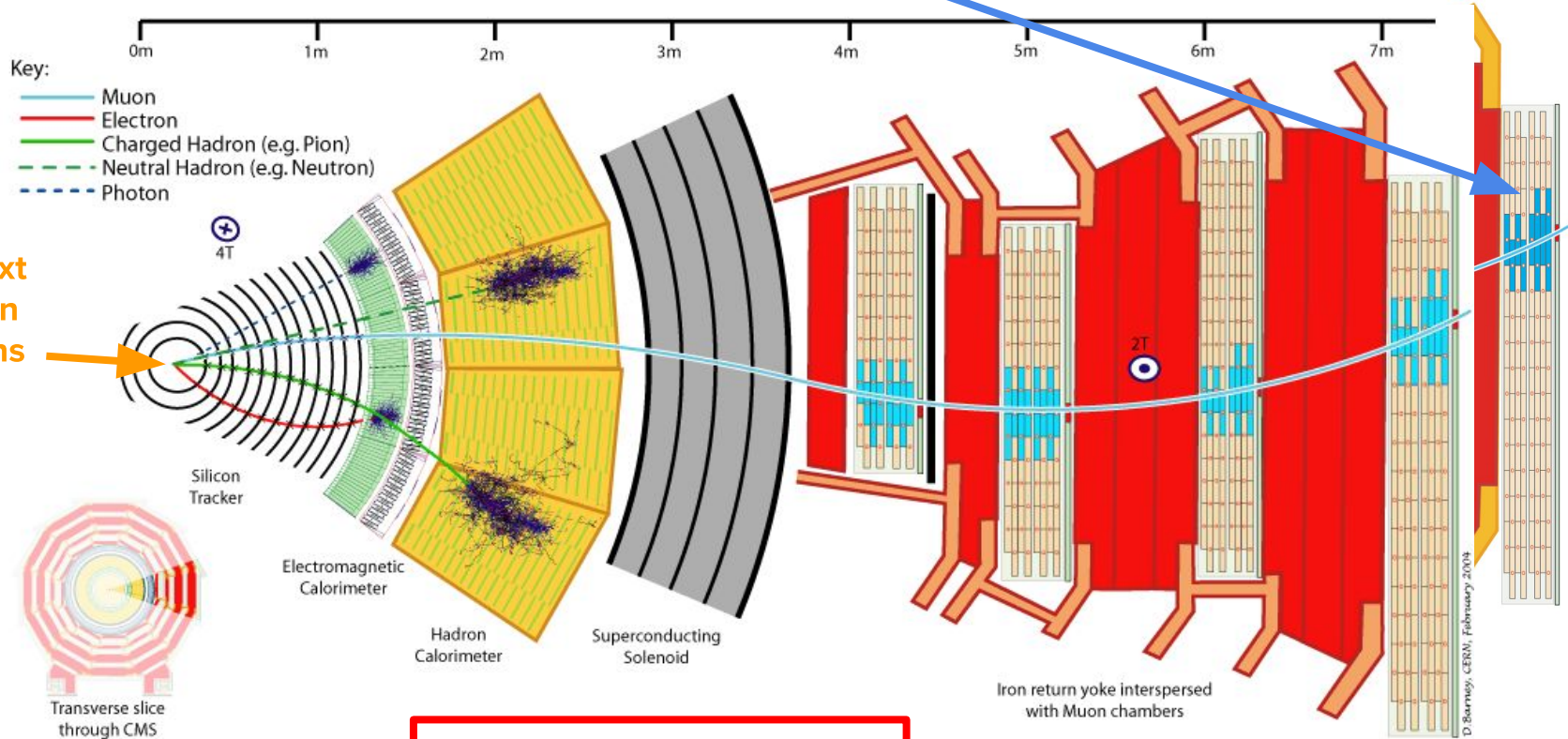
Need 1) a magnetic field and 2) interaction with material

Detectors at the large hadron collider: onion-like



Generic order of detector types: some measurements destructive!

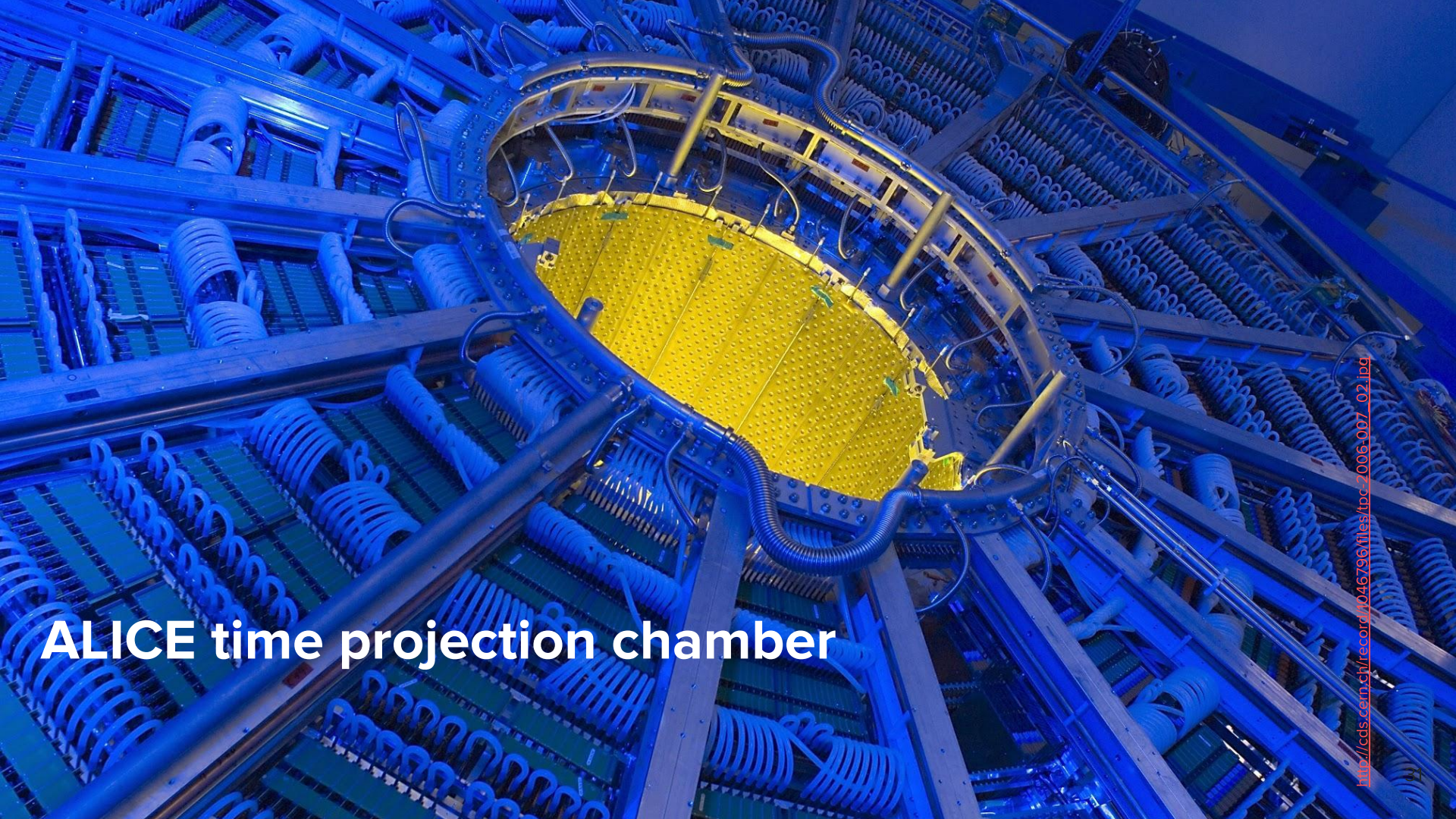
Note when the muon arrives here



The next collision happens here:

$$25 \text{ ns} \cdot c \approx 7.5 \text{ m}$$

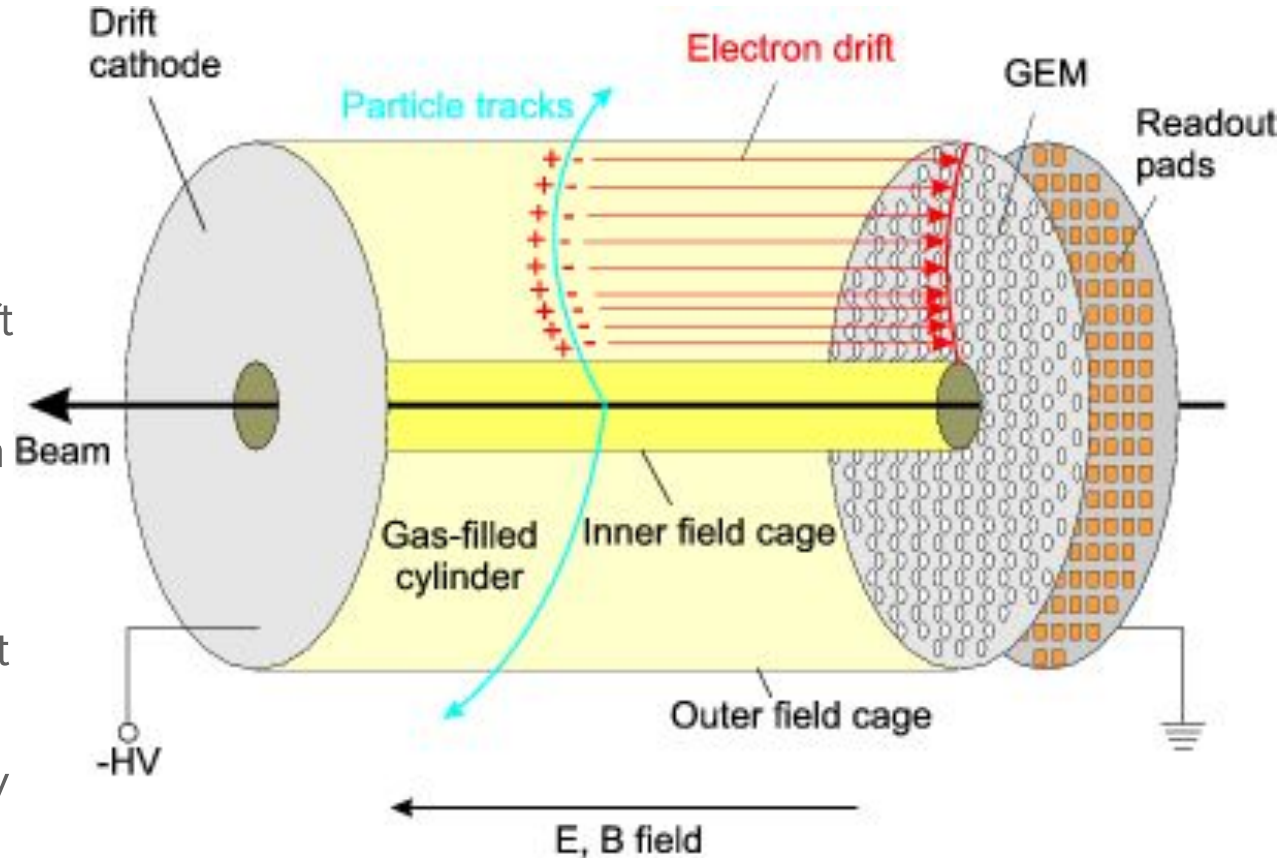
Ionization of a gas



ALICE time projection chamber

TPC

1. Ionization of gas in **chamber** with electric field causes electron drift
2. Signal gets amplified, in this case by gas electron multipliers \rightarrow electron avalanche
3. Readout pads can detect signal that can be **projected** onto trajectory
4. z (along beam) information from **timing**



Energy loss in matter

Bohr vs Bethe formula for energy loss

Bethe:

$$-\left\langle \frac{dE}{dx} \right\rangle = \frac{4\pi}{m_e c^2} \cdot \frac{nz^2}{\beta^2} \cdot \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \cdot \left[\ln \left(\frac{2m_e c^2 \beta^2}{I \cdot (1 - \beta^2)} \right) - \beta^2 \right]$$

Electron density n (N_e):

$$n = \frac{N_A \cdot Z \cdot \rho}{A \cdot M_u}$$

$$I = hf_{\text{avg}}$$

Bohr (from your pre-work):

$$-\frac{dE}{dx} = \frac{4\pi z^2 e^4}{m_e v^2} N_e \ln \frac{\gamma^2 m_e v^3}{ze^2 f_{\text{avg}}}$$

ρ is the density of the material, Z its [atomic number](#), A its [relative atomic mass](#), N_A the [Avogadro number](#) and M_u the [Molar mass constant](#).

Bohr vs Bethe formula for energy loss

Bethe:

$$-\left\langle \frac{dE}{dx} \right\rangle = \frac{4\pi}{m_e c^2} \cdot \frac{nz^2}{\beta^2} \cdot \left(\frac{e^2}{4\pi\epsilon_0} \right)^2 \cdot \left[\ln \left(\frac{2m_e c^2 \beta^2}{I \cdot (1 - \beta^2)} \right) - \beta^2 \right]$$

Dominates for lower non-relativistic energies: dE/dx decreases until a minimum is reached at around $v=0.96c$.

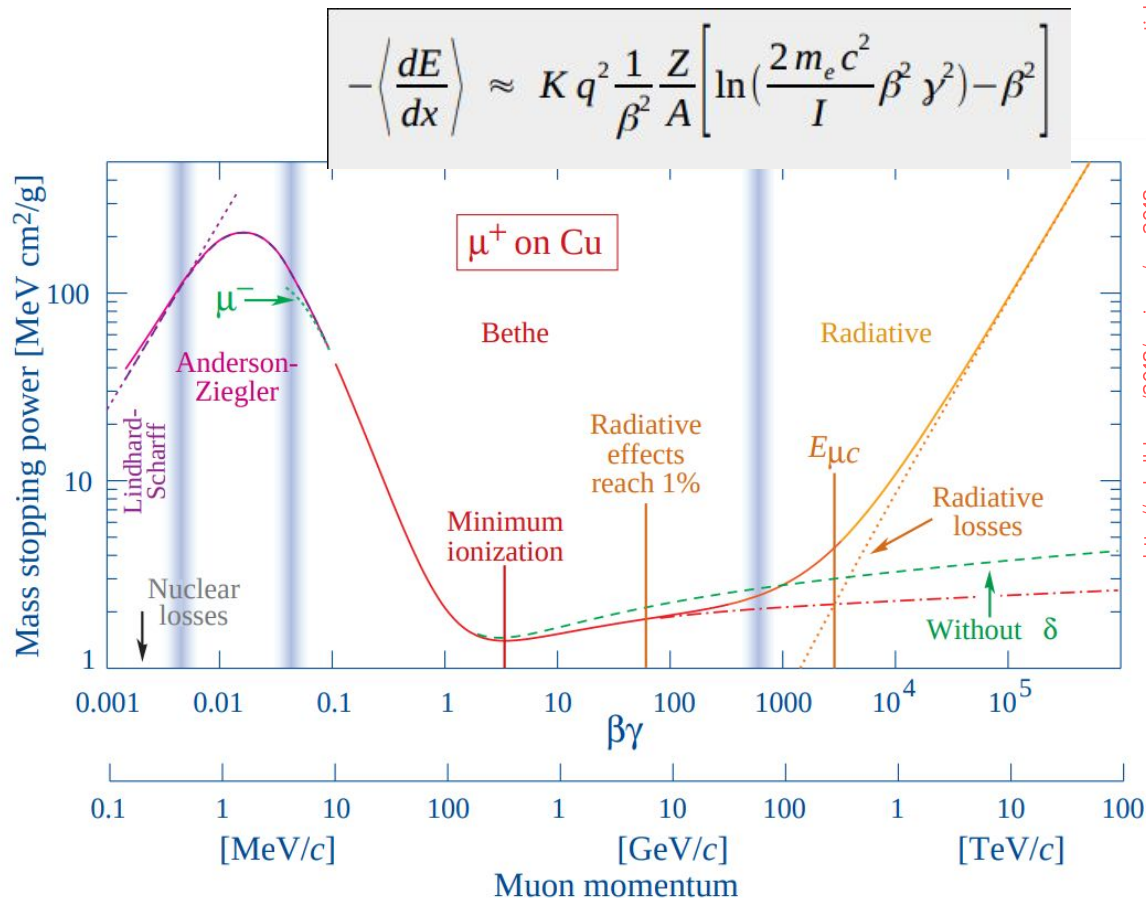
At much larger energies: dE/dx rises again due to log term -- save a **density**

correction: outer electrons are shielded by inner electrons polarized by electric field of incoming particle

Ionization loss

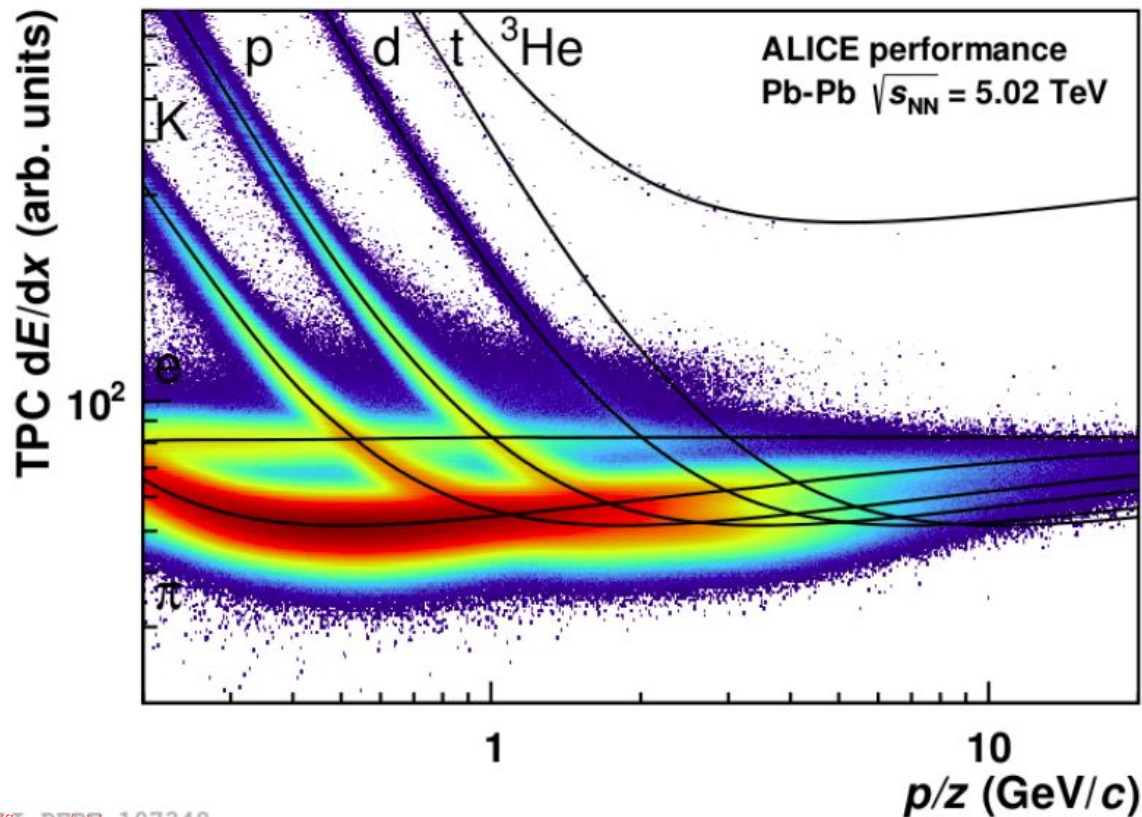
- Can measure ionization loss dE/dx
- K is a coefficient:
 $K = .307 \text{ MeV mol}^{-1}\text{cm}^{-2}$
- I is the mean excitation energy

Depends on charge, atom number, ionization energy, density



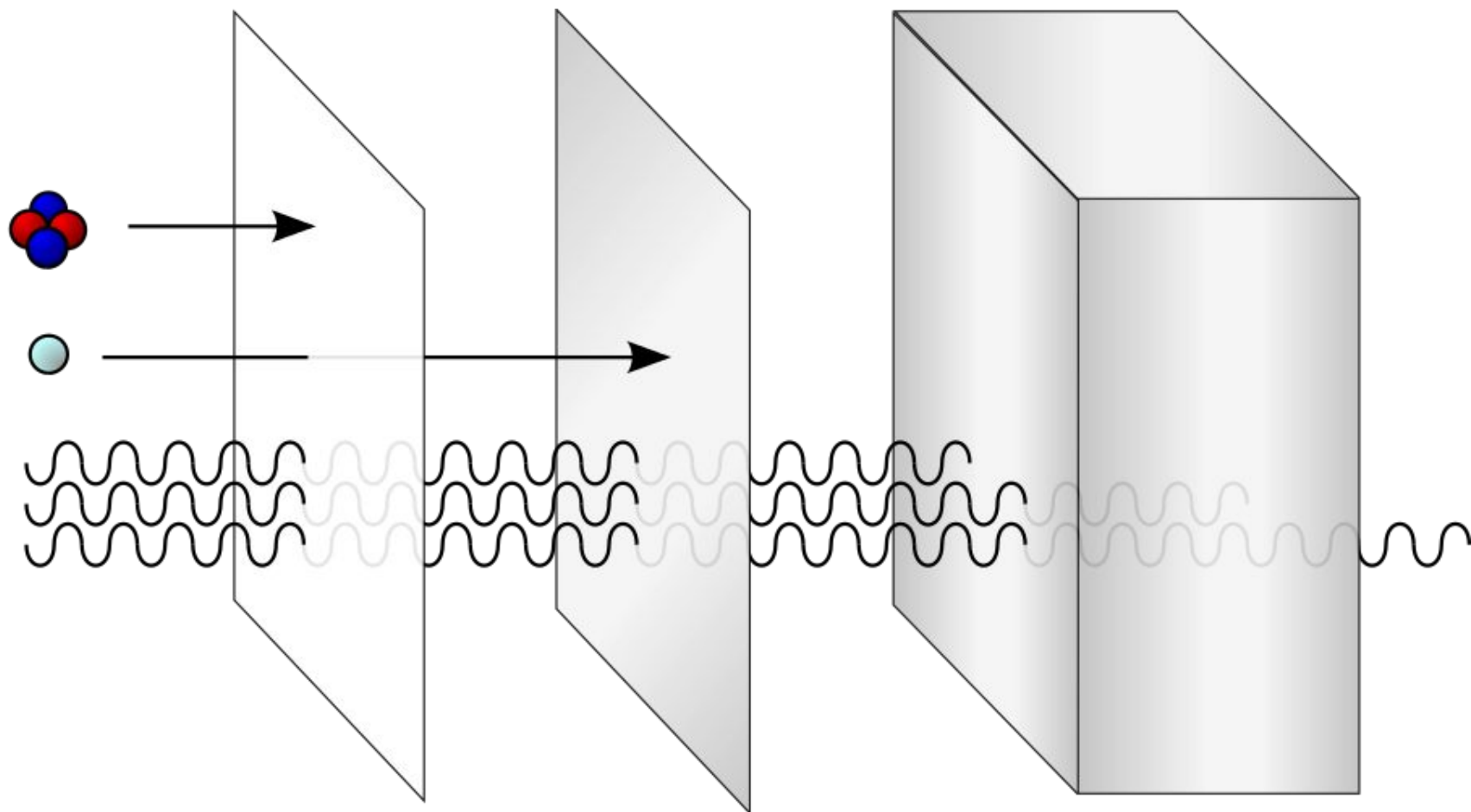
Identification of particles with the ALICE TPC

- Every point is one measurement!
- Can identify particles for low momenta
- For higher momenta, all particles behave like a minimum ionizing particle (MIP)



We use the interaction of particles with material to detect, measure and identify them

α
 β
 γ



Paper

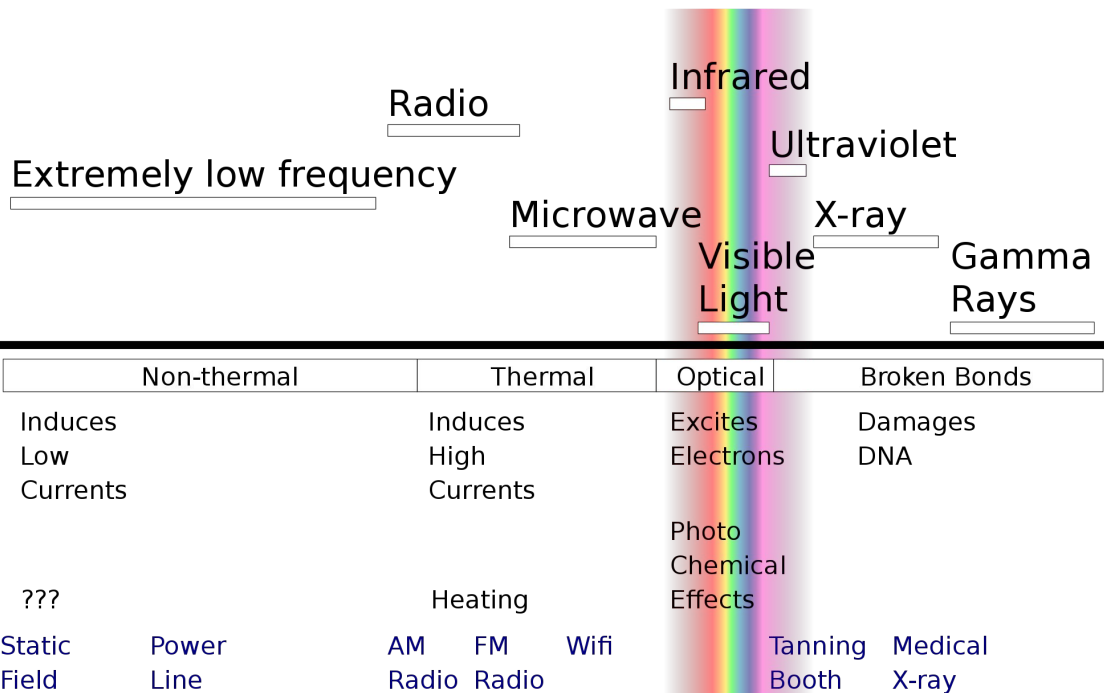
Aluminium

Lead

Ionizing vs non-ionizing radiation

Non-ionising

Ionising



Radiation: energy in the form of waves or particles

Ionizing radiation: radiation that, when passing through matter, strips electrons off atoms or molecules. **This can be damaging to living tissue.**

Types of ionizing radiation

Naturally:

- Cosmic rays
- Beta decay
- Positrons
- Electrons
- Protons
- Alpha particles (helium nuclei)

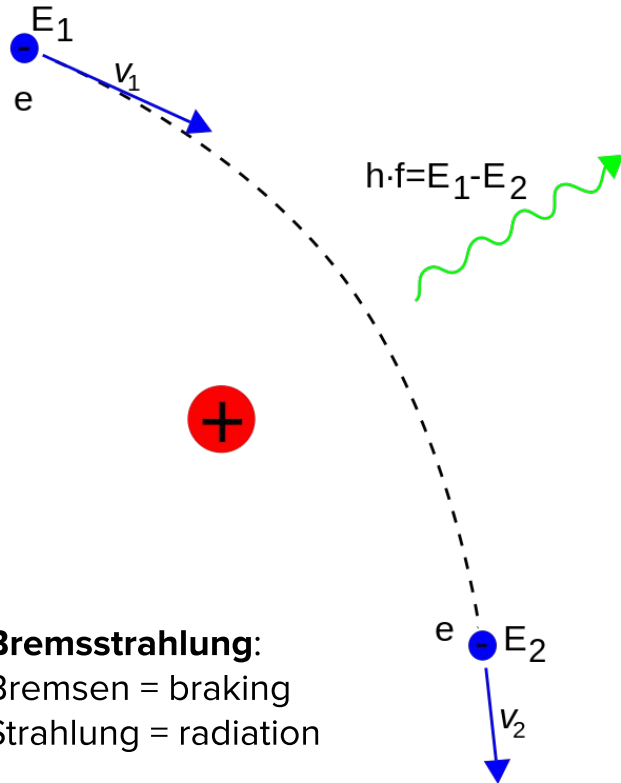
Artificially:

- X-rays
- Particle accelerators

Direct from charged particles:

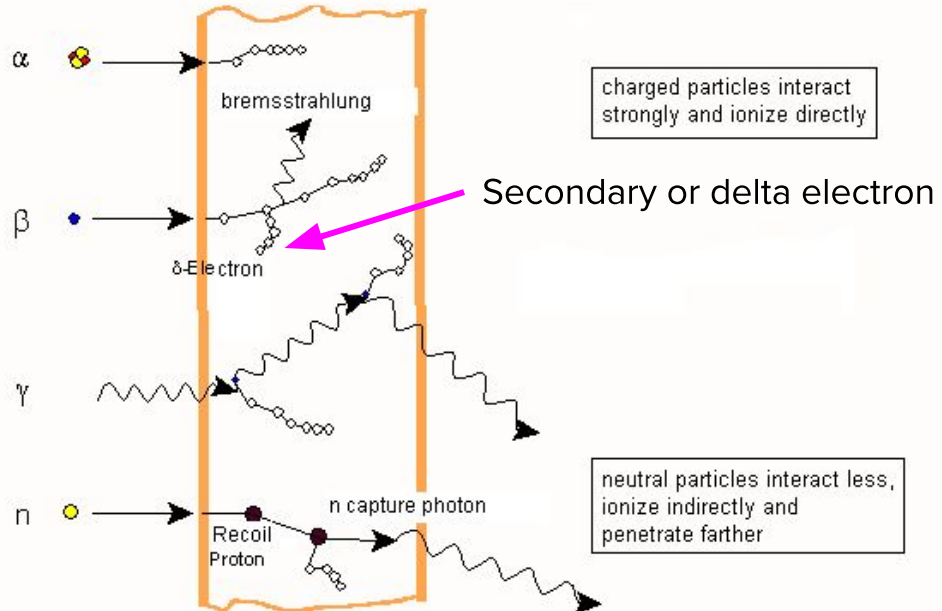
- Heavy charged particles: elastic scattering
- Electrons: Brehmsstrahlung (= braking radiation)
- Electrons: secondary or delta-electrons
- Positrons: electron-positron annihilation
- ...

Stopping particles

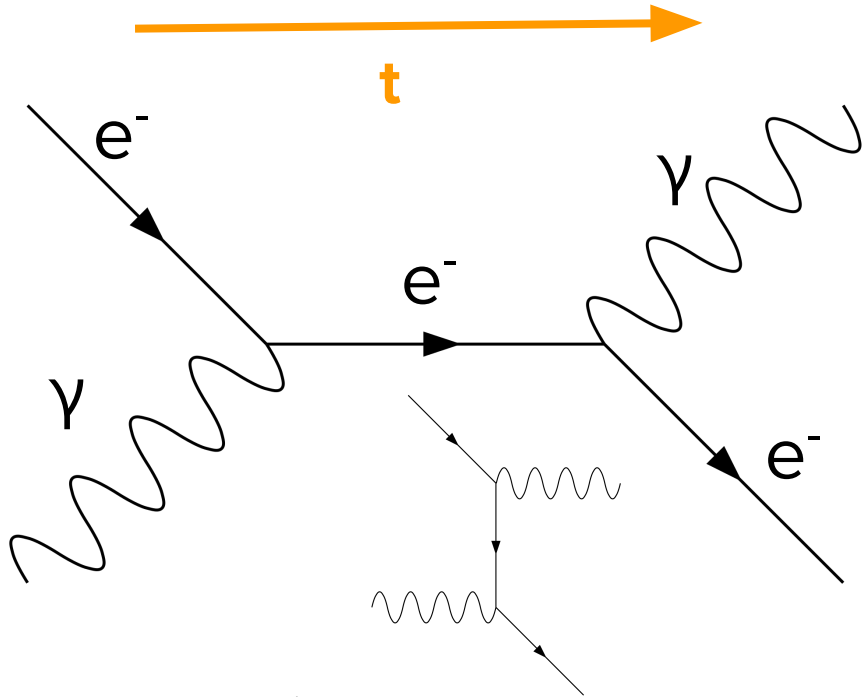


Bremsstrahlung:
Bremsen = braking
Strahlung = radiation

Interaction of ionizing Radiation with Matter

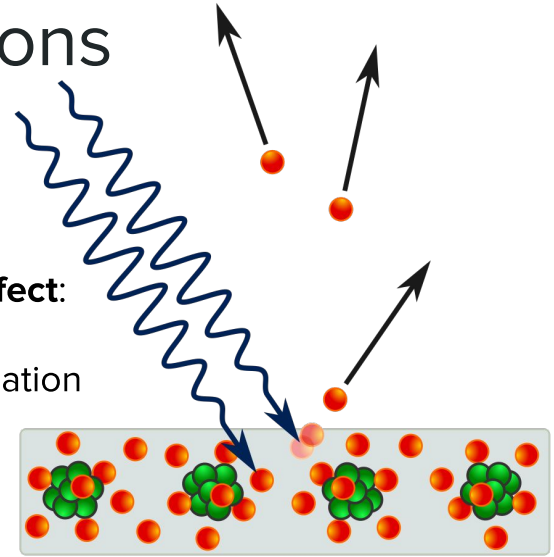


Indirect ionizing radiation from photons

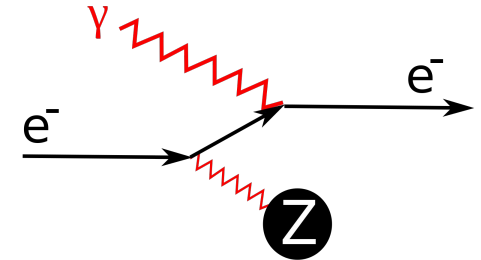


Compton scattering of a photon and electron. To first order only.

Photoelectric effect: absorption of a photon and ionization of an atom.

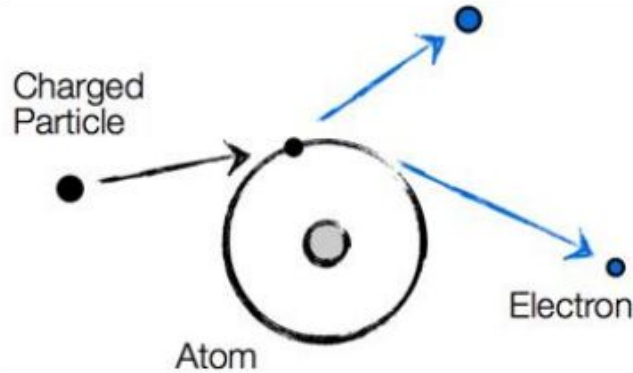


Or: an electron interacting with a photon and thereby changing its energy level in the atom Z

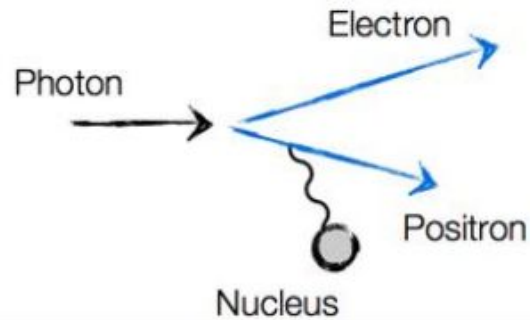


Interaction with matter

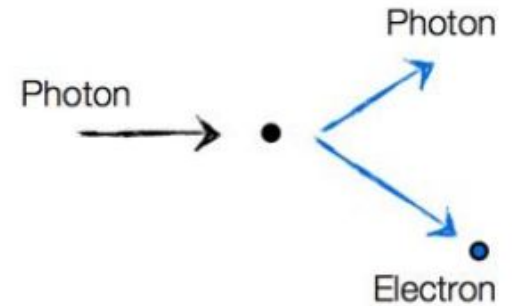
ionisation



Electron-positron
pair production

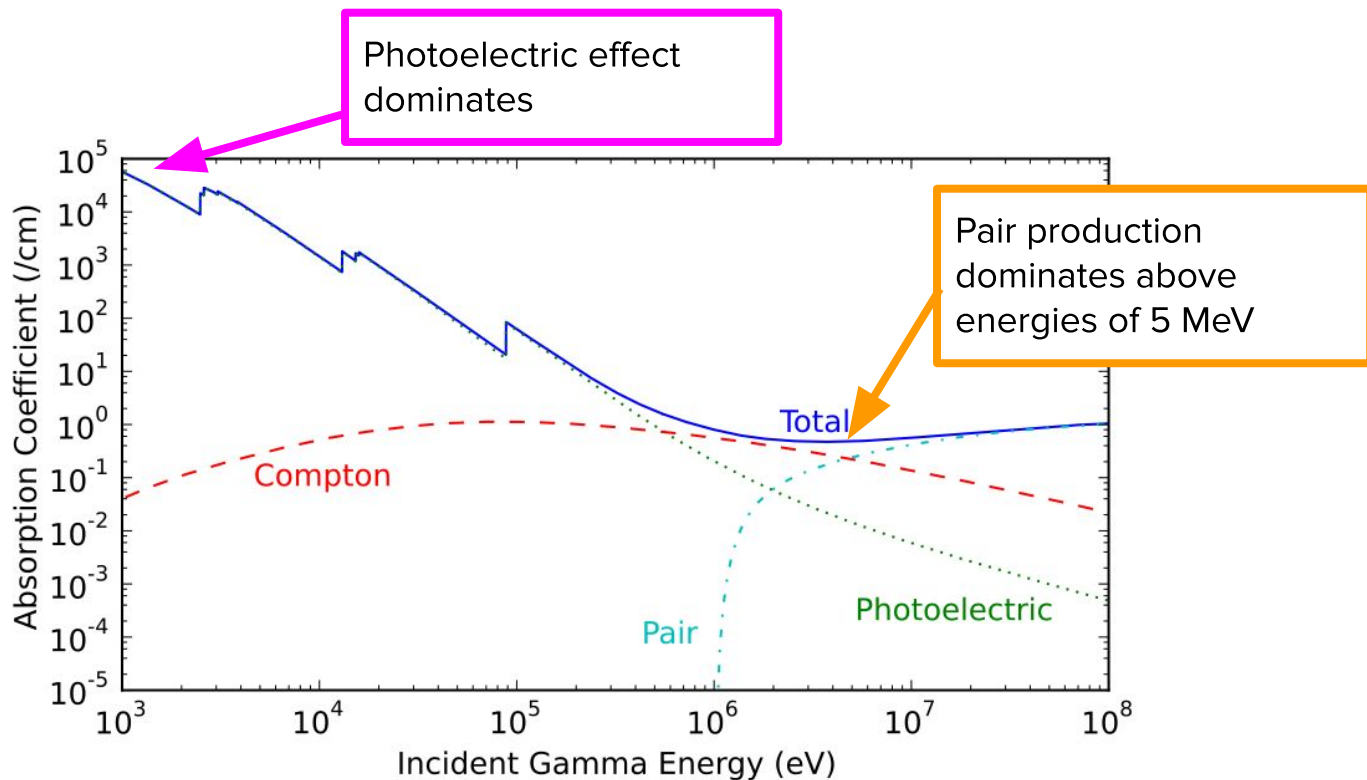


Compton
scattering



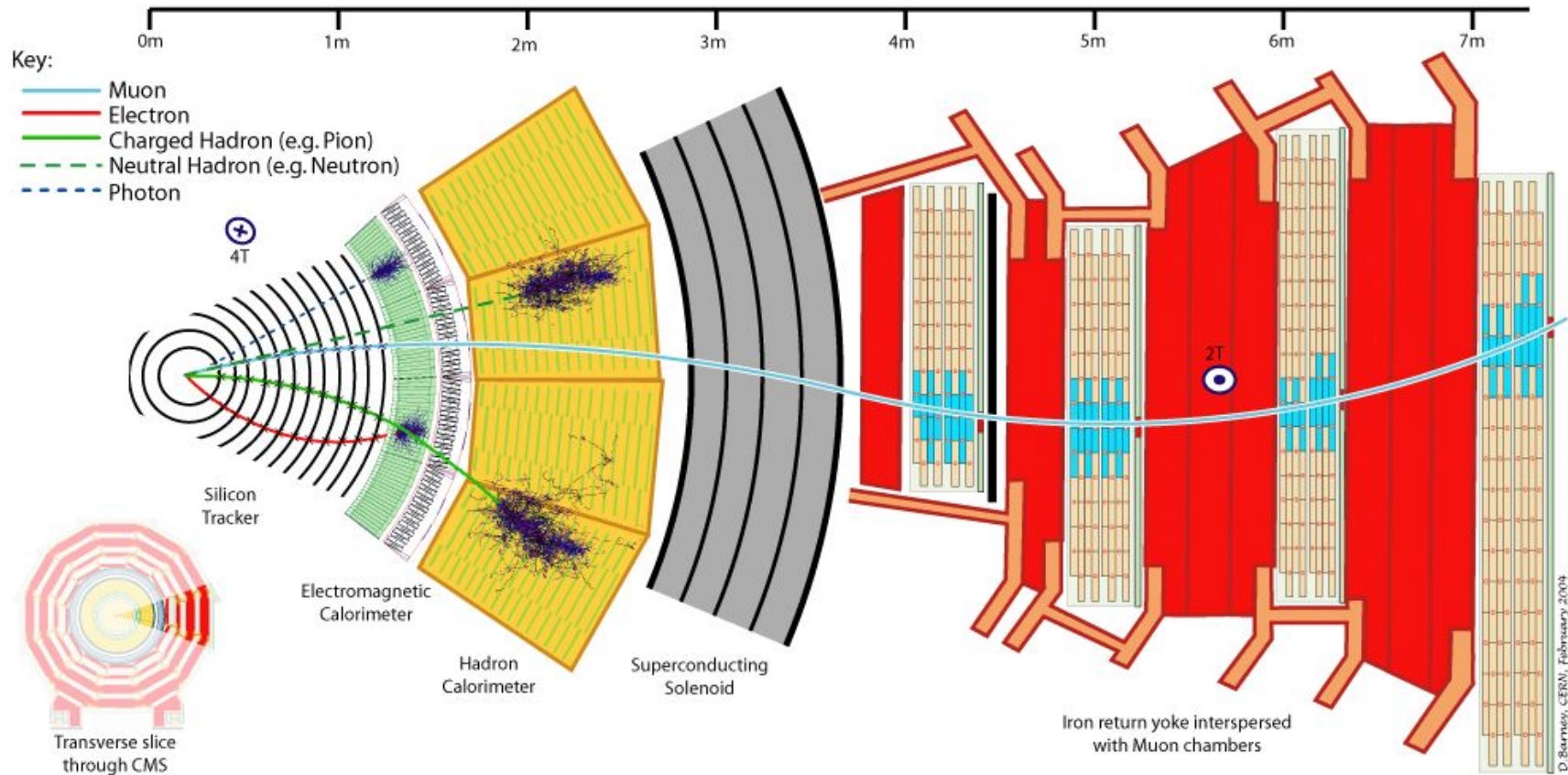
Different photon interactions in lead

See also the particle data group review on the [passage of radiation through matter](#)



Intended energy loss: using matter to
detect particles

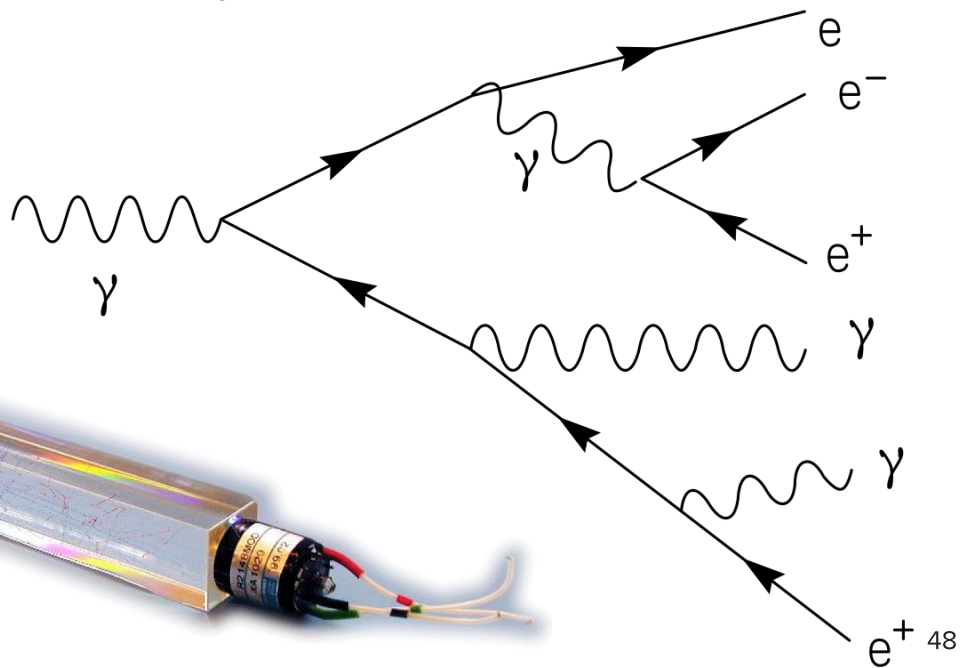
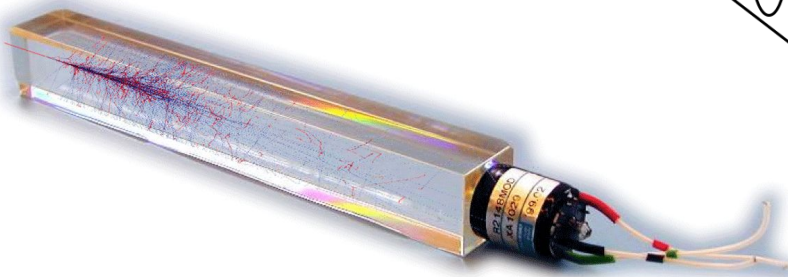
Where are the “destructive measurements”?



Electromagnetic calorimeter

- Electromagnetic shower by interaction with material
- Depth of shower in a material is determined by
 - Energy
 - Critical energy where Bremsstrahlung rate = ionization rate
 - Radiation length of material

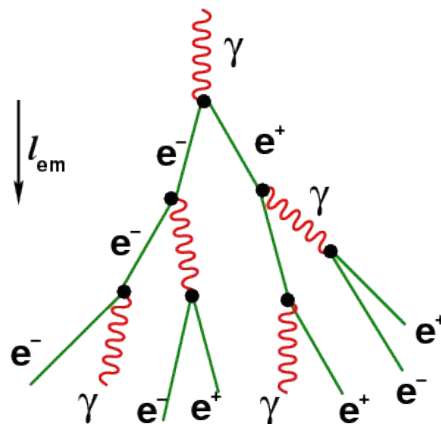
$$X = X_0 \frac{\ln(E_0/E_c)}{\ln 2}$$



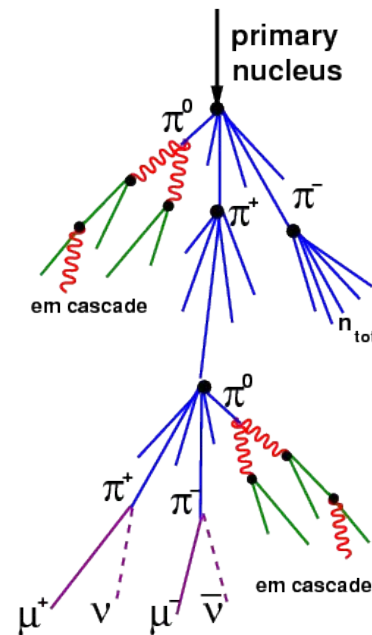
Hadronic calorimeter

| | λ_{int} [cm] | X_0 [cm] |
|--------|-----------------------------|------------|
| Szint. | 79.4 | 42.2 |
| LAr | 83.7 | 14.0 |
| Fe | 16.8 | 1.76 |
| Pb | 17.1 | 0.56 |
| U | 10.5 | 0.32 |
| C | 38.1 | 18.8 |

em cascade



hadronic cascade



$$\left. \begin{aligned} X_0 &\sim \frac{A}{Z^2} \\ \lambda_{\text{int}} &\sim A^{1/3} \end{aligned} \right\} \Rightarrow \frac{\lambda_{\text{int}}}{X_0} \sim A^{4/3}$$

$$\lambda_{\text{int}} \gg X_0$$

Hadronic calorimeters are much thicker: larger shower depth!

Energy loss: practical examples

Protons vs X-rays

You have seen: x-rays
and proton beams for
curing cancer

*Curing Cancer with
Proton Beams – with
Suzie Sheehy*

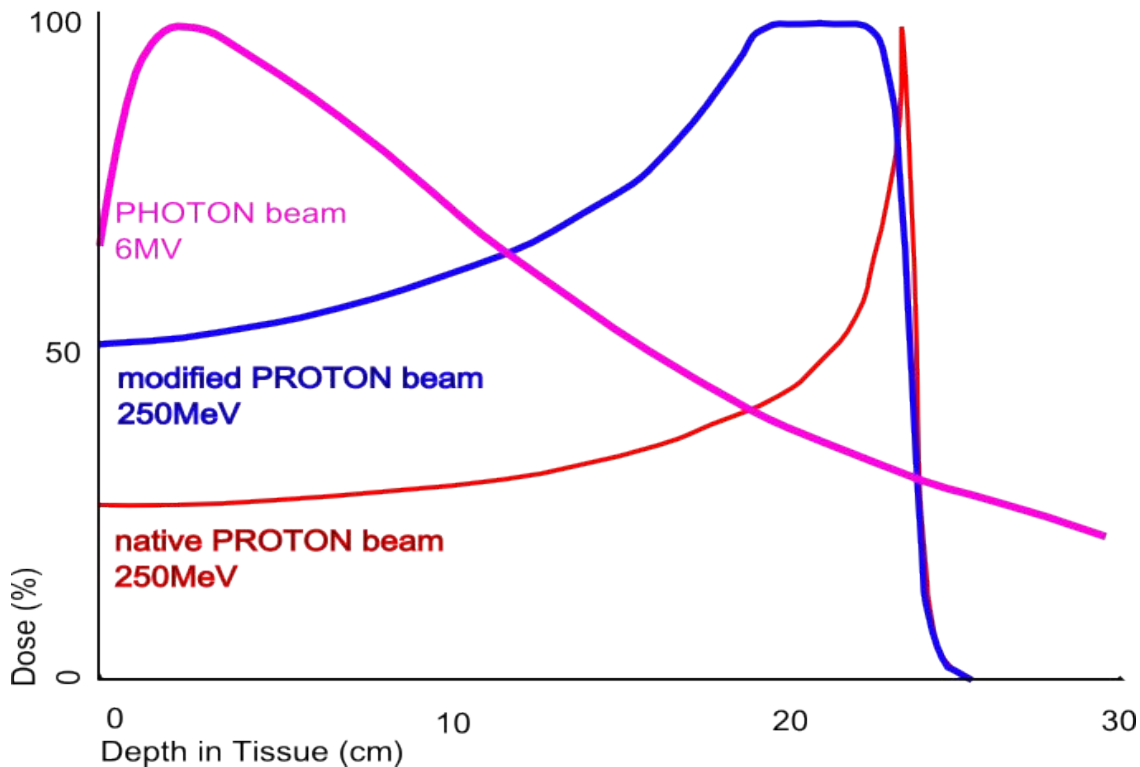
[https://www.youtube.com
/watch?v=ZQ7kyocqjiho](https://www.youtube.com/watch?v=ZQ7kyocqjiho)

Why do protons stop and
X-rays don't? What are
X-rays?



Bragg peak

As the particle loses more energy on its path, the higher the probability of interaction with the material.



Minimum ionizing particle

Typically, 23000 e^- are collected in a 300 μm silicon detector from a minimum ionizing particle.

A 388 MeV muon is a minimum ionizing particle -- using $v = 0.96 c$ for where particles become minimum ionizing.

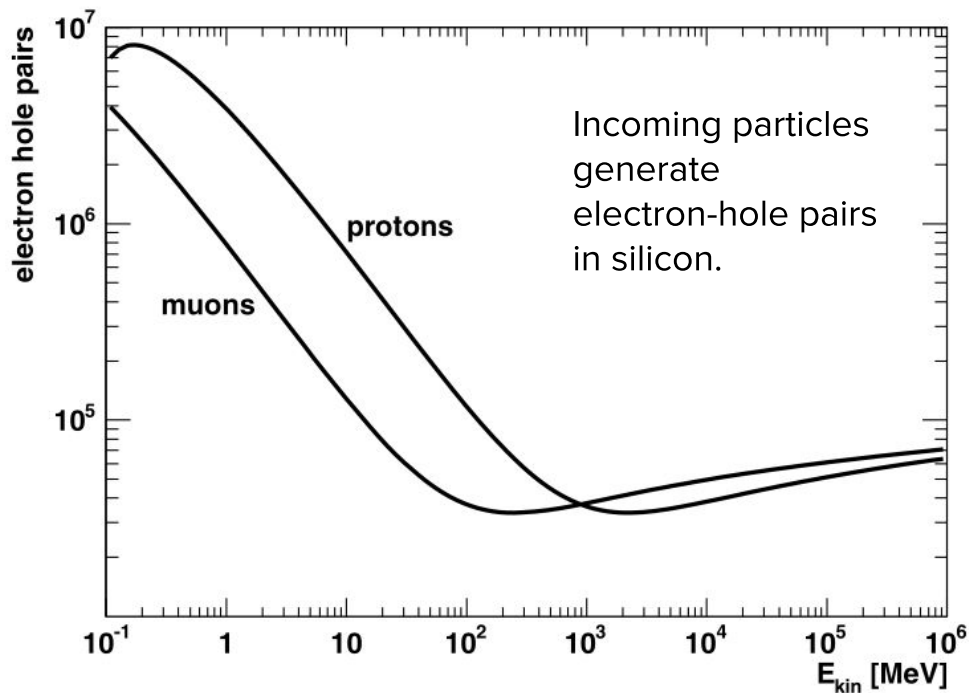


Fig. 2.2. Number of electron-hole pairs generated in a 300- μm -thick silicon layer

Unintended energy loss: we don't want matter when detecting particles!

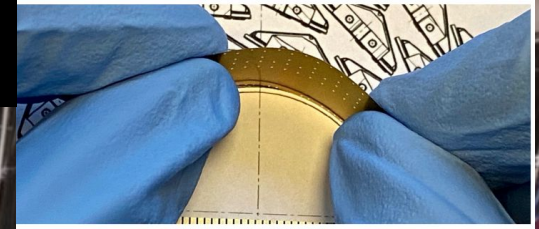
Multiple scattering

- Material is used for detection but also deflects charged particles
- Mainly through Coulomb force, also strong interaction for hadrons
- X_0 or radiation length indicates
- The radiation length X_0 of silicon is 9.36 cm \leftrightarrow 21.82 g/cm²

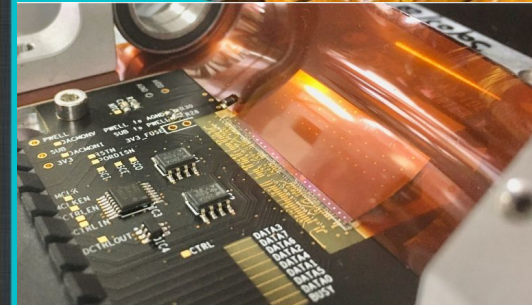
$$\theta_{\text{plane}}^{\text{rms}} = \frac{13.6 \text{ MeV}}{\beta pc} z \sqrt{\frac{x}{X_0}} \left[1 + 0.038 \ln \left(\frac{x}{X_0} \right) \right],$$

A pixel detector built for LHC has a thickness of about 2% of a radiation length per layer, changing the trajectory of a 1 GeV particle by an angle rms of $\approx 0.1^\circ$.

ALICE inner tracking system: towards 0.1% X_0



Nog
dichter
bij de
oerknal



Het goud is overigens geen goud, maar polyimide-folie met ragdunne koperen voedingskabels voor de sensoren. Dun genoeg om vrijkomende

zijn ontstaan. ITS moet de deeltjes betrappen die uit die ziedende oersoep ontsnappen en de fysici vertellen wat er daarbinnen precies gaande is.

geven, is de verwachting. De nieuwe inner tracker doet dat van nog dichterbij dan de eerdere versie. Daardoor kunnen de kernbotsingen preciezer worden bekeken.

dezelfde plak silicium zitten. Dat scheelt kabels en elektronica in de detector.

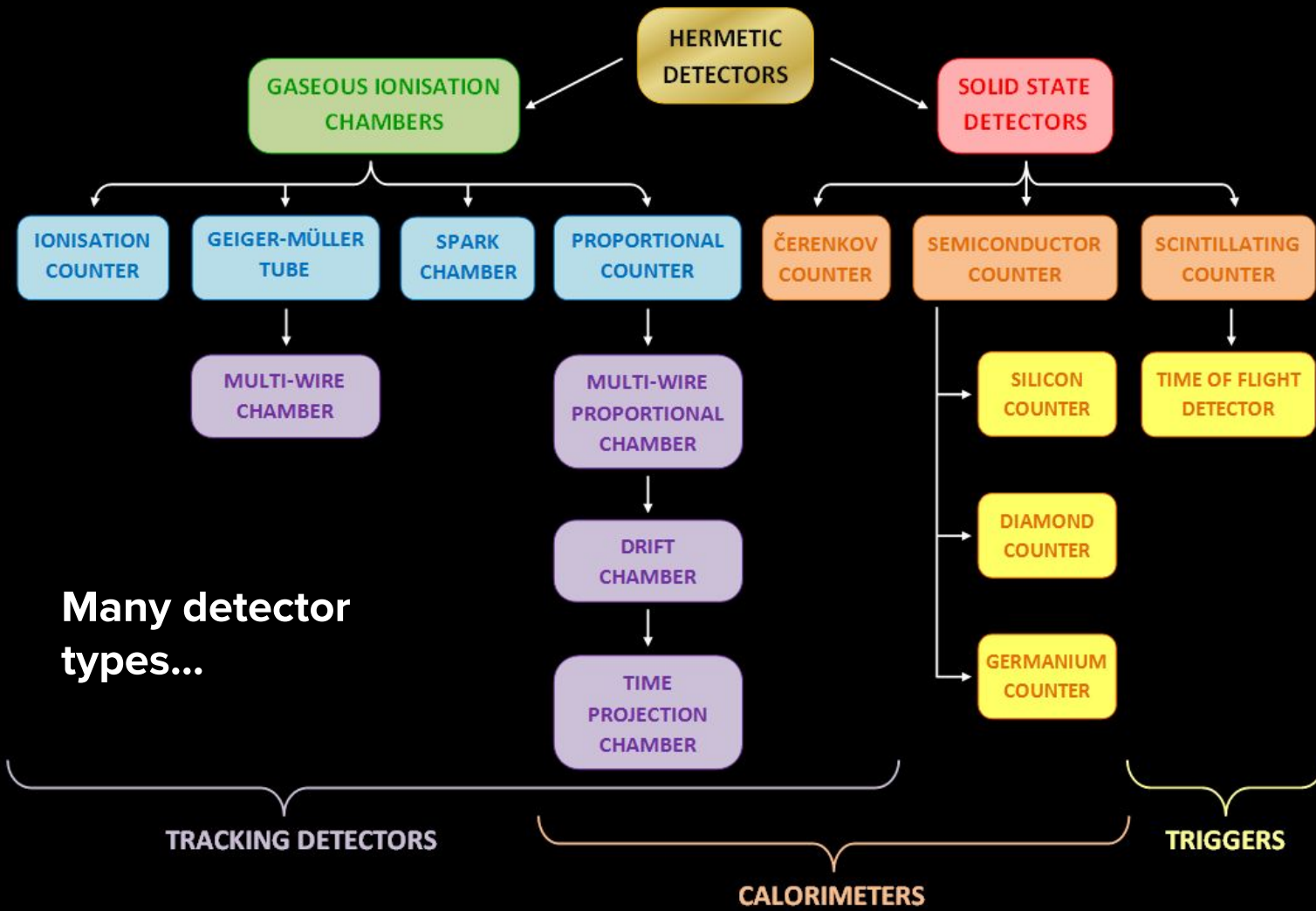
In de komende meetperiode kan ITS gemakkelijk honderd keer zoveel meetgegevens verzamelen als alles wat ALICE in

ne versies worden vervangen. Daar is ook de trigger-apparatuur bij die beslist welke botsingen bijzonder genoeg zijn om vast te leggen. Het computersysteem dat data verzamelt en toegankelijk maakt, wordt eveneens vernieuwd.

De upgrade-periode is een hectische tijd. Het binnenste van de grote ondergrondse detector is vorig jaar meteen

kleine honderd sensorduigen. Een kwart van alle duigen, die in de lagen nummer 6 en 7, zijn gemaakt op Nikhef in Amsterdam. Daar lijmden leden van het ALICE-team met eindeloos geduld de koeling en de sensoren stuk voor stuk handmatig op de ijle koolstofvezel dragers. Deze sensorduigen zijn vorig najaar al in trillingsvrije kratten van Amsterdam

Passage of radiation through matter
can be used to measure as well as to
stop particles



Many detector types...

https://upload.wikimedia.org/wikipedia/commons/c/c0/Detectors_summary_3.png

**Stay
tuned!**

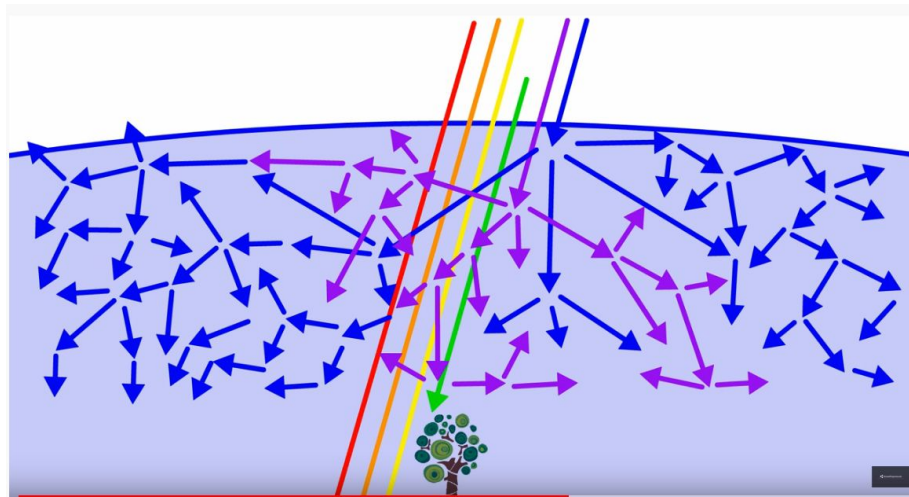
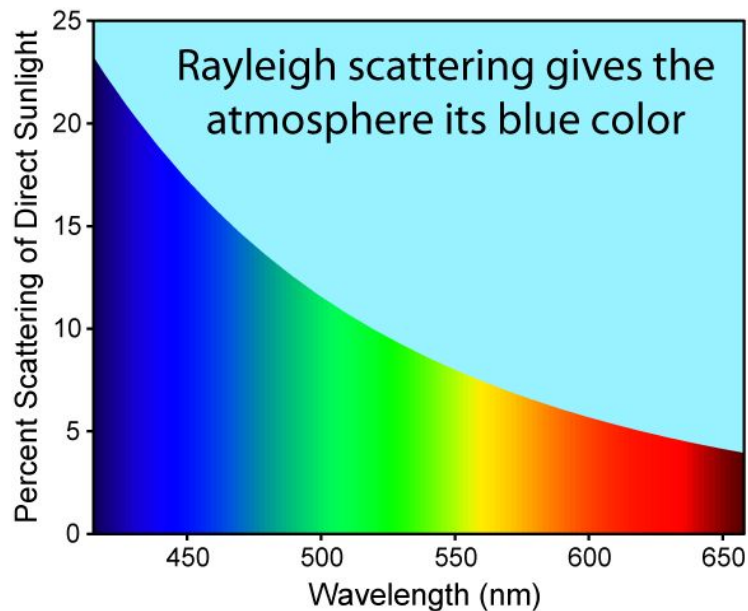
Additional material

Other sources used

- Slides from Erik Butz (see also [here](#)), [Simon Spannagel](#), [Freya Blekman](#), [Peter Schleper](#), [Erika Garutti](#), [Ingrid-Maria Gregor](#)
- Book on [Pixel detectors](#) by Rossi, Fischer, Rohe, Wermes
- Book on [Semiconductor Detector Systems](#) by Helmuth Spieler, see also [online notes](#)
- Particle data group review: [passage of radiation through matter](#)
- Book on [particle radiation](#) by Konrad Kleinknecht
- Book on [Techniques in experimental particle physics](#) by WR Leo

See also links behind pictures in the slides for more sources.

Remark about photons: blue sky



Rayleigh scattering cross section:

$$\sigma_s = \frac{2\pi^5}{3} \frac{d^6}{\lambda^4} \left(\frac{n^2 - 1}{n^2 + 2} \right)^2$$

$$\frac{\sigma_{blau}}{\sigma_{rot}} = \frac{1/\lambda_{blau}^4}{1/\lambda_{rot}^4} = \left(\frac{650 \text{ nm}}{450 \text{ nm}} \right)^4 \approx 4.4$$